A controlled transition

FROM: Classical field simulations

TO: Full quantum dynamics

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Classical field and positive P simulations

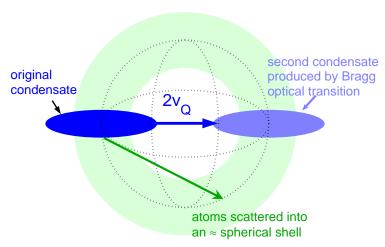
- Both give simulations of quantum dynamics beyond a linearised Hamiltonian
- Both are numerically tractable
- Each method has its drawbacks
- A hybrid method can alleviate some of these

Classical field ("truncated Wigner") simulations

EXAMPLE SYSTEM: BEC COLLISION

Mean field fails for scattered atoms

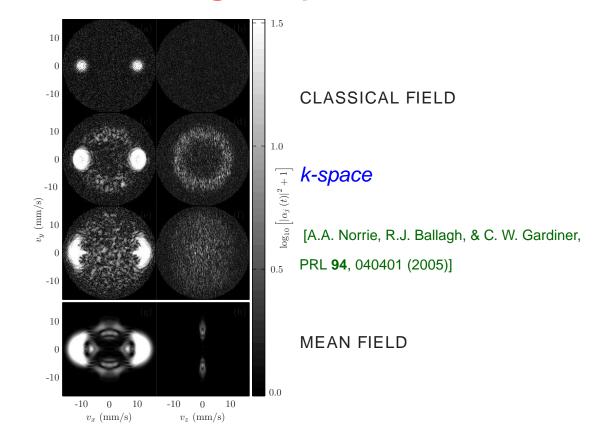
x-space



Experiments as per e.g.

[J. M. Vogels, K. Xu, & W. Ketterle, PRL **89**, 020401 (2002)] and many others

$$\widehat{H} = \text{K.E.} + \frac{g}{2} \int \left[\widehat{\Psi}^{\dagger}(x) \right]^2 \widehat{\Psi}(x)^2 d^3x$$

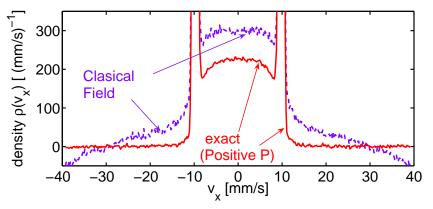


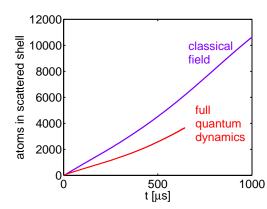
$$\psi(x,0) = \phi(x,0) + \frac{\eta(x)}{\sqrt{2}}$$
 (add $\frac{1}{2}$ virtual particle)

$$i\hbar \frac{d\psi(x)}{dt} = \left[-\frac{\hbar^2}{2m} \nabla^2 + g \left| \psi(x) \right|^2 \right] \psi(x)$$
 (mf GP equation)

Works well for large N= 6,000,000 particles

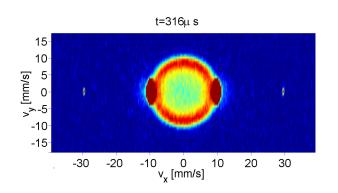
Less particles: N = 150,000





Full quantum dynamics (positive-P) calculations

$$\psi(x,0) = \phi(x,0)$$
 ; $\widetilde{\psi}(x,0) = \phi(x,0)$



[PD & P. D. Drummond, PRL 98, 120402 (2007)]

two fields no virtual initial particles

$$i\hbar \frac{d\psi(x)}{dt} = \left[-\frac{\hbar^2}{2m} \nabla^2 + g \psi(x) \widetilde{\psi}^*(x) + \sqrt{\frac{g}{i\hbar}} \, \xi(x,t) \right] \psi(x)$$

$$i\hbar \frac{d\widetilde{\psi}(x)}{dt} = \left[-\frac{\hbar^2}{2m} \nabla^2 + g \psi(x) \widetilde{\psi}^*(x) + \sqrt{\frac{-g}{i\hbar}} \, \widetilde{\xi}(x,t) \right] \widetilde{\psi}(x)$$

mean field GP + noise

Classical field has errors for N= 150,000 particles Positive P only works for short times

Hybrid representation

- Consruct a smooth transition parametrised by parameter λ
- $\lambda = 0$ is classical Field, $\lambda = 1$ is positive-P
- Want transition to be monotonic!

AFTER SOME TRIES, OBTAIN:

$$\psi(x,0) = \widetilde{\psi}(x,0) = \phi(x,0) + \sqrt{\frac{f_1(\lambda)}{2}\eta(x)}$$
 initial conditions

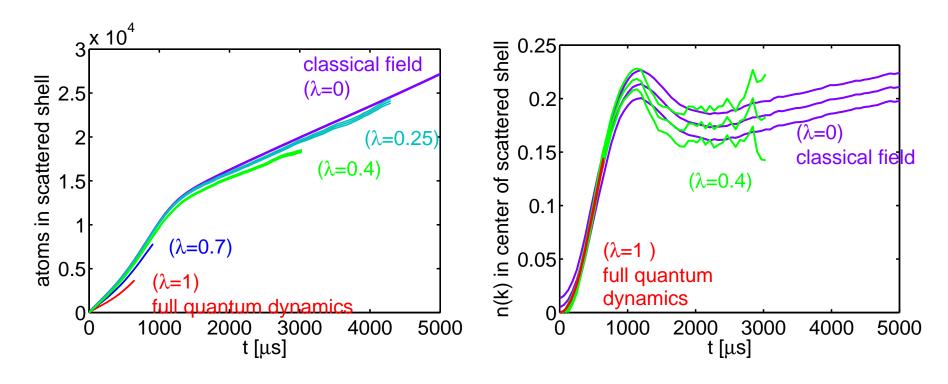
$$i\hbar \frac{d\psi(x)}{dt} = \left[-\frac{\hbar^2}{2m} \nabla^2 + g \psi(x) \widetilde{\psi}^*(x) + f_2(\lambda) \sqrt{\frac{g}{i\hbar}} \, \xi(x,t) \right] \psi(x)$$

$$i\hbar \frac{d\widetilde{\psi}(x)}{dt} = \left[-\frac{\hbar^2}{2m} \nabla^2 + g \psi(x) \widetilde{\psi}^*(x) + f_2(\lambda) \sqrt{\frac{-g}{i\hbar}} \, \widetilde{\xi}(x,t) \right] \widetilde{\psi}(x)$$

evolution

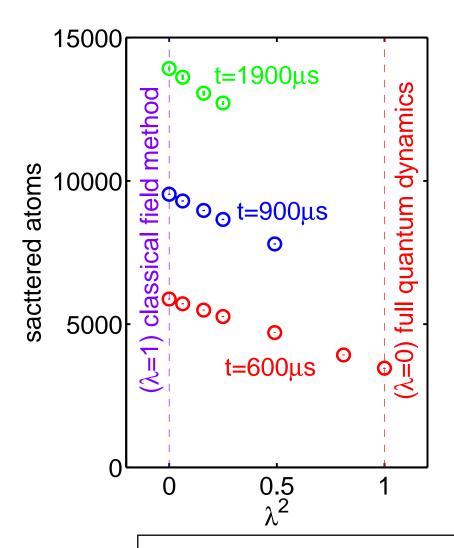
$$\overline{n}(x) = \left\langle \widehat{\Psi}^{\dagger}(x) \widehat{\Psi}(x) \right\rangle = \left\langle \widetilde{\psi}^{*}(x) \psi(x) - \frac{f_{1}(\lambda)}{2} \right\rangle_{\text{ensemble}} \quad \text{observables}$$

Resulting simulations



- Why combine bad features of both methods?
- Can get an idea of the size of the correction to the classical field method.
 e.g: total scattered number overestimated by several thousand
- Can verify some observables at long times
 e.g. density in main scattered shell is calculated correctly with classical field method

Estimating full quantum dynamics



- Calculate for various values of λ
- Here the observable varies linearly with λ^2
- Allows extrapolation to $\lambda = 1$ (full quantum dynamics) for much longer times

CONCLUSIONS:

- Quantitative assessment of classical field method accuracy
- Extrapolation to full QD for much longer times