

Few-fermion thermometry

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Tomasz Sowiński

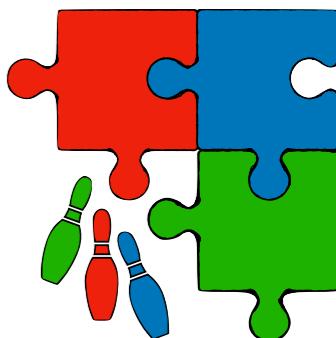
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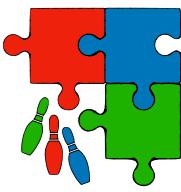
FEW-BODY PROBLEMS
THEORY GROUP



FewBody.ifpan.edu.pl



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motivation

Selim Jochim's experimental group
in Heidelberg

6Li	
Atomic number (Z)	3
Nucleons (Z+N)	6
Total electronic spin (S)	1/2
Total nuclear spin (I)	1
Hyperfine states (F=S+I)	1/2 or 3/2

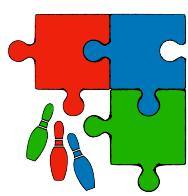
$$\begin{aligned}|\uparrow\rangle &= |F = 3/2, m_F = -3/2\rangle \\ |\downarrow\rangle &= |F = 1/2, m_F = 1/2\rangle\end{aligned}$$

message from experiments

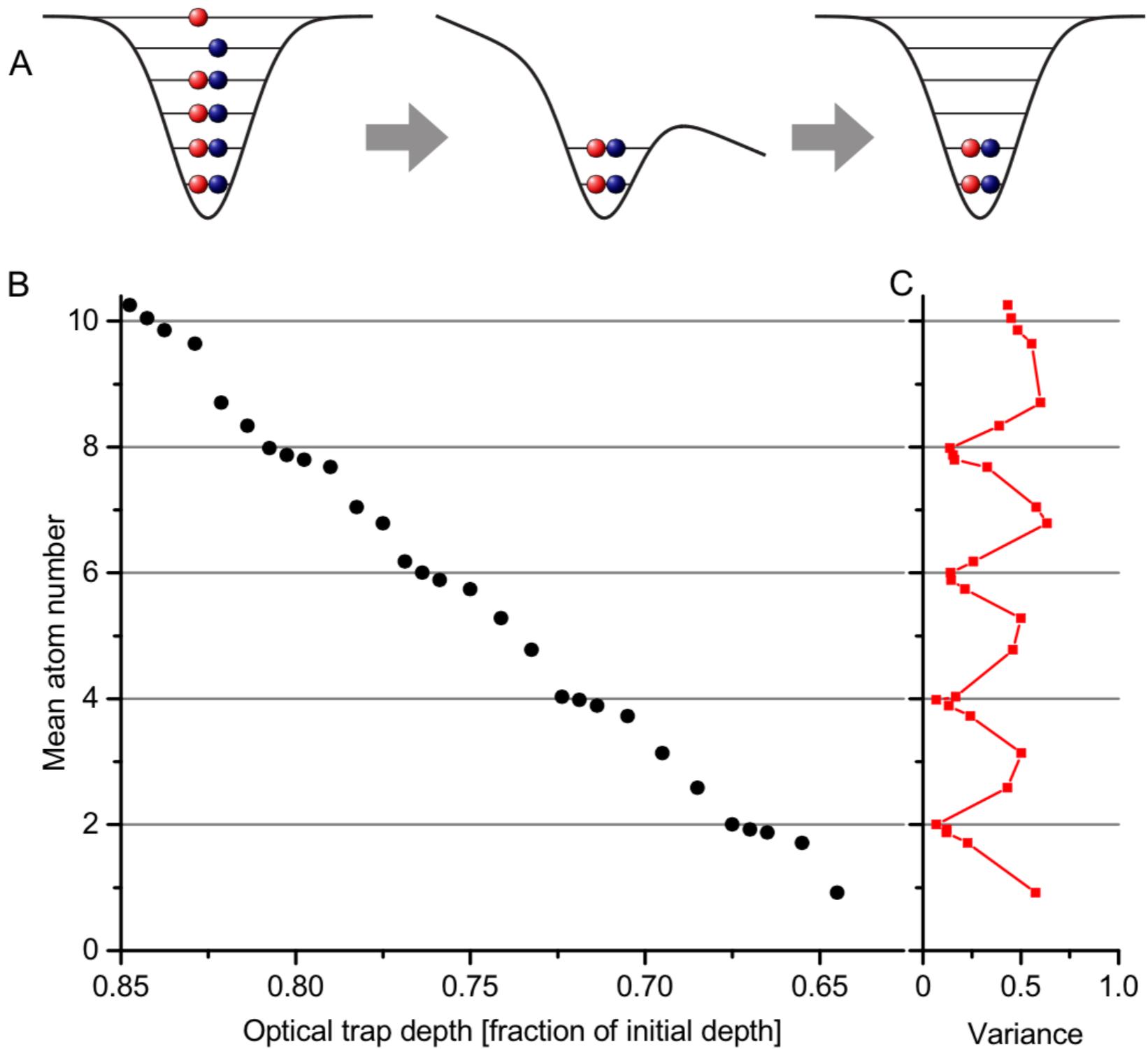
- two flavours of spinless fermions
- well controlled number of particles
- one-dimensional confinement
- strong interactions between flavours

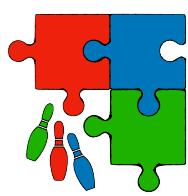
... what for???

- extreme control on the system
- extreme precision of measurements
- from 'one' to 'many' crossover

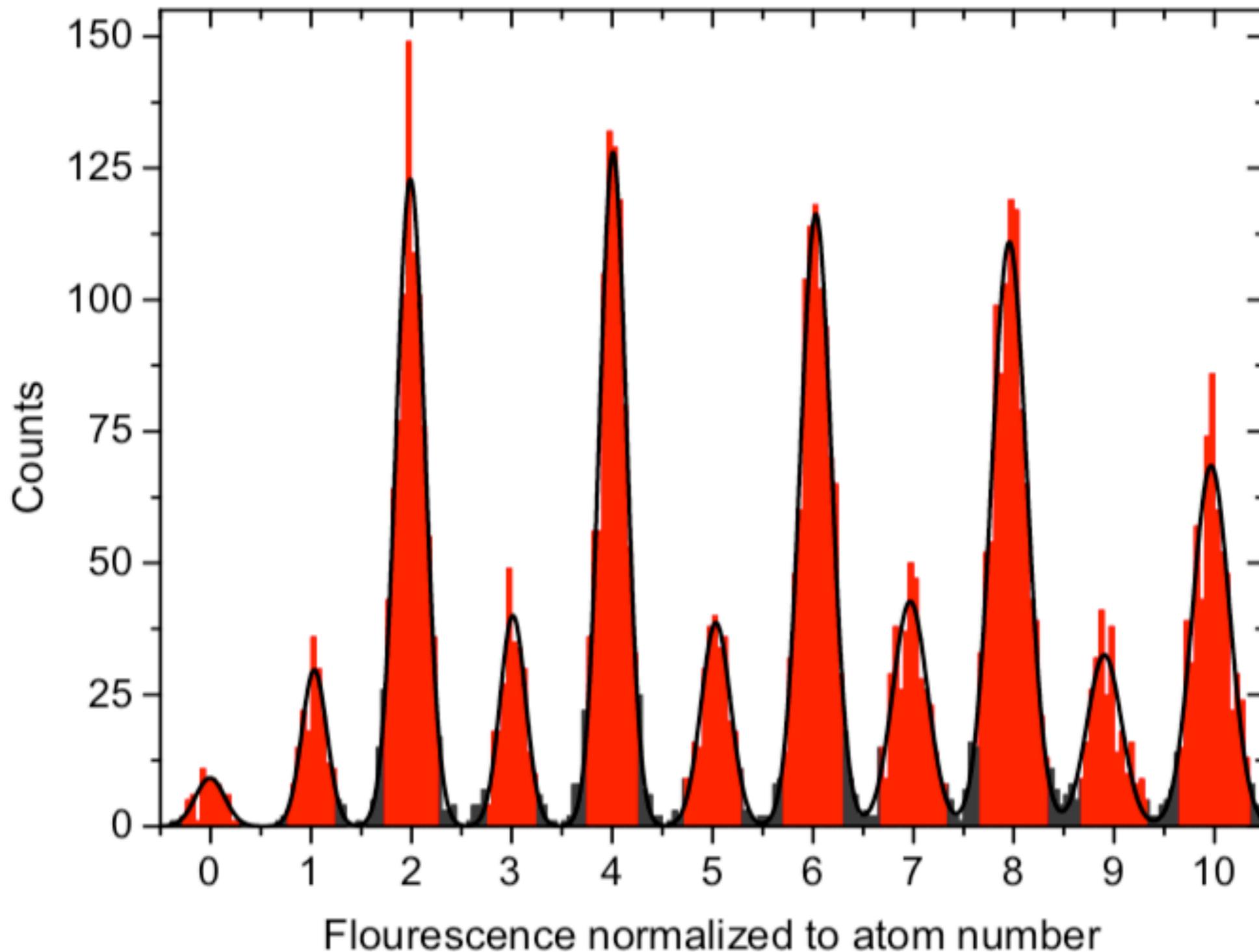


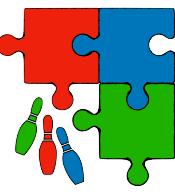
message from experiments





message from experiments





the model

- two distinguishable flavours of fermions (\uparrow and \downarrow)
- both flavours may have different masses
- both flavours confined in one-dimensional trap
- opposite spins do interact via sort range δ -like potential

$$\hat{\mathcal{H}} = \sum_{\sigma} \int dx \hat{\Psi}_{\sigma}^{\dagger}(x) \left[-\frac{\hbar^2}{2m_{\sigma}} \frac{d^2}{dx^2} + V_{\sigma}(x) \right] \hat{\Psi}_{\sigma}(x)$$
$$+ g \int dx \hat{\Psi}_{\downarrow}^{\dagger}(x) \hat{\Psi}_{\uparrow}^{\dagger}(x) \hat{\Psi}_{\uparrow}(x) \hat{\Psi}_{\downarrow}(x)$$

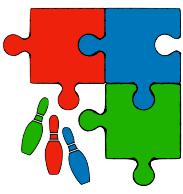
(anti-)commutation relations

$$\left\{ \hat{\Psi}_{\sigma}(x), \hat{\Psi}_{\sigma}^{\dagger}(x') \right\} = \delta(x - x')$$

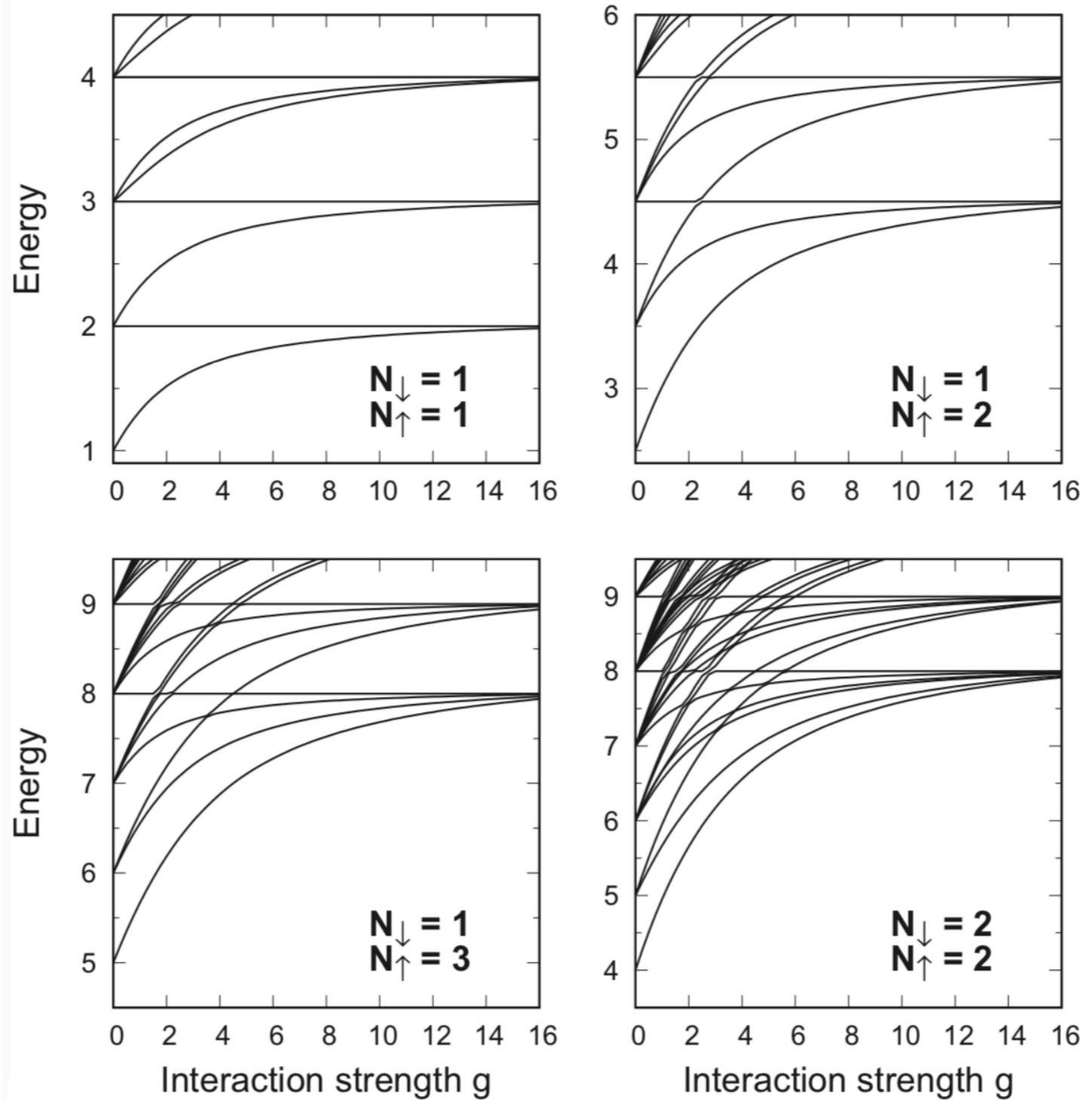
$$\left\{ \hat{\Psi}_{\sigma}(x), \hat{\Psi}_{\sigma}(x') \right\} = 0$$

$$[\hat{\mathcal{N}}_{\uparrow}, \hat{\mathcal{H}}] = [\hat{\mathcal{N}}_{\downarrow}, \hat{\mathcal{H}}] = 0$$

$$\hat{\mathcal{N}}_{\sigma} = \int dx \hat{\Psi}_{\sigma}^{\dagger}(x) \hat{\Psi}_{\sigma}(x)$$

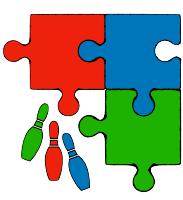


many-body spectrum



**Notice the quasi-degeneracy
of the spectrum for strong repulsions!**

**Utilise it
for precise thermometry!**



thermal sensitivity

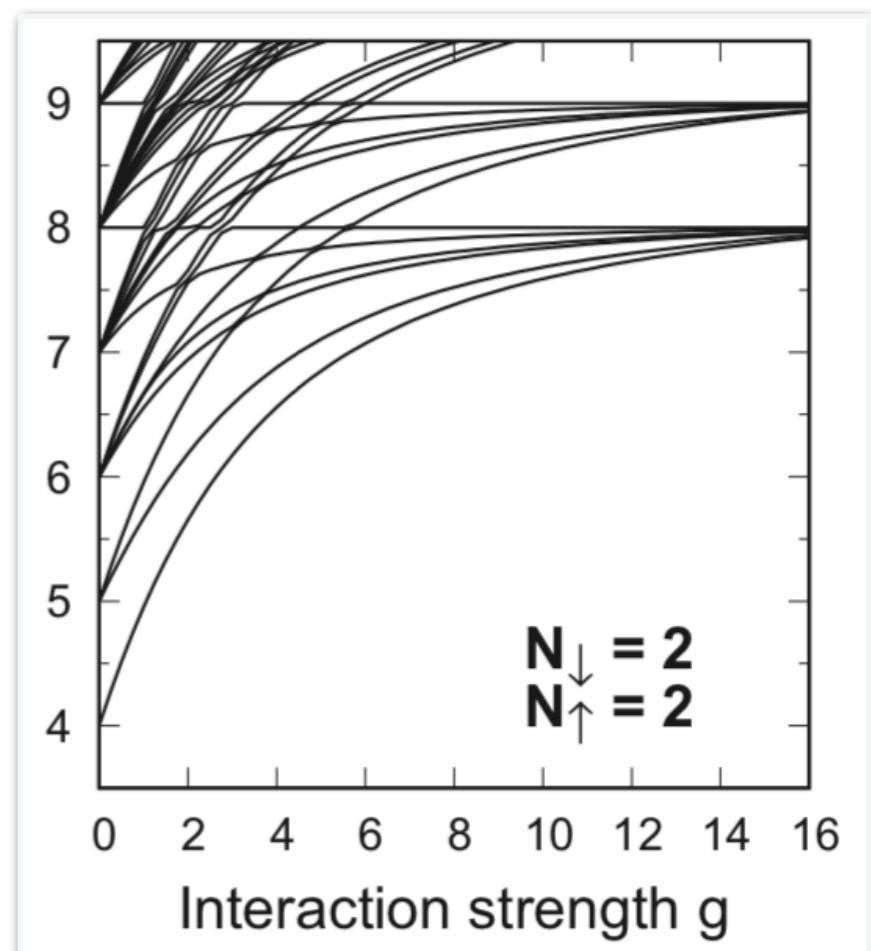
Let us assume that the system is in thermal equilibrium with some other system of unknown temperature T

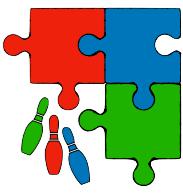
$$\hat{\rho}_T = \frac{1}{Z} e^{-\hat{\mathcal{H}}/k_B T} = \sum_i P_i |E_i\rangle\langle E_i|$$

$P_i = \frac{1}{Z} e^{-E_i/k_B T}$

Question: how precisely one can determine temperature T by performing measurements on the system?

Answer: it depends on measurements one has in the arsenal!





thermal sensitivity

$$\hat{\rho}_T = \frac{1}{Z} e^{-\hat{\mathcal{H}}/k_B T} = \sum_i P_i |\mathcal{E}_i\rangle\langle\mathcal{E}_i|$$

Measurements: Some set of probabilities as functions of T.

The limit is determined by the Fisher information.

$$\{p_i(T)\} \rightarrow \mathcal{F}(\{p_i\}) = \sum_i \frac{1}{p_i} \left(\frac{dp_i}{dT} \right)^2$$

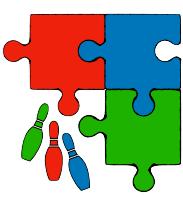
$$\Delta T^2 \geq \frac{1}{n \mathcal{F}(\{p_i\})}$$

IF: One can measure Gibbs probabilities

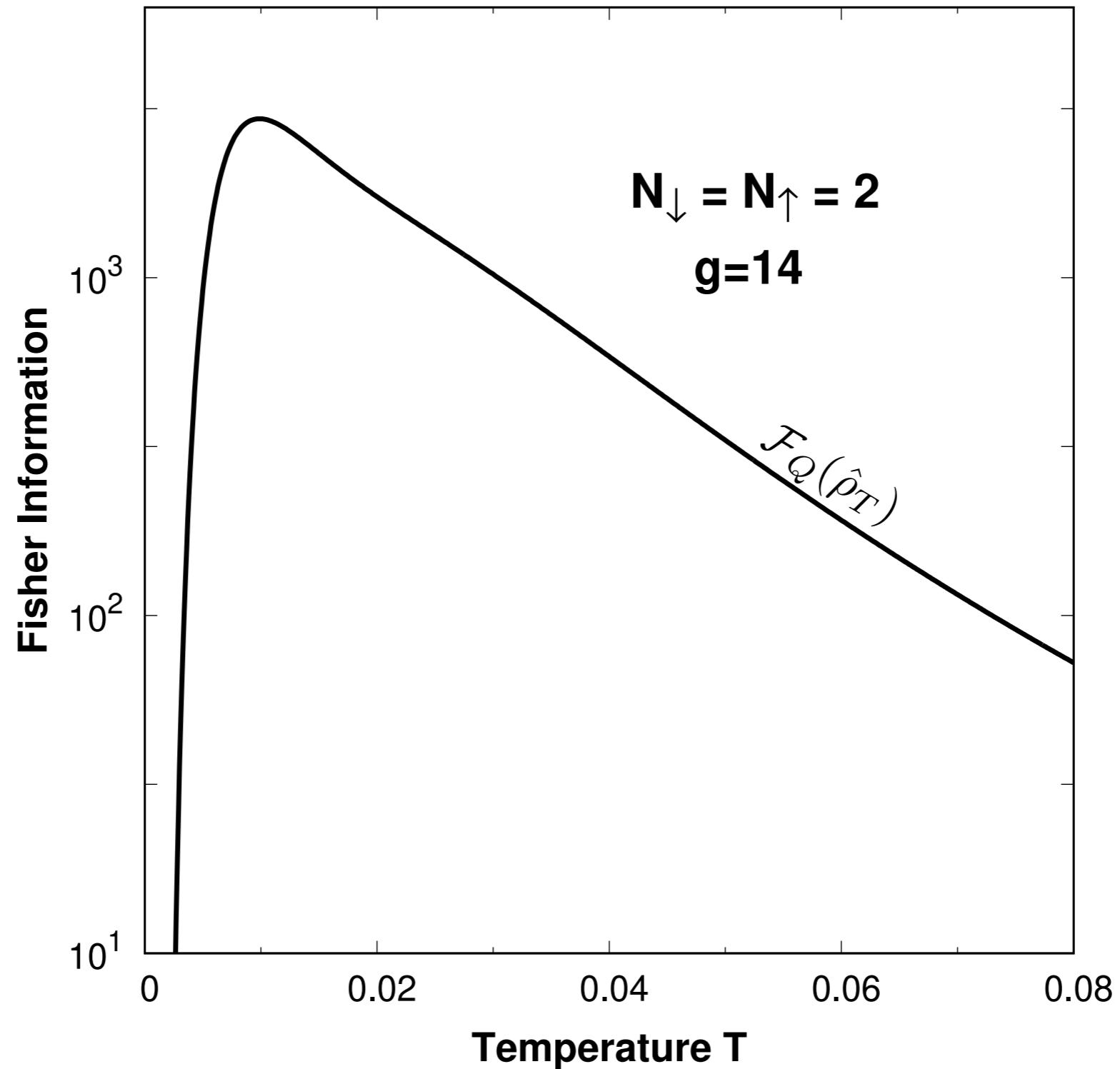
$$\{p_i(T)\} = \{P_i(T)\}$$

then the Fisher information is the largest possible!

$$\mathcal{F}_Q(\hat{\rho}_T) = \frac{\Delta \hat{\mathcal{H}}^2}{T^4} = \frac{\text{Tr}[\hat{\rho}_T \hat{\mathcal{H}}^2] - \text{Tr}[\hat{\rho}_T \hat{\mathcal{H}}]^2}{T^4}$$

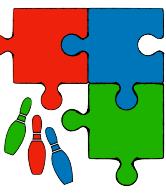


quantum Fisher information

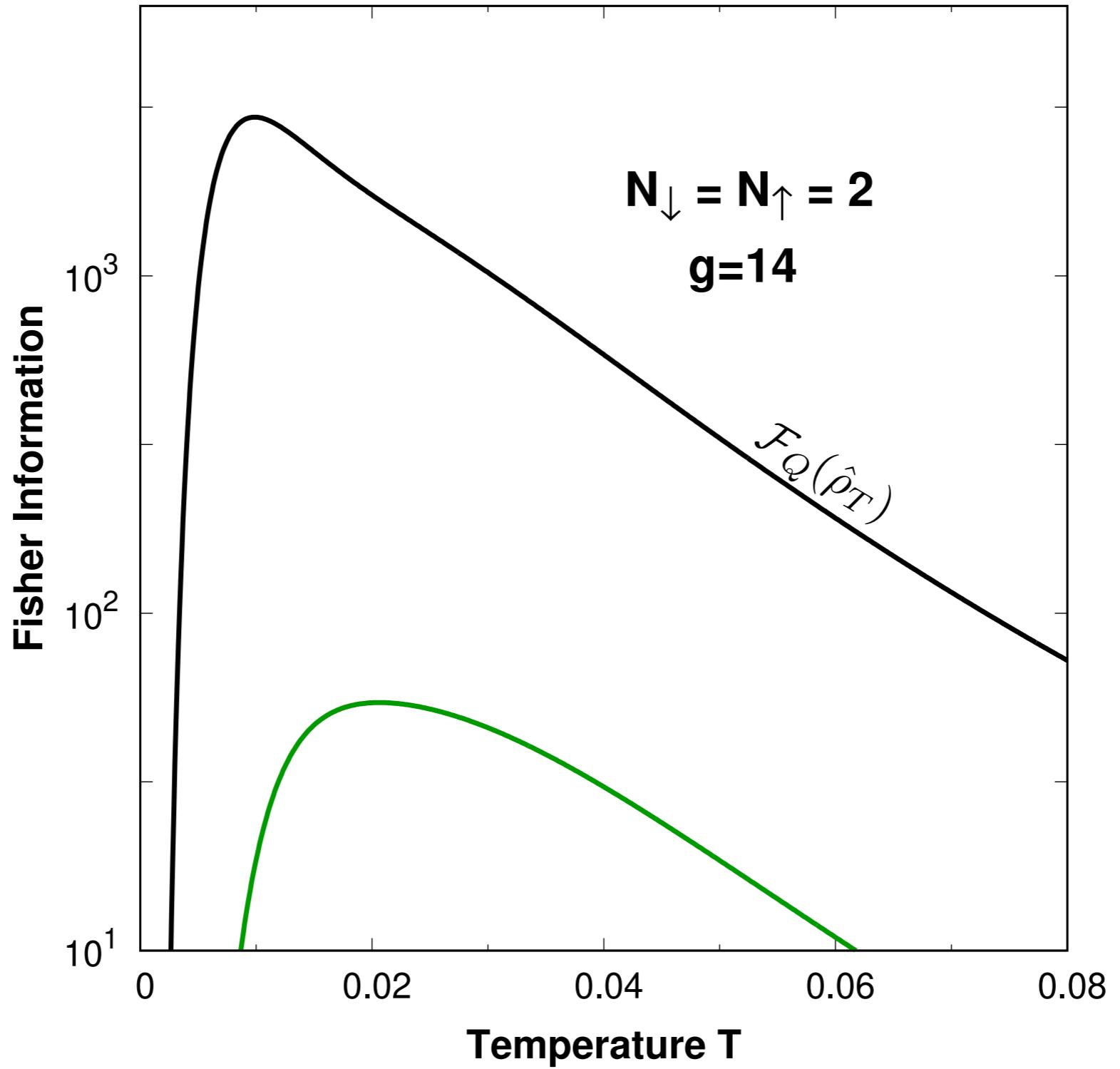


$$\mathcal{F}_Q(\hat{\rho}_T) = \frac{\Delta \hat{\mathcal{H}}^2}{T^4} = \frac{\text{Tr}[\hat{\rho}_T \hat{\mathcal{H}}^2] - \text{Tr}[\hat{\rho}_T \hat{\mathcal{H}}]^2}{T^4}$$

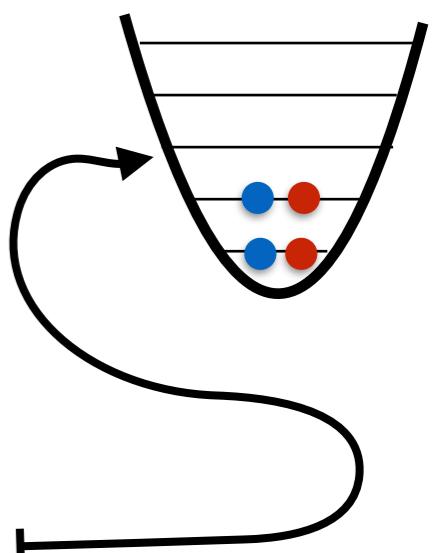




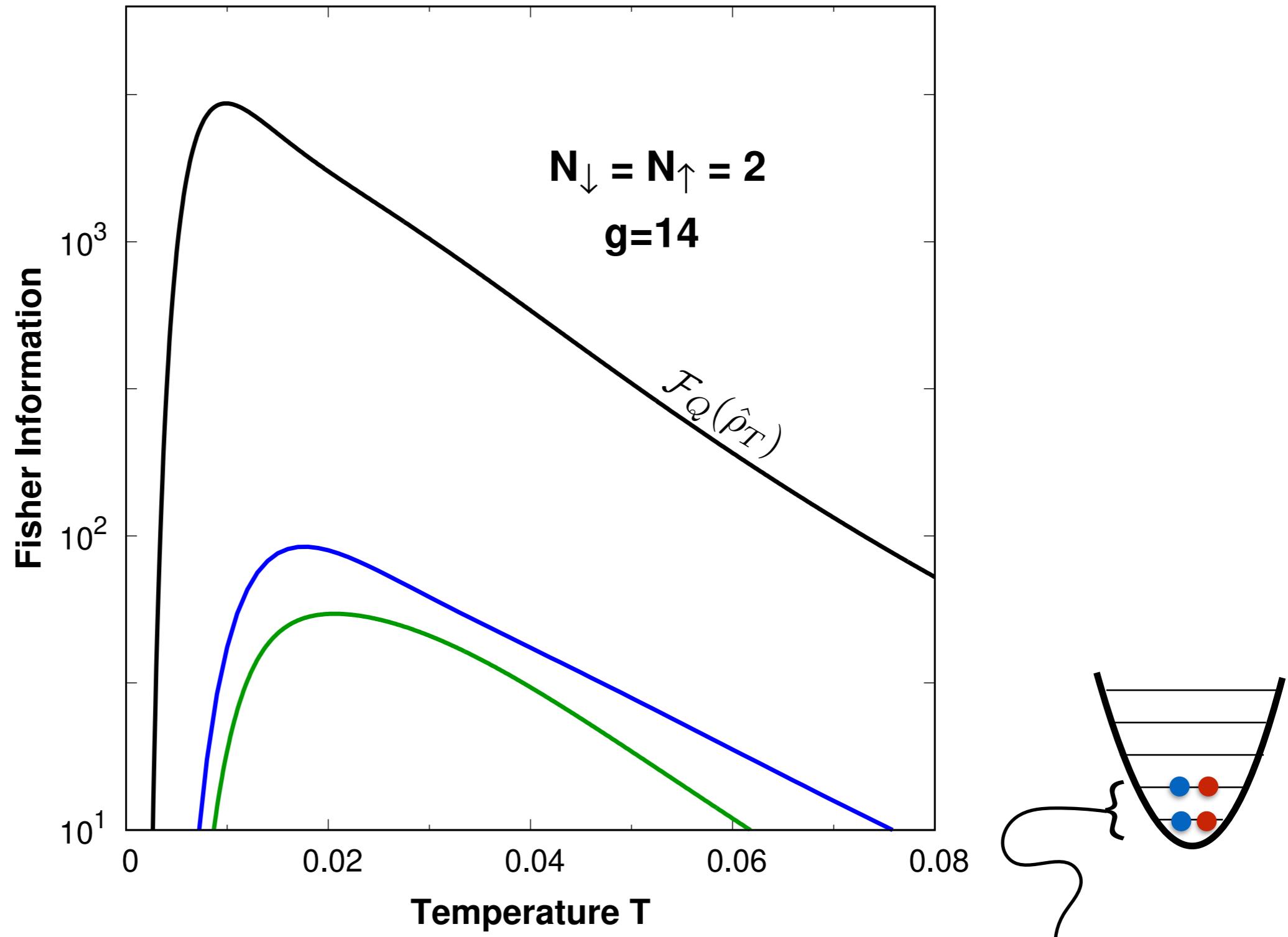
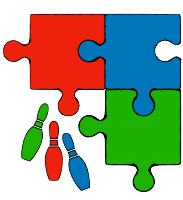
the simplest scenario



The probability of finding a particle just above the Fermi level ↪

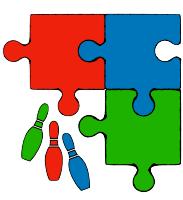


Occupations of the lowest orbitals

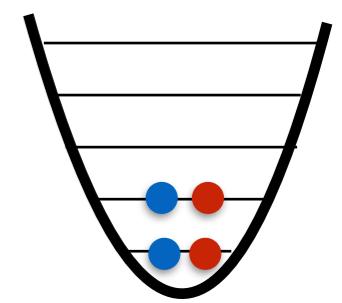
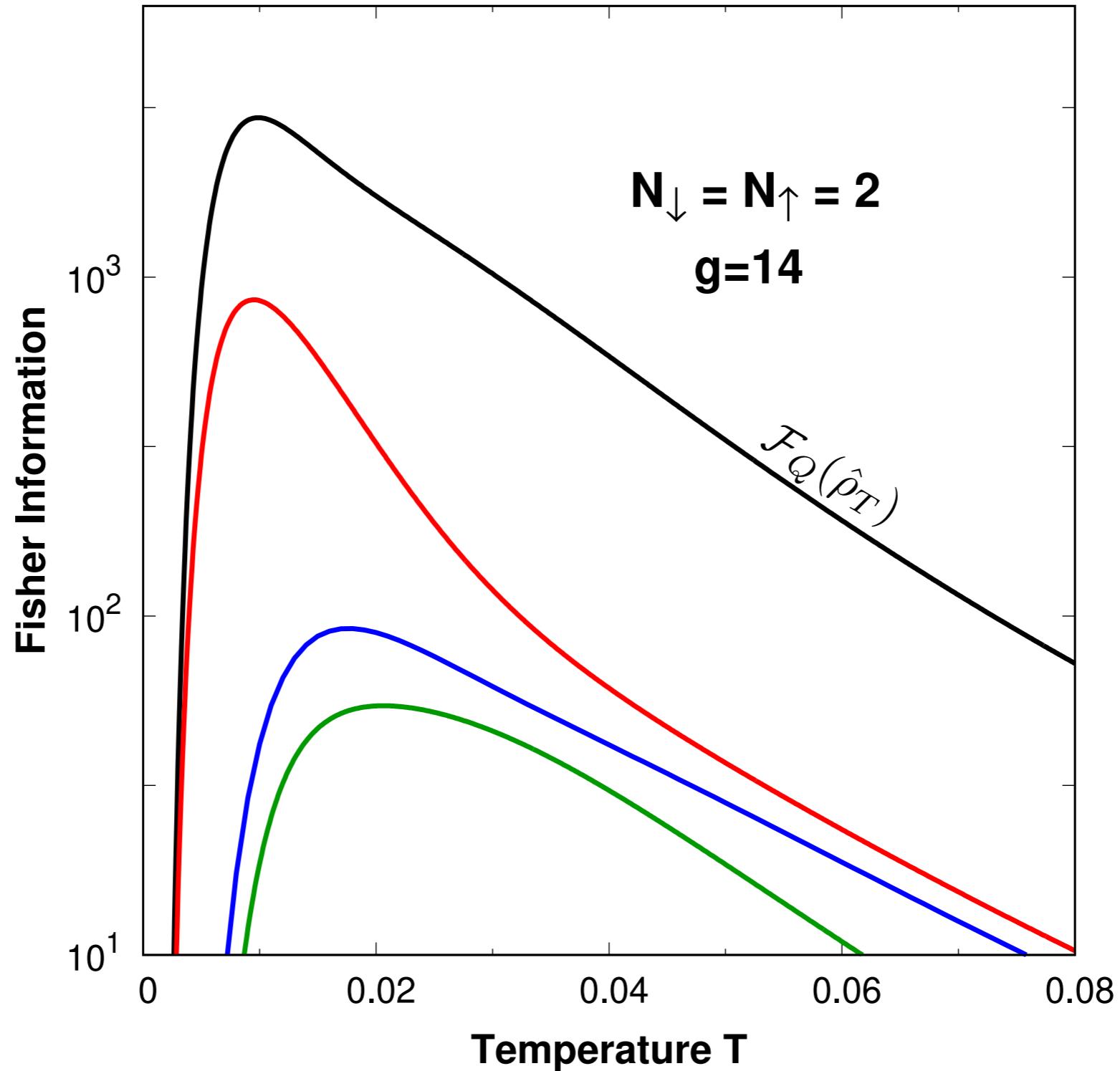


Probabilities p_0, p_1, p_2, p_3, p_4 that in two the lowest orbitals one can find
0, 1, 2, 3, or 4 particles.





All occupations are known

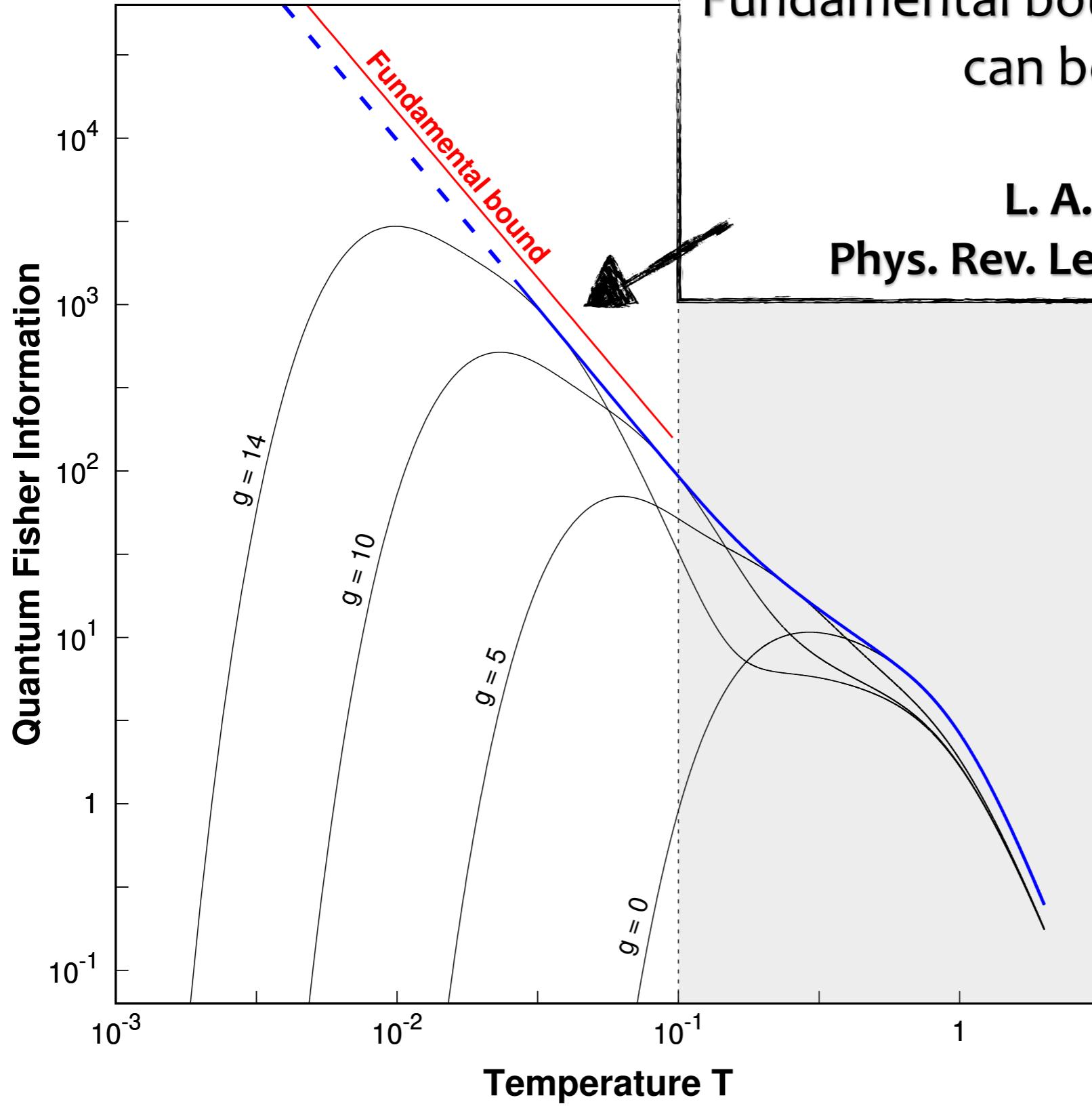


When we know **all** probabilities
of finding **four particles** in different configurations



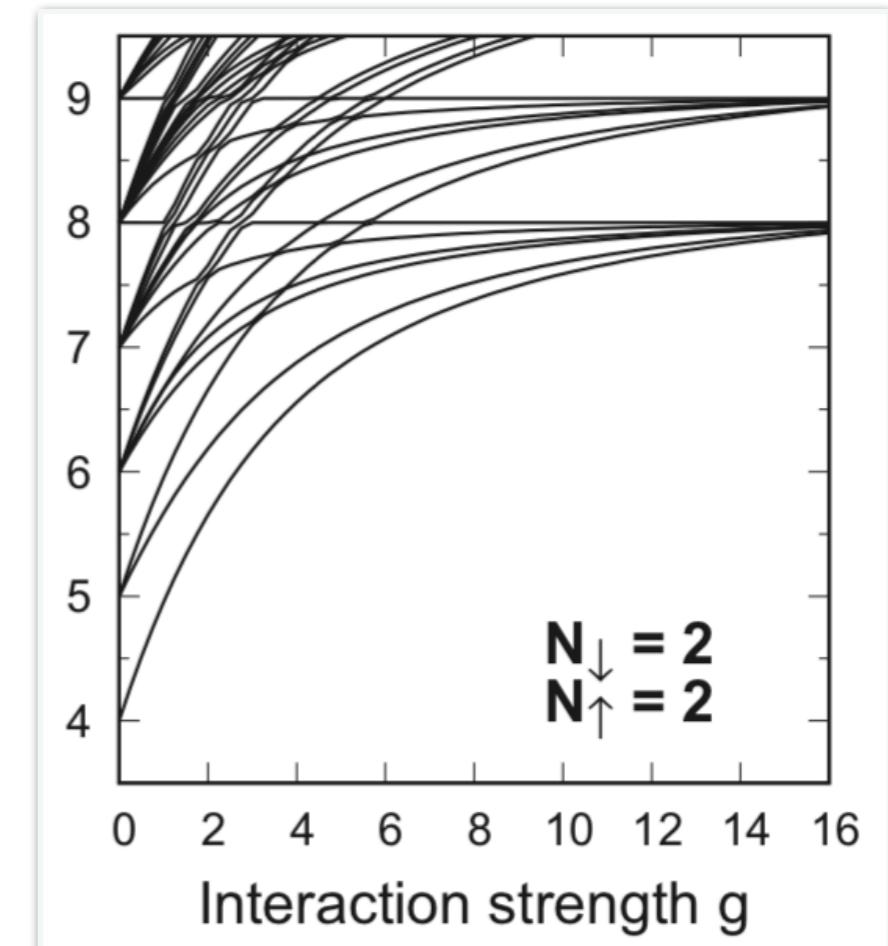


the fundamental bound



Fundamental bound when level spacings
can be engineered

L. A. Correa et al.
Phys. Rev. Lett. 114, 220405 (2015).



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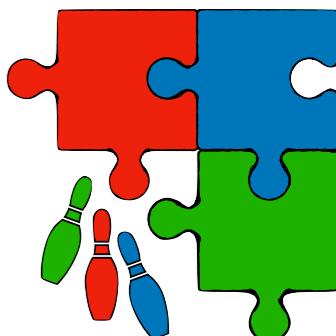
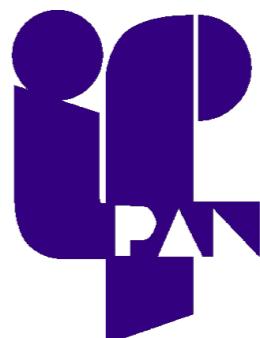
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