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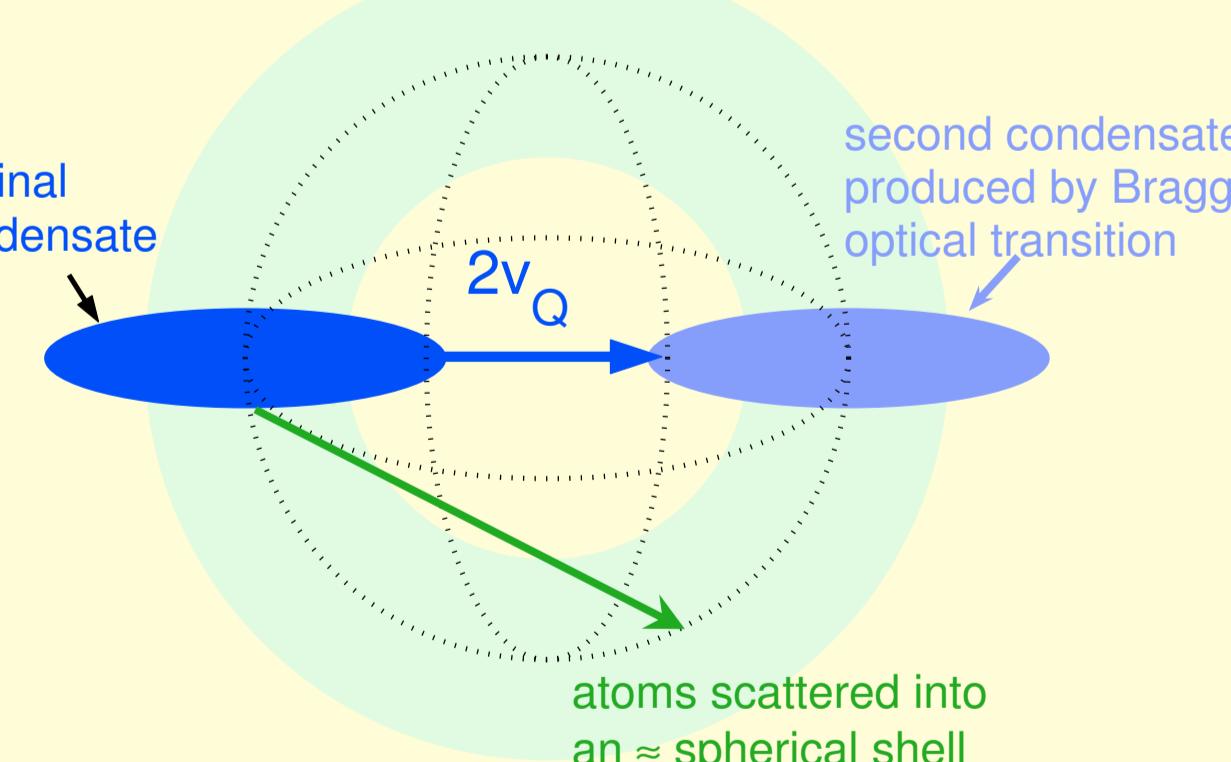
Quantum dynamics: THE ISSUE

- **COMPLETE METHODS** – Intractable or run for a limited time.
(*positive-P, stochastic wavefunction,...*)
- **TRACTABLE METHODS** – Incomplete (Sometimes appreciably, sometimes not).
(*GP, PGPE, c-field,...*)

A (partial) Remedy

1. Suppose we have:
 - A complete method \mathcal{Q} that is tractable only for a short time.
 - Approximate methods $\mathcal{A}, \mathcal{B}, \dots$ that are tractable for long times.
2. Suppose \mathcal{A} and \mathcal{B} are similar enough to \mathcal{Q} that an ad-hoc *continuous* blend of the evolution equations can be made:
 - $\mathcal{H}_A = \lambda \mathcal{Q} + (1 - \lambda) \mathcal{A}$
 - $\mathcal{H}_B = \lambda \mathcal{Q} + (1 - \lambda) \mathcal{B}$
 With blending parameter $\lambda \in [0, 1]$. Here $\lambda \rightarrow 1$ as the equations approach full QD.
3. A judicious choice of hybrid evolution equations produces simulations that *last much longer than the exact method \mathcal{Q}* , but have *smaller systematics than \mathcal{A} or \mathcal{B}* .
4. With λ continuous, one can extrapolate to $\lambda = 1$ obtain an estimate of the full quantum mechanical result.
5. With two different hybrids \mathcal{H}_A and \mathcal{H}_B , wild extrapolations are prevented.
6. The uncertainty of the final predictions can be well characterised from their statistics.

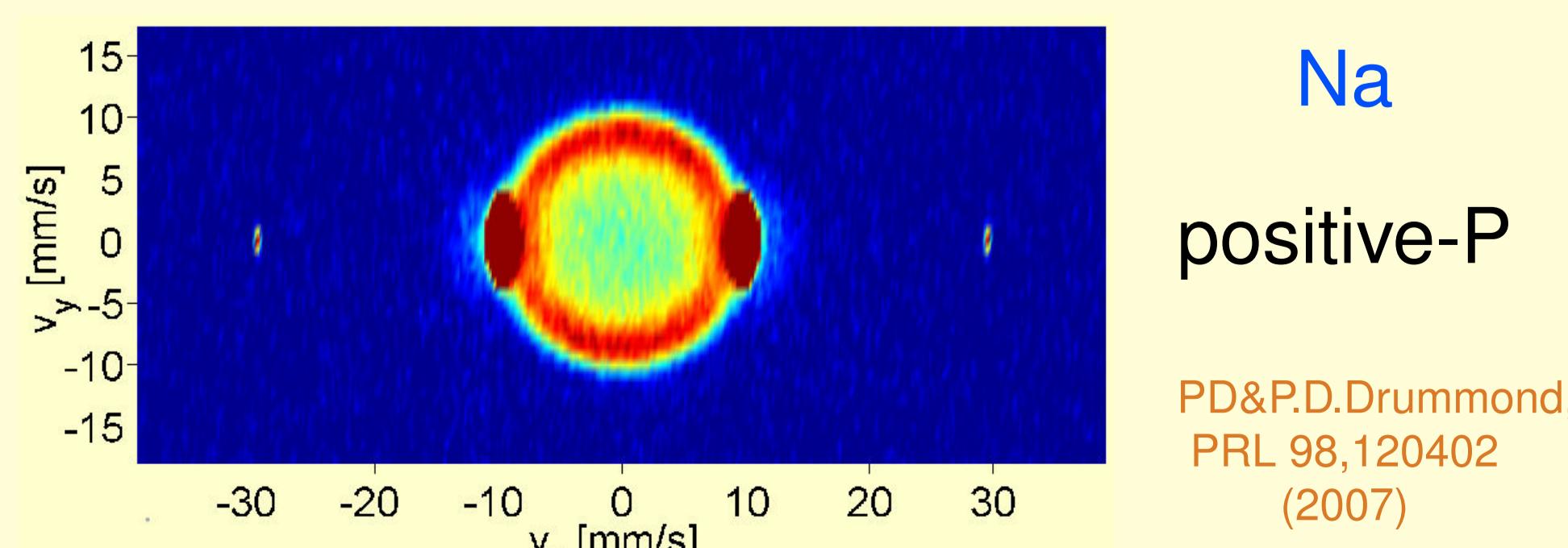
Example: BEC collision



Some experiments:

- He* A.Perrin et al., PRL 99, 150405 (2007).
- Na A.P.Chikkatur et al, PRL 85, 483 (2007). J.M.Vogels et al, PRL 89, 020401 (2007).
- Rb N.Katz, PRL 95, 220403 (2005).

$N = 1.5 \times 10^5$ atoms, lattice $\sim 10^6$ points.



Q: Complete dynamics: Positive P distribution

Random walk of samples $\psi(x)$ and $\tilde{\psi}(x)$

[Equivalent to full QD in the ∞ samples limit.]

$$i\hbar \frac{d\psi(x)}{dt} = \left[-\frac{\hbar^2}{2m} \nabla^2 + g \psi(x) \tilde{\psi}^*(x) + \sqrt{\frac{g}{i\hbar}} \xi(x, t) \right] \psi(x)$$

$$i\hbar \frac{d\tilde{\psi}(x)}{dt} = \left[-\frac{\hbar^2}{2m} \nabla^2 + g \psi(x) \tilde{\psi}^*(x) + \sqrt{\frac{-g}{i\hbar}} \tilde{\xi}(x, t) \right] \tilde{\psi}(x)$$

= mean field GP + noise fields $\xi, \tilde{\xi}$

$$\bar{n}(x) = \langle \tilde{\Psi}^\dagger(x) \tilde{\Psi}(x) \rangle = \langle \tilde{\psi}^*(x) \psi(x) \rangle_{\text{ensemble}}$$

A : Truncated Wigner

GP evolution + $\frac{1}{2}$ virtual particle in IC

$$\psi(x, 0) = \phi_{GP}(x, 0) + \frac{\eta(x)}{\sqrt{2}} \quad \eta(x) : \text{gaussian noise}$$

$$i\hbar \frac{d\psi(x)}{dt} = \left[-\frac{\hbar^2}{2m} \nabla^2 + g |\psi(x)|^2 \right] \psi(x)$$

$$\bar{n}(x) = \left\langle |\psi(x)|^2 - \frac{1}{2} \right\rangle_{\text{ensemble}}$$

B : GP mean field

\mathcal{H}_A : tWigner / +P blend

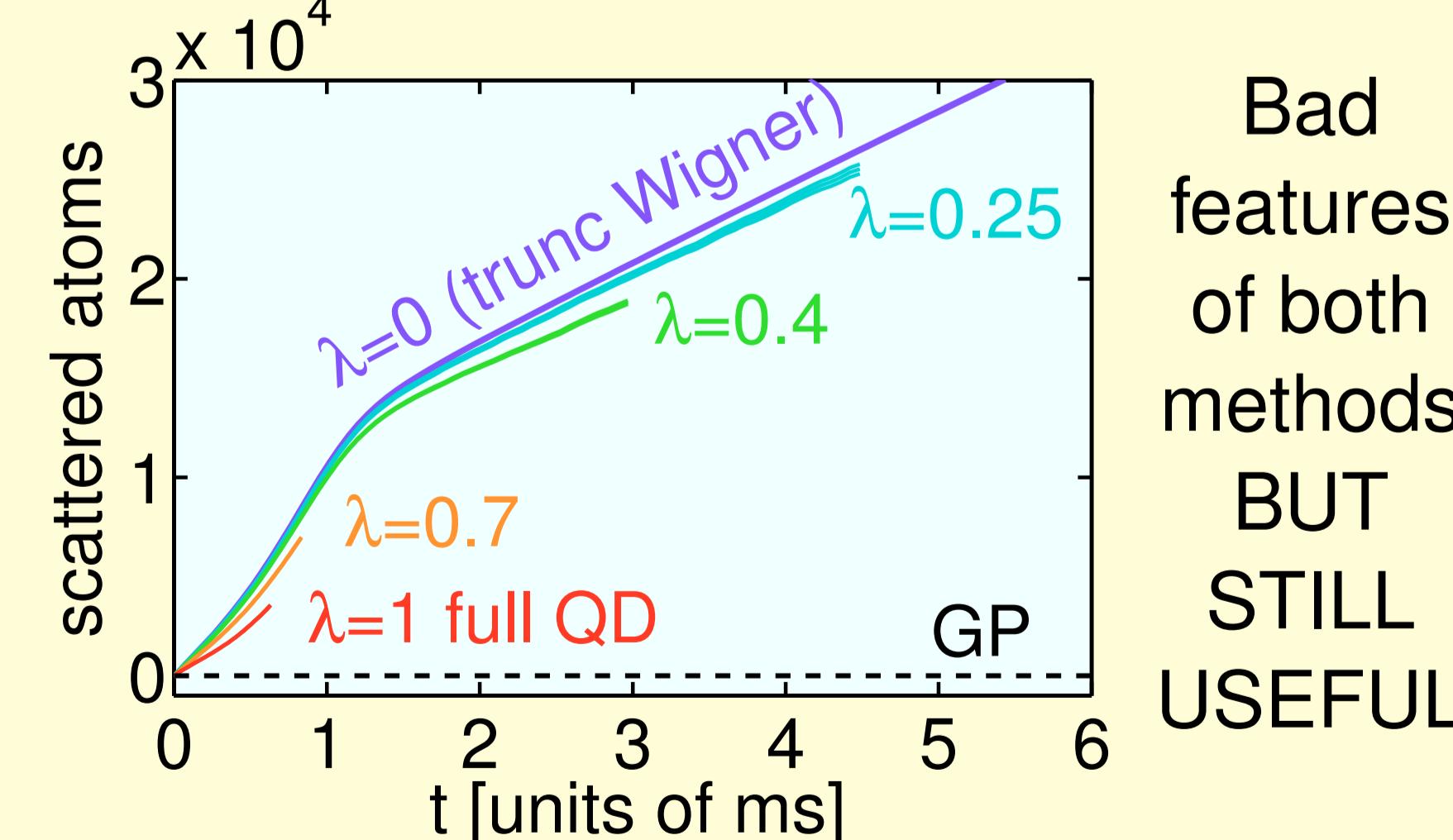
$\lambda = 0$: Wigner, $\lambda = 1$: complete QD

$$\psi(x, 0) = \tilde{\psi}(x, 0) = \phi_{GP}(x, 0) + \sqrt{\frac{1 - \lambda}{2}} \eta(x)$$

$$i\hbar \frac{d\psi(x)}{dt} = \left[-\frac{\hbar^2}{2m} \nabla^2 + g \psi(x) \tilde{\psi}^*(x) + \sqrt{\frac{\lambda g}{i\hbar}} \xi(x, t) \right] \psi(x)$$

$$i\hbar \frac{d\tilde{\psi}(x)}{dt} = \left[-\frac{\hbar^2}{2m} \nabla^2 + g \psi(x) \tilde{\psi}^*(x) + \sqrt{\frac{-\lambda g}{i\hbar}} \tilde{\xi}(x, t) \right] \tilde{\psi}(x)$$

$$\bar{n}(x) = \left\langle \tilde{\psi}^*(x) \psi(x) - \frac{1 - \lambda}{2} \right\rangle_{\text{ensemble}}$$

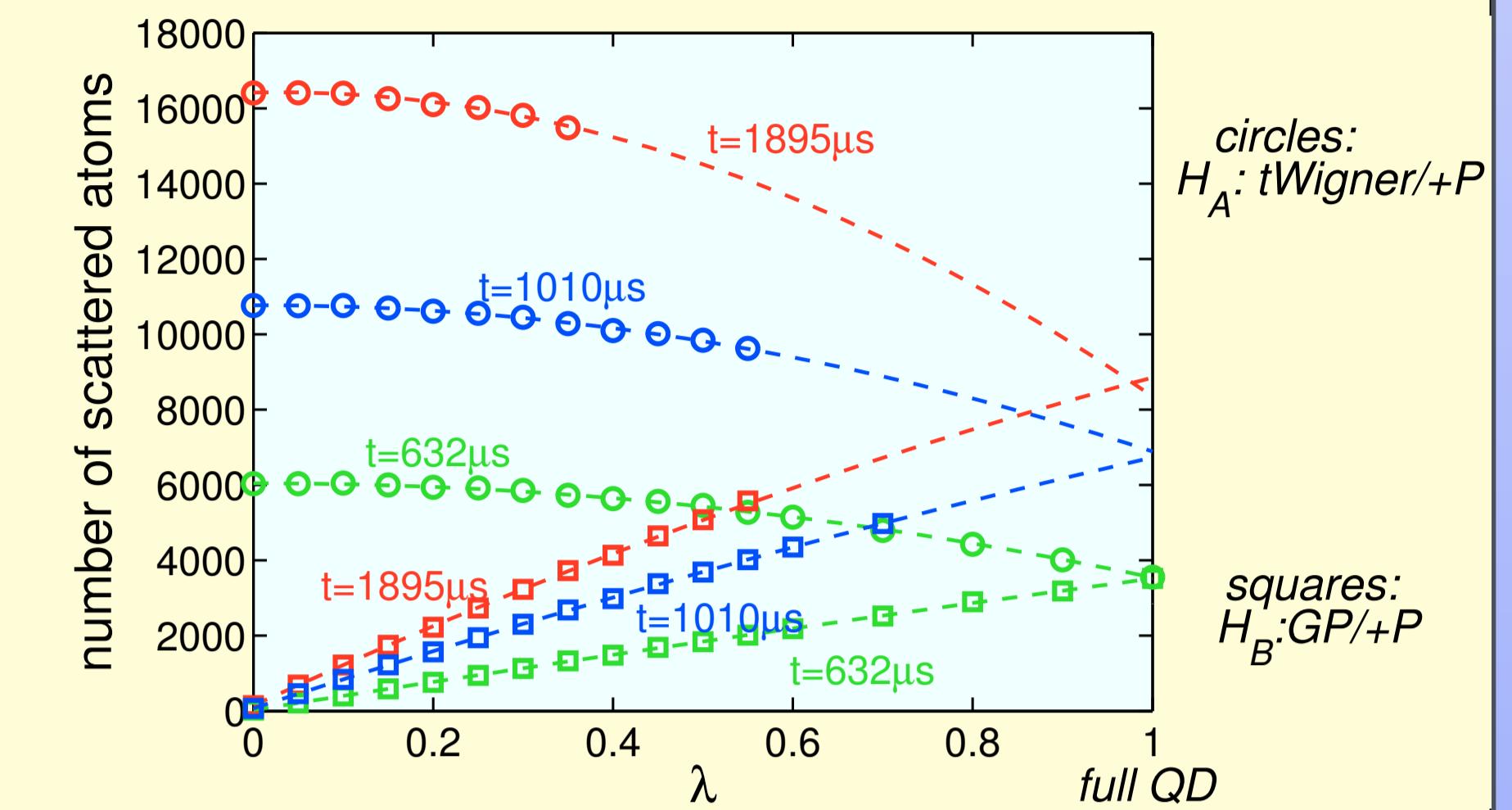


\mathcal{H}_B : GP / +P blend

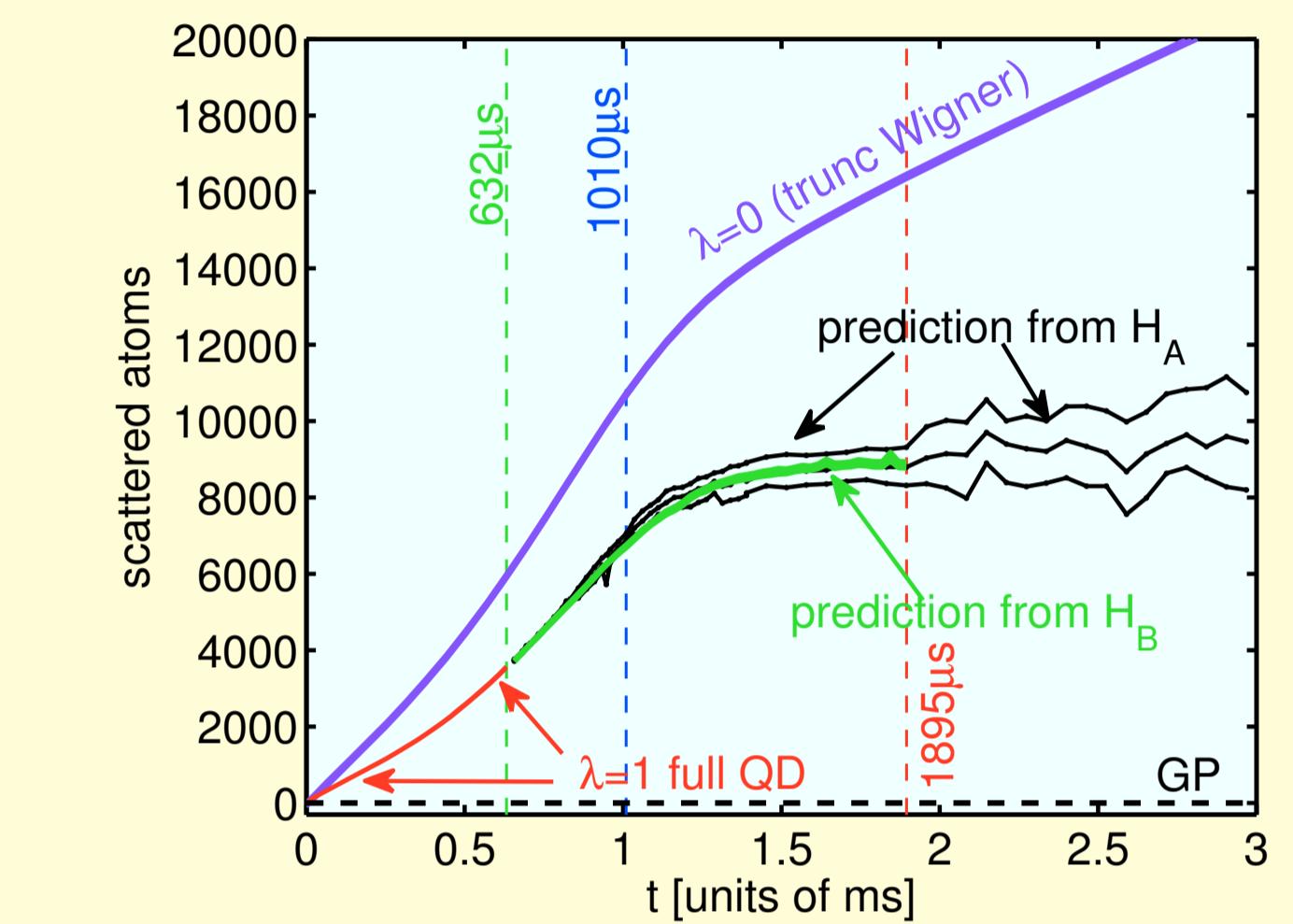
Identical to positive-P except noise is multiplied by $\sqrt{\lambda}$ as per $\sqrt{\frac{\lambda g}{i\hbar}} \xi(x, t)$, etc.

No. of halo scattered atoms

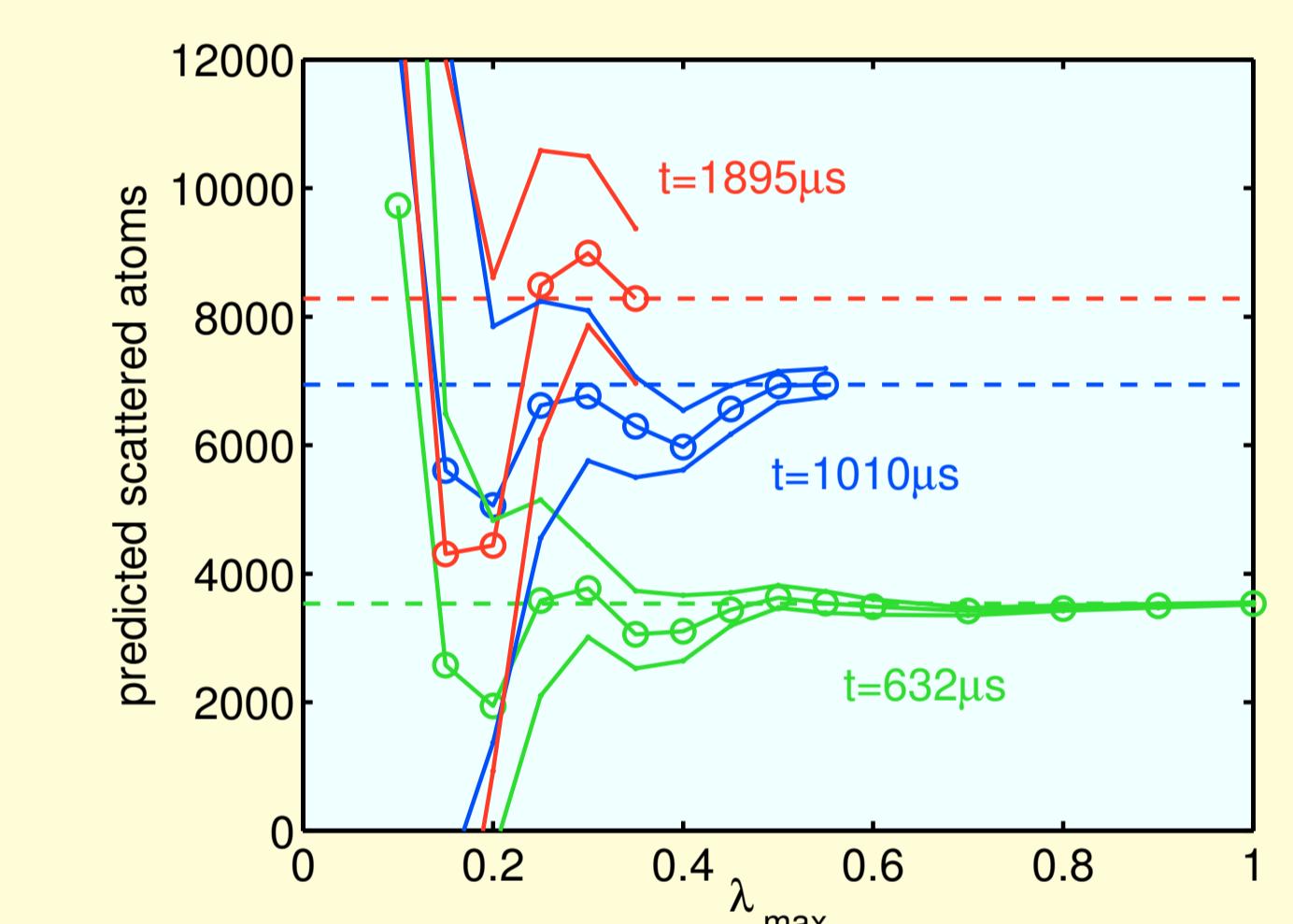
λ -EXTRAPOLATION:



FINAL PREDICTIONS: (1σ error bars)

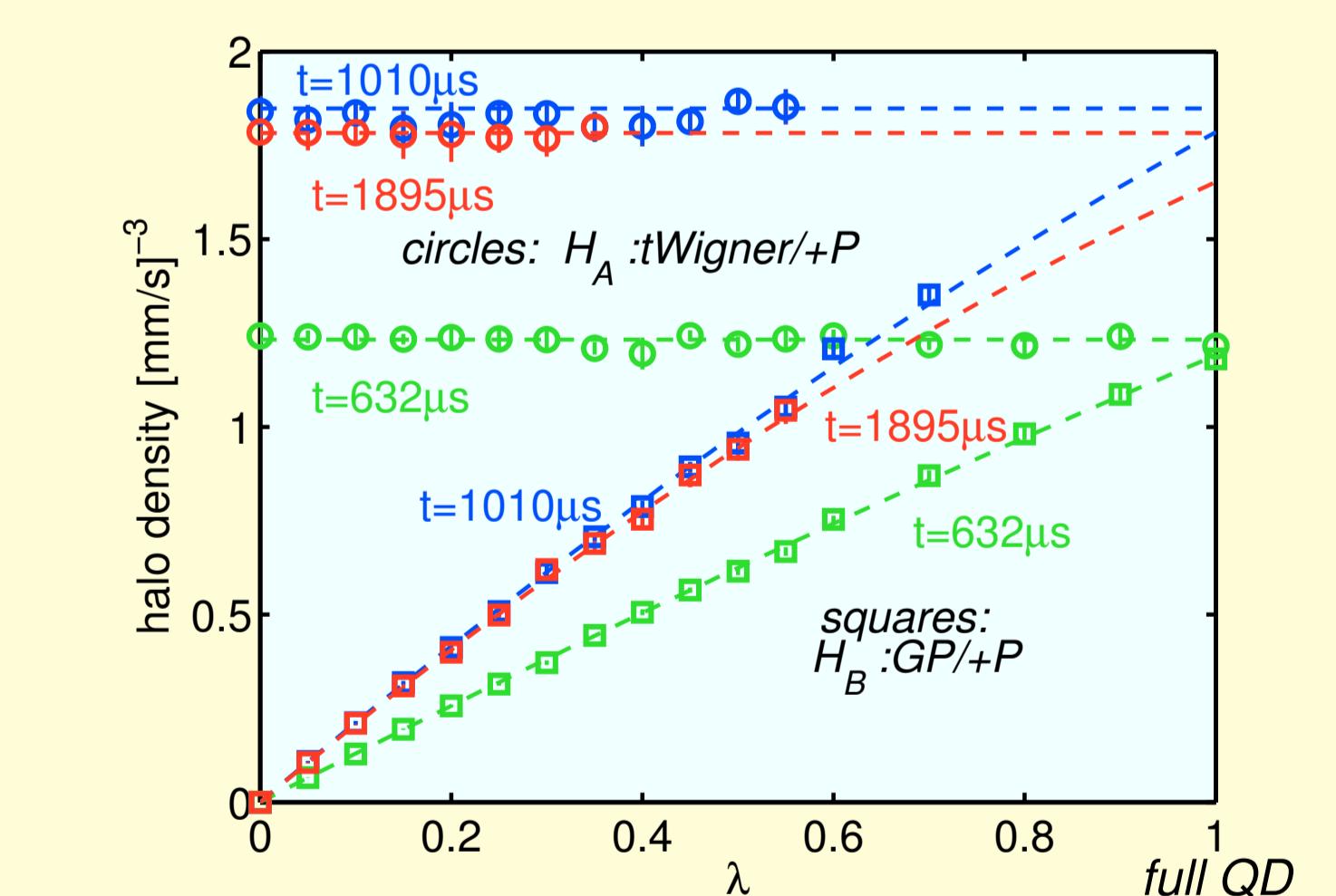


CHECK: EXTRAPOLATION USING A LIMITED λ RANGE $\lambda \in [0, \lambda_{\max}]$

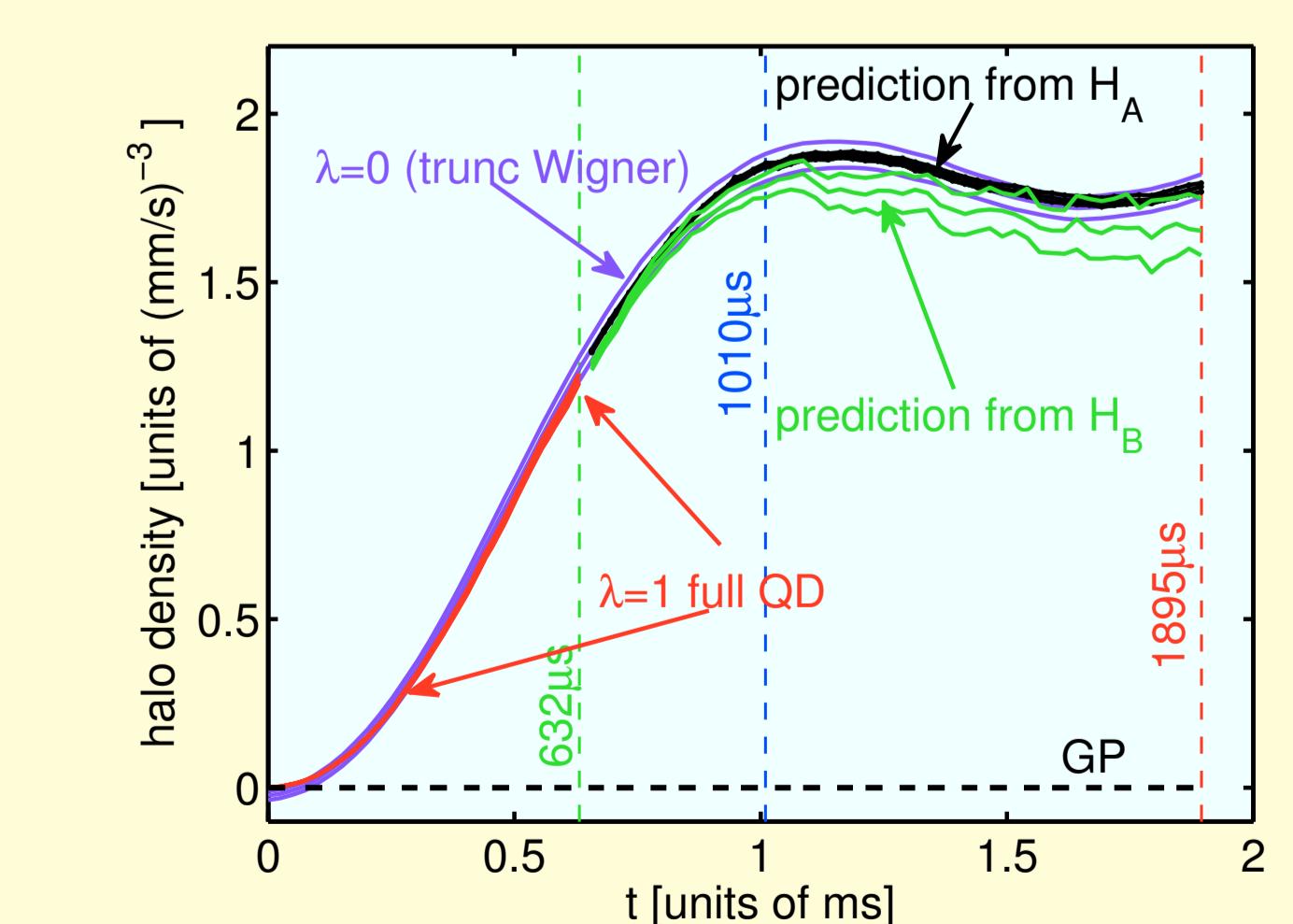


Peak halo density

λ -EXTRAPOLATION:



FINAL PREDICTIONS:



P. Deuar, arXiv:0903.1309, to appear in PRL.

This research was supported by the European Community under the contract MEIF-CT-2006-041390.