

Semiconductor few-electron quantum dots as spin qubits for quantum computing

“Delft Spin Qubit Team”

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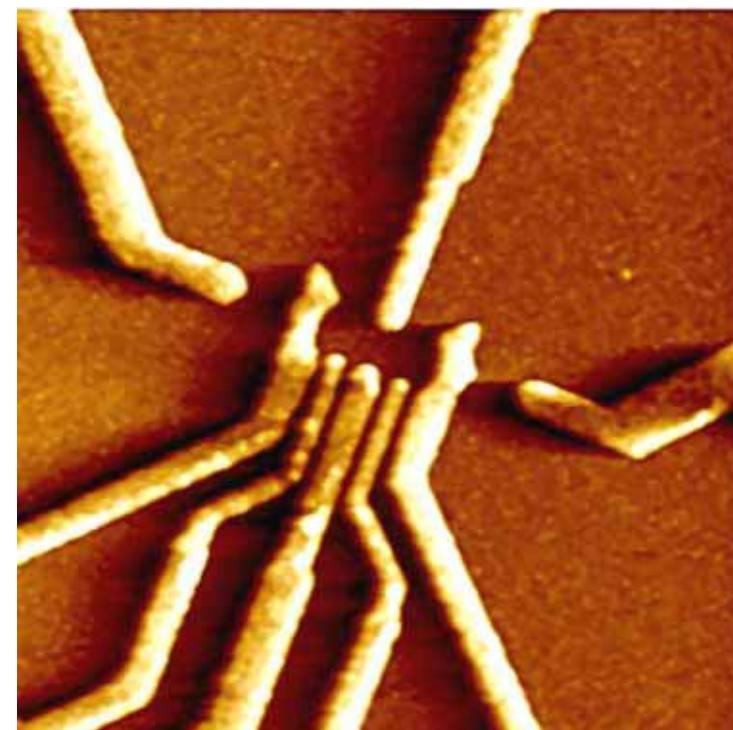
Leo Kouwenhoven

Collaborators

Tarucha group (Tokyo)

Loss group (Basel)

Fujisawa, Hirayama (NTT)

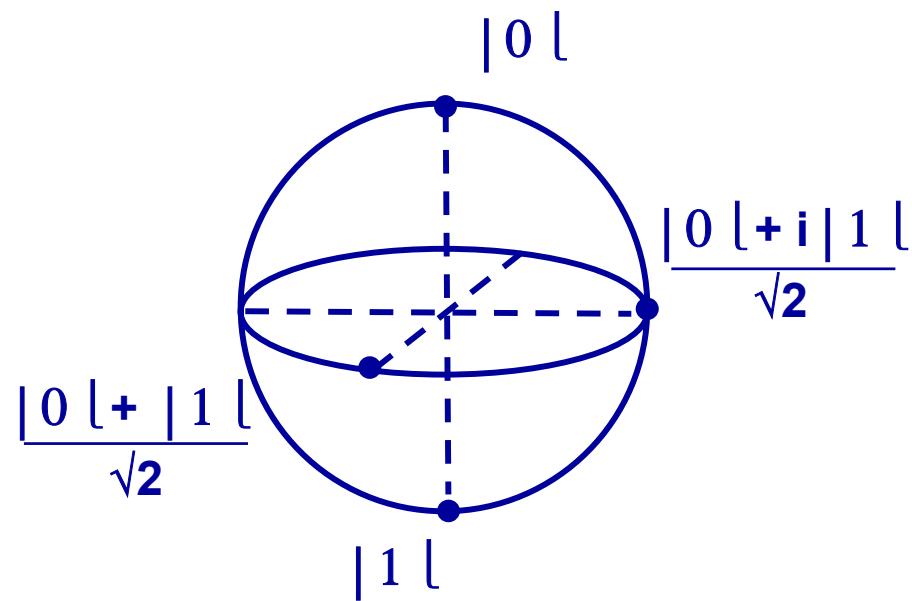


Outline

- I. Introduction to quantum computing**
- II. Quantum dots and spin**
- III. One-electron quantum dots as spin qubits**
- IV. Outlook**

Part I.

Introduction to quantum computing



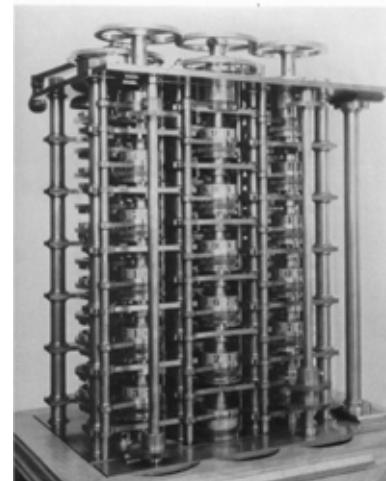
“Hard” problems

$$15 = 3 \times 5$$

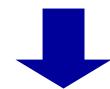
$$91 = \dots \times \dots ?$$

$$437 = \dots \times \dots ?$$

???



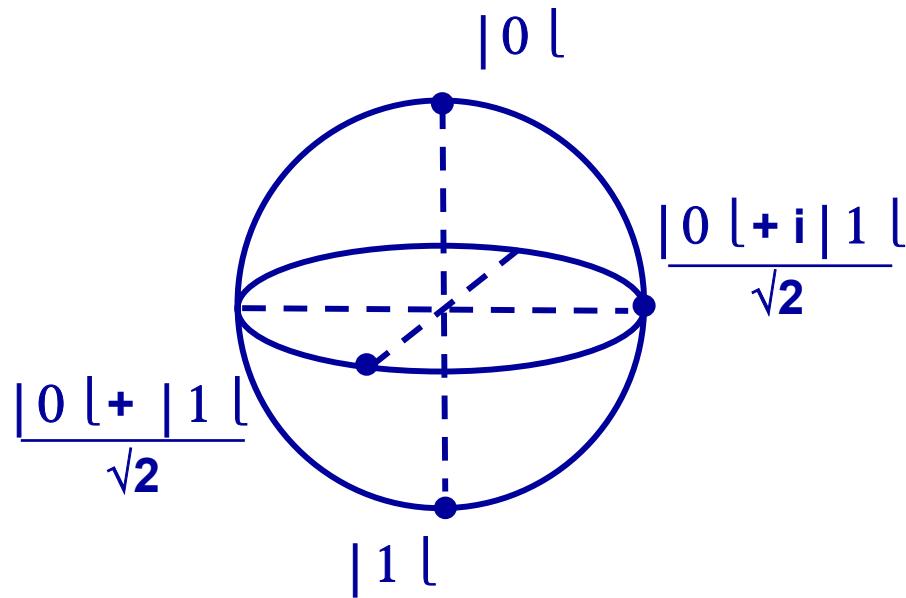
Courtesy IBM Corporation



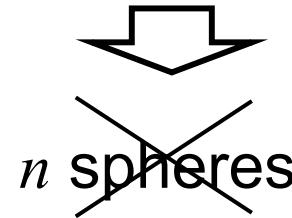
factoring takes exponential effort

“Hard” problems are beyond the reach of any machine relying on the classical laws of physics

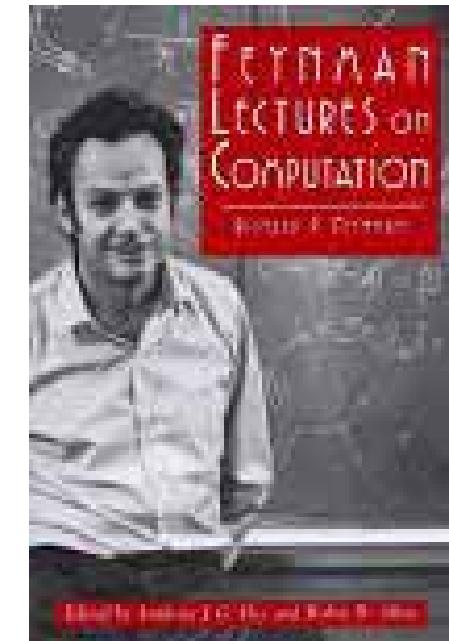
Complexity of Quantum Systems



n coupled quantum bits



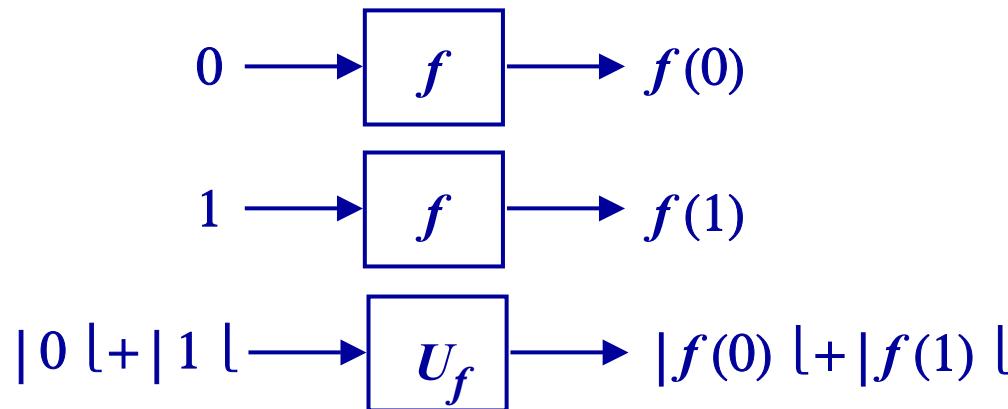
2^n degrees of freedom!



Could a quantum computer *efficiently* simulate quantum systems ?

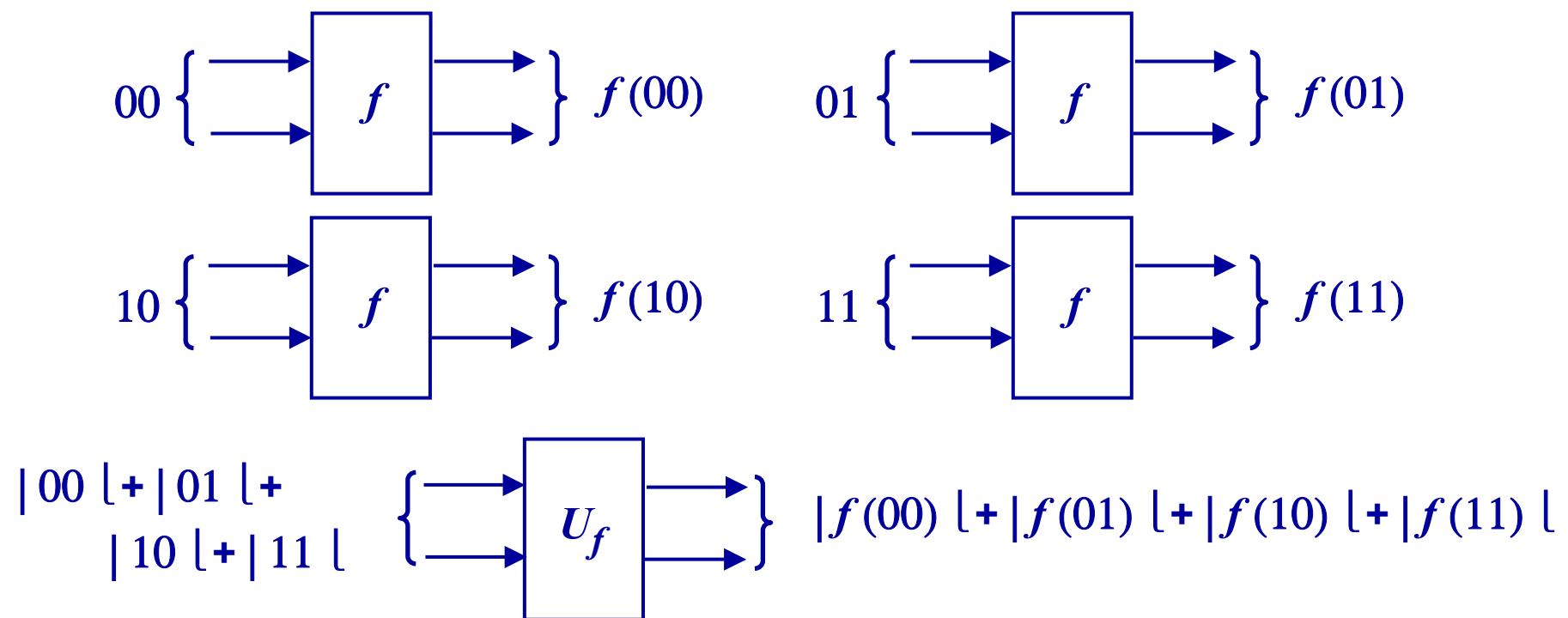
Could it be used to solve hard problems ?

Quantum Parallelism



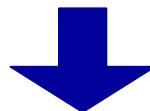
Computational power:
classically $\propto n$
quantum $\propto 2^n$

D. Deutsch, 1985



Quantum algorithms

Measurement of $|f(0)\rangle + |f(1)\rangle$ gives either $f(0)$ or $f(1)$



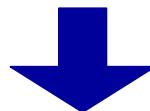
The exponential power appears inaccessible ...

Nevertheless: quantum algorithms make computational speed-ups possible !

- Exponential for factoring integers (P. Shor 1994)
- Quadratic for unstructured searches (L. Grover 1996)
- Exponential for quantum simulations (S. Lloyd 1996)

Quantum error correction

Decoherence destroys quantum parallelism



The exponential power appears limited in time ...

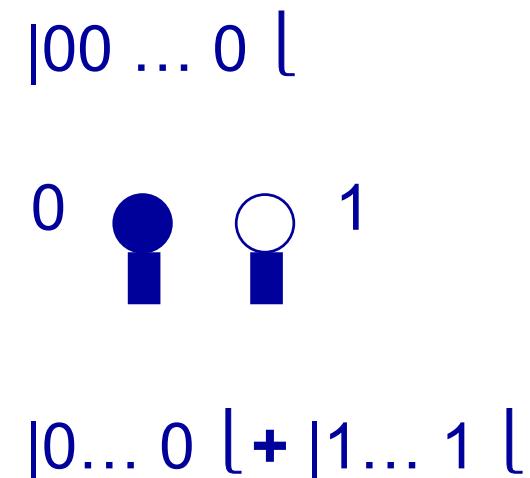
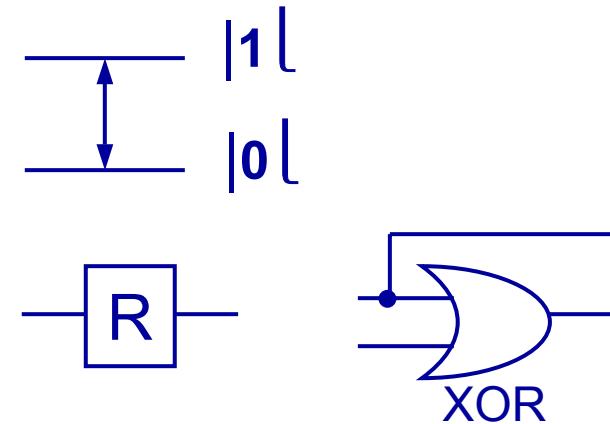
Nevertheless: quantum error correction makes arbitrarily long quantum computations possible !

- Quantum error correction (P. Shor 1996, A. Steane 1996)
- Accuracy threshold (D. Aharonov 1997, A. Kitaev 1997, ...): $\sim 10^{-4}$

Implementation of quantum computers

D. DiVincenzo

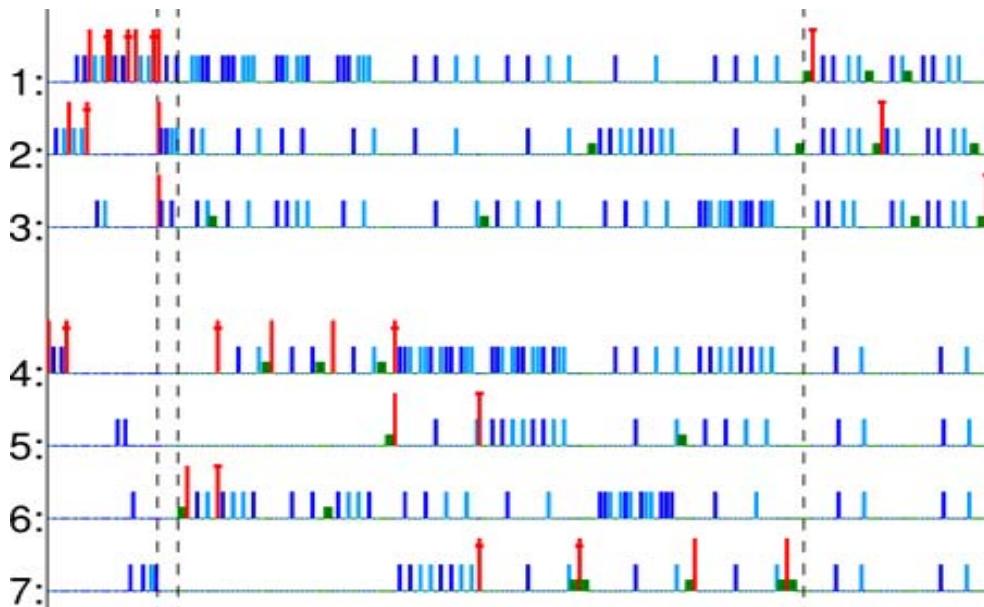
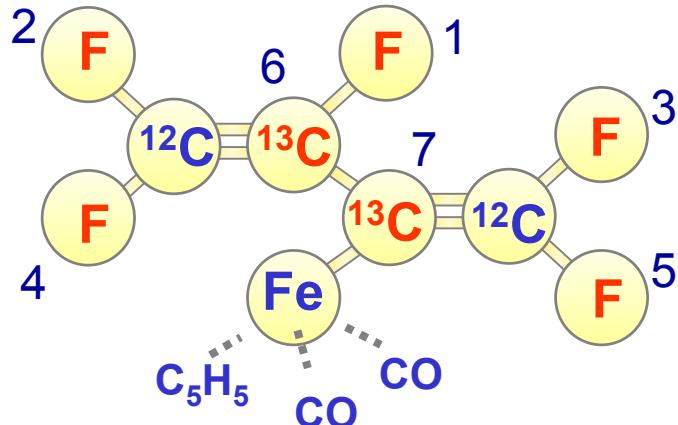
1. Well-defined qubits
2. Universal set of quantum gates
3. Initialization to a pure state
4. Qubit-specific measurement
5. Long coherence times



Key challenge: combine access to qubits (initialization, control, readout) with high degree of isolation (coherence)

Factoring 15 with nuclear spins

Vandersypen et al., *Nature* 414, 883 (2001)



$$15 = 3 \times 5$$

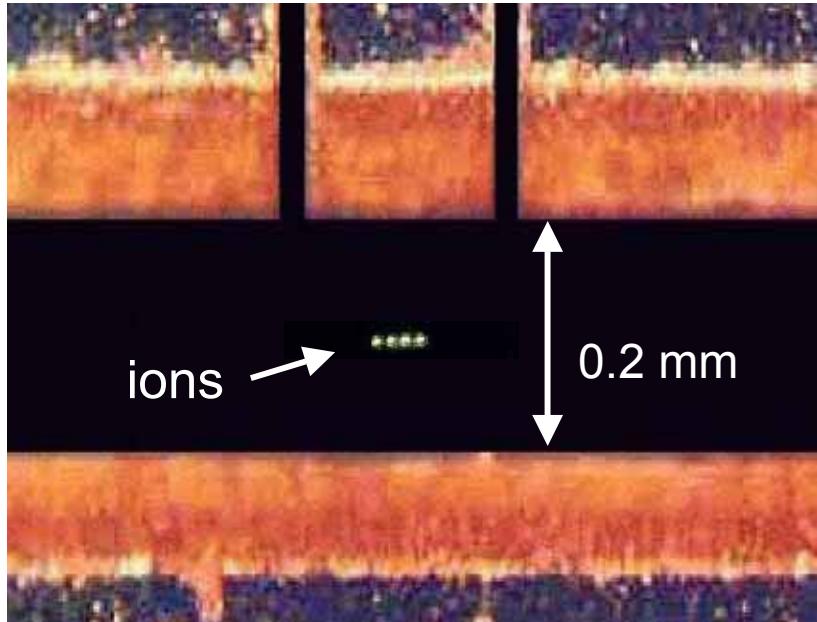
Proof-of-principle of quantum computing

But:

No practical path for scaling liquid NMR to many more qubits

Other techniques

Trapped ions



Courtesy D. Wineland, NIST

Scalability is a big problem!

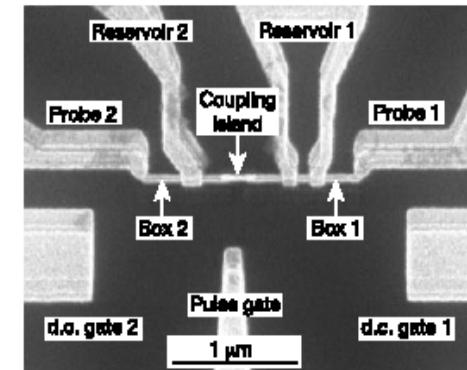
Neutral atoms in optical traps



Courtesy J. Kimble, Caltech

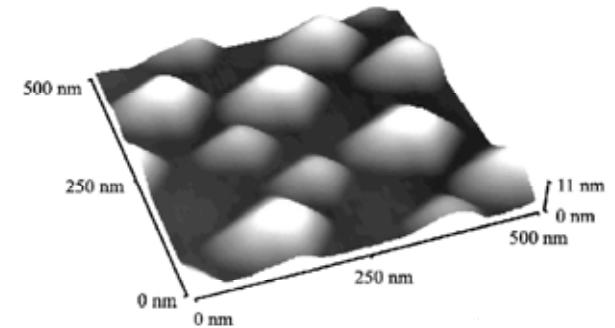
Charge quantum bits in the solid state

- *Superconducting charge qubits*
 - One- and two-qubit operations
 - Single-shot read-out

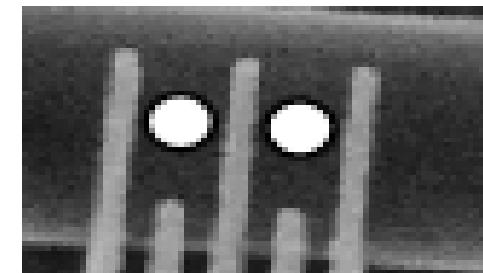


Courtesy Y. Nakamura, NEC

- *Semiconductor quantum dots*
 - Excitons in self-assembled quantum dots
 - Double dot charge qubit

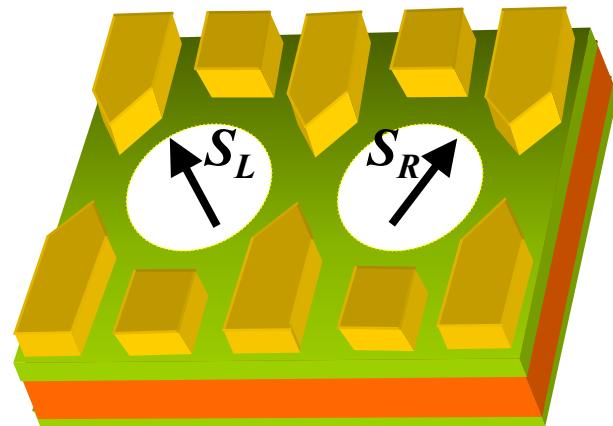


- Charge easy to manipulate and read out
- But coherence time only ~ ns



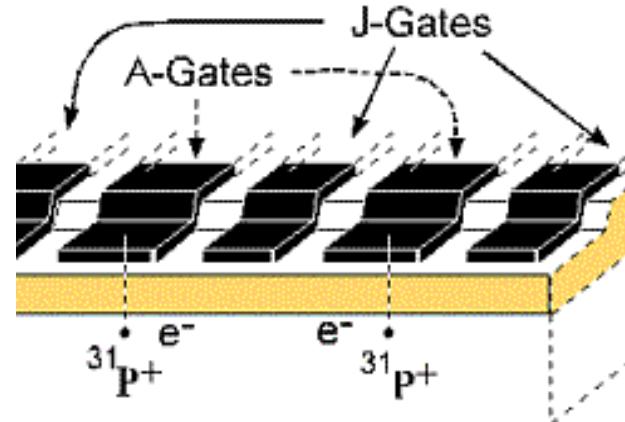
“Scalable” single-spin qubit proposals

Electron spin in gated quantum dots (electrical)



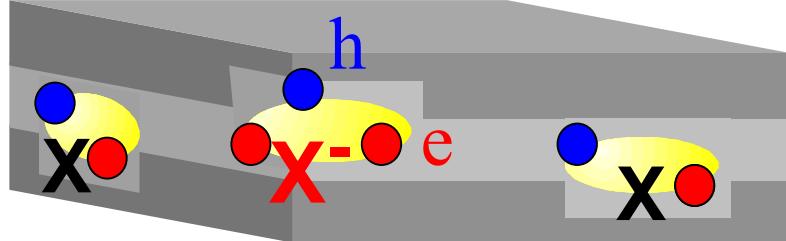
Loss & DiVincenzo, PRA 1998

Nuclear spin of P donors in Si



Kane, Nature 1998

Electron spin in self-assembled quantum dots (optical)



Gammon et al.
Imamoglu et al., PRL 1999

Electron spin in a lateral quantum dot

Why electron spin in a lateral quantum dot?

- Electron spin is a natural 2-level system
- Well isolated from environment ♥
long coherence times expected
- Lateral quantum dots are very flexible and controllable systems, should be scalable

Of course, not only application-driven:

- Spin physics at the fundamental quantum limit: 1 spin !

Part II:

Quantum dots and spin

2D Artificial Atoms



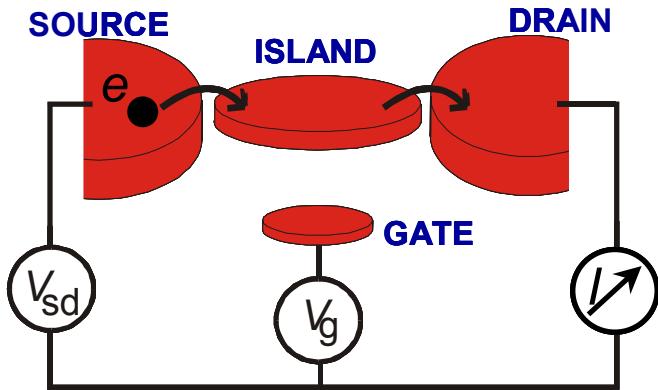
2D Artificial Atoms

1 Ta				2 Ha
3 Et	4 Au			5 Ko
7 Sa	8 To	9 Ho		10 Mi
13	14	15	16 Wi	17 Fr
			18 El	19
				20 Da

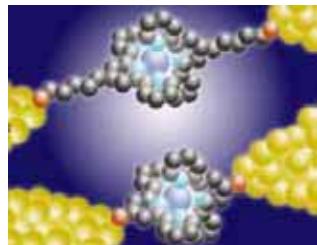
Quantum dots

- Small box occupied by electrons (holes)
- Coupled via tunnel barriers to source and drain reservoirs
- Coupled capacitively to gate electrode(s)
- Box (island) has discrete energy spectrum

♥ **artificial atom!**

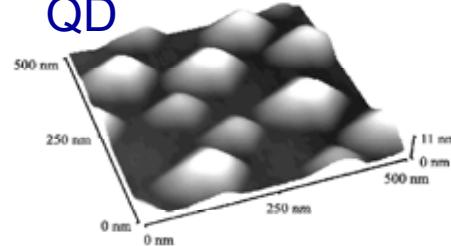


single molecule



1 nm

self-assembled QD



10 nm

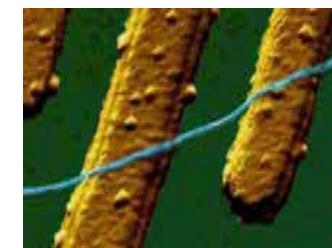
control & flexibility

lateral QD



100 nm

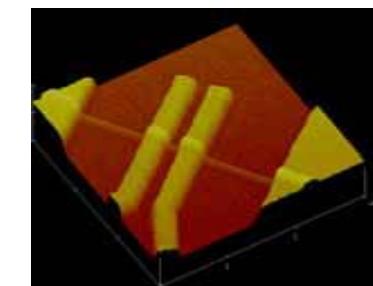
nanotube



metallic (superconducting) nanoparticle

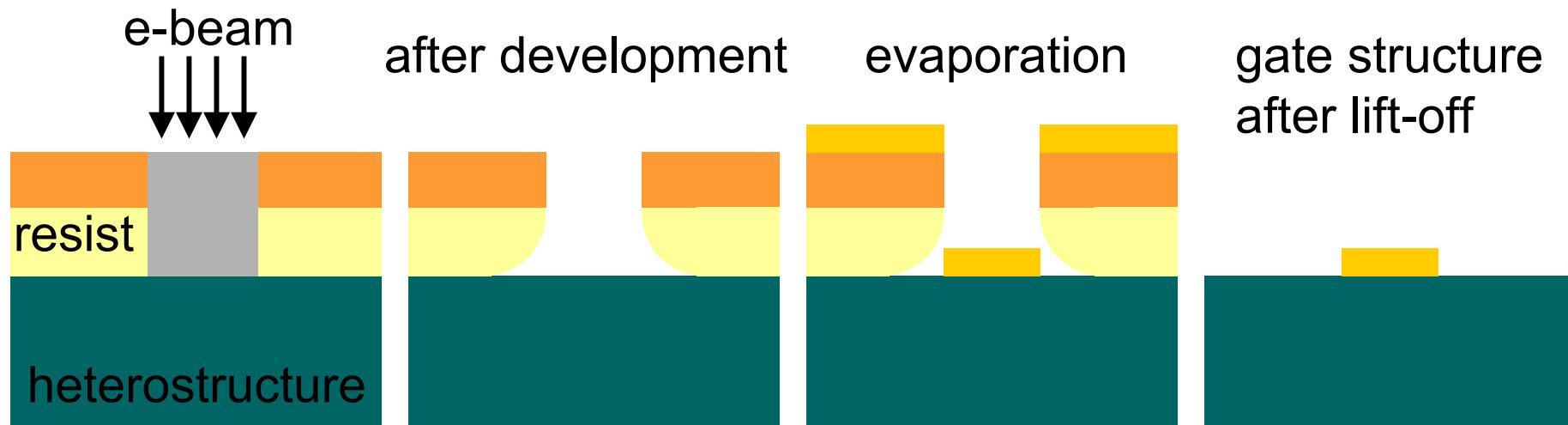
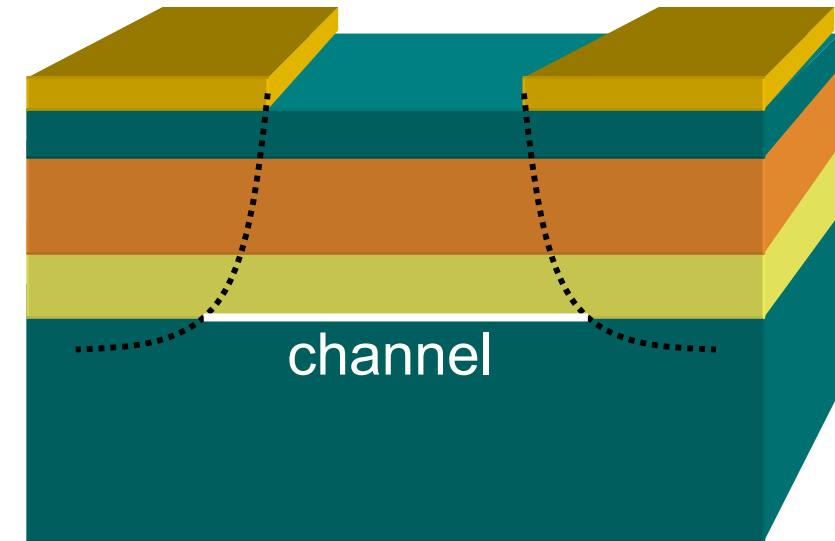


vertical QD

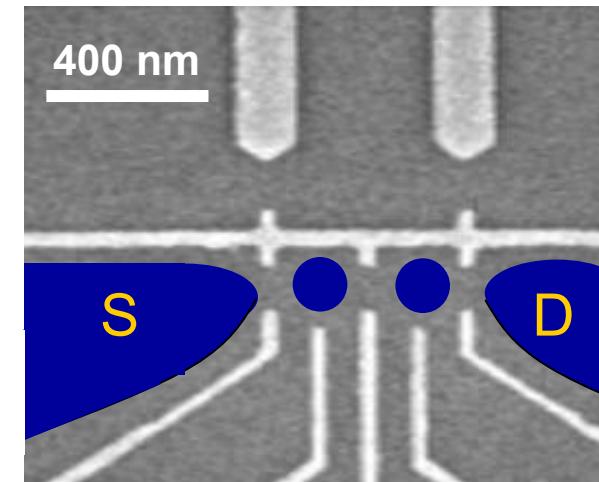
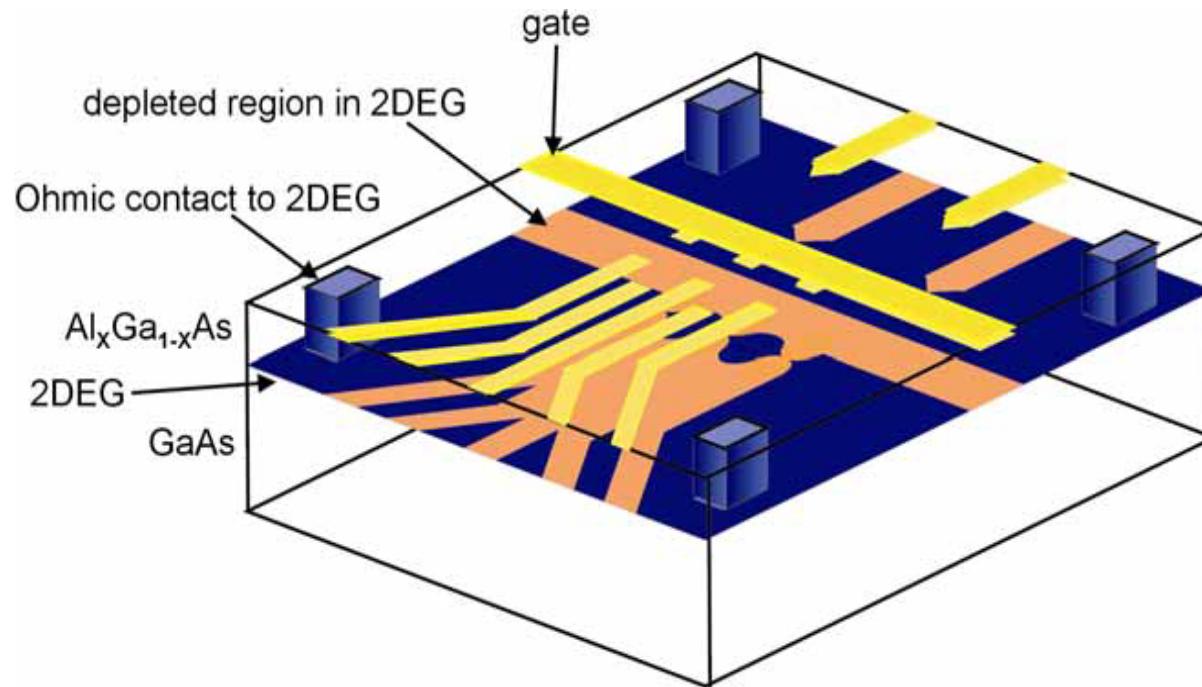


nanowire

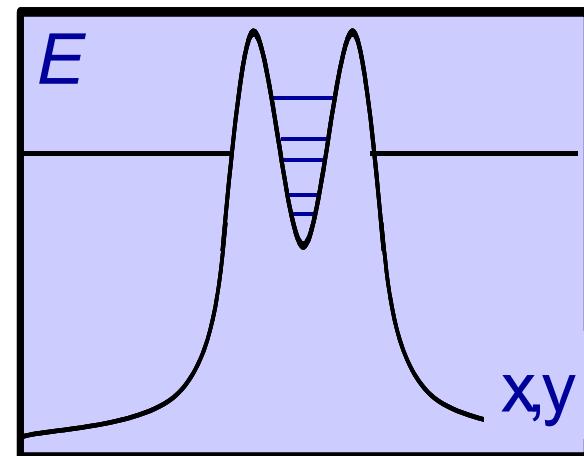
Heterostructure processing



Lateral QD fabrication

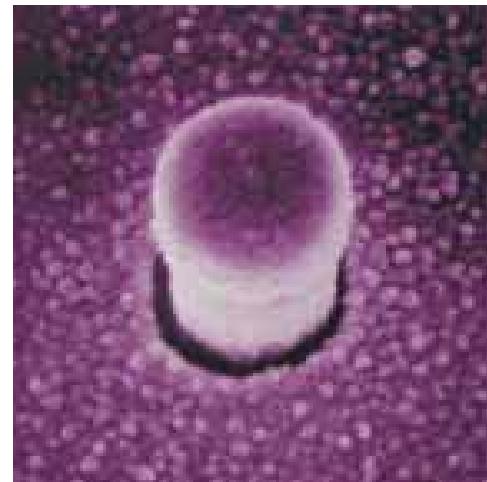


- High-mobility 2DEG ($\sim 10^6 \text{ cm}^2/\text{Vs}$)
- Density $\sim 10^{15} \text{ m}^{-2}$ $\heartsuit \lambda_F \sim 30 \text{ nm}$
- Resolution gate structure $\sim 20 \text{ nm}$
- Dot size $\sim 100 \text{ nm}$
- Comparable to electron wavelength
 \heartsuit discrete energy spectrum

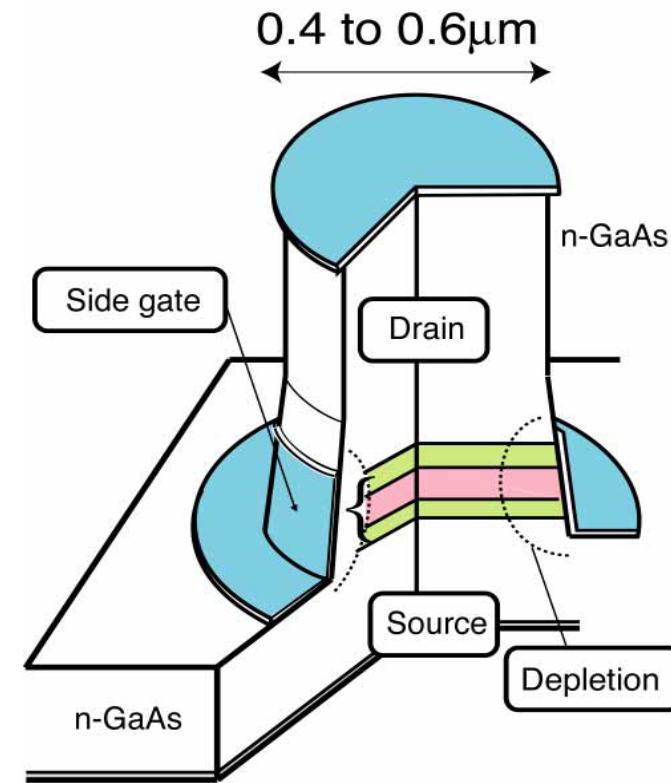
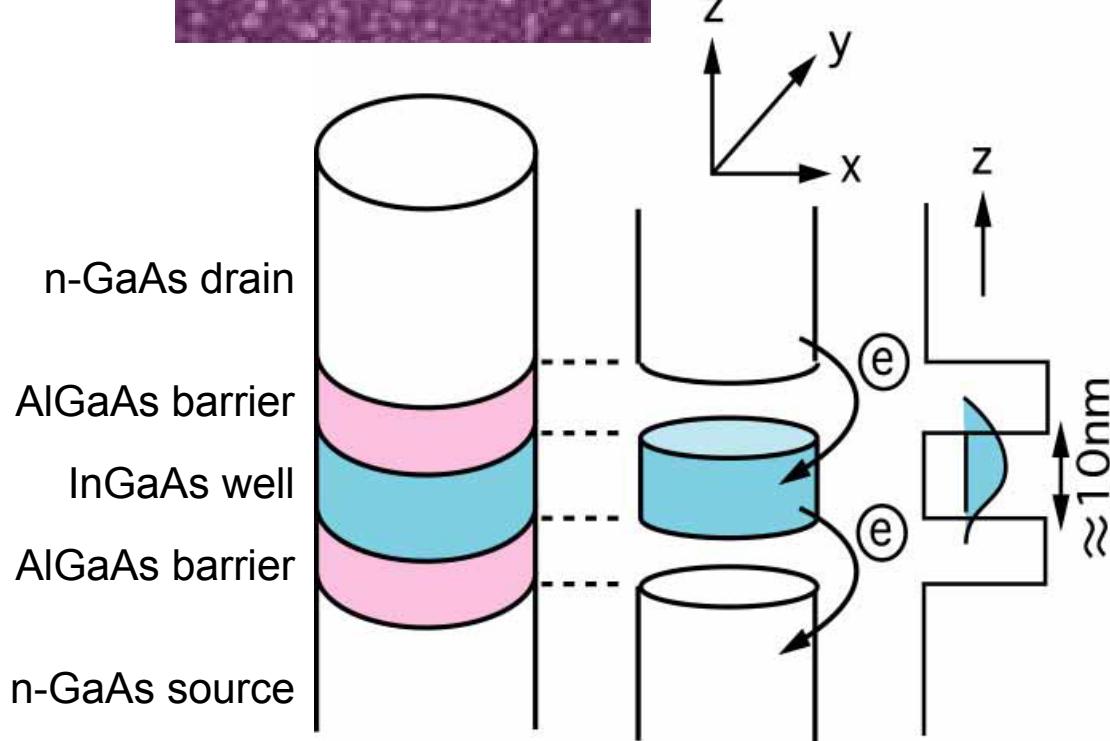


Vertical QD fabrication

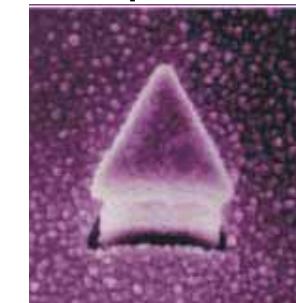
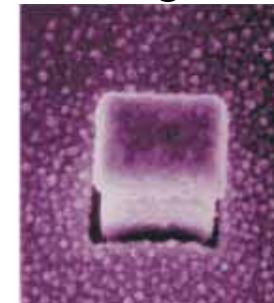
Tarucha PRL 77, 3613 ('96)



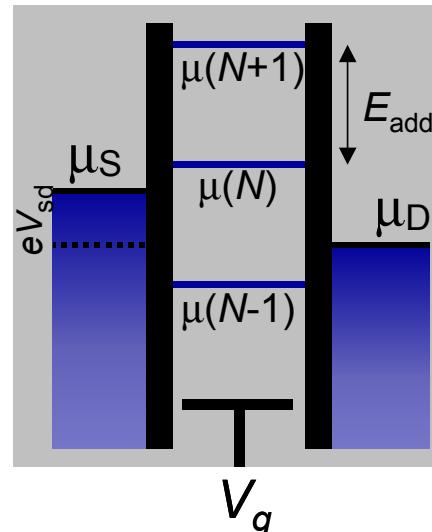
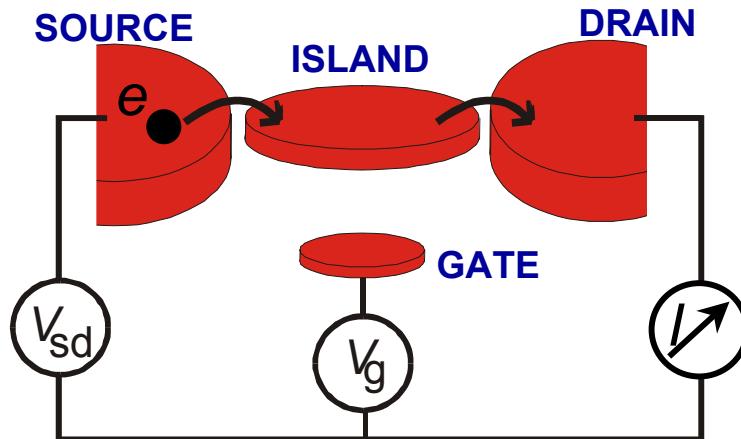
Circular pillar
made from double
barrier structure
by dry and wet
etching



Other geometries possible



Coulomb Blockade in transport



charging energy:

$$E_C = e^2/C_\Sigma$$

$$C \propto R,$$

R (radius) 100 nm

♥ $E_C \sim \text{few meV}$

CONSTANT INTERACTION MODEL

$$U(N) = \frac{(-Ne + C_g V_g + q_0)^2}{2C_\Sigma} + E_{int}(N)$$

$$C_\Sigma = C_L + C_R + C_g$$

$$\mu(N) \equiv U(N) - U(N-1) = E_C [(N - \frac{1}{2} - q_0/e) - C_g V_g / e] + E_N \quad E_C = e^2/C_\Sigma$$

$$E_{add} \equiv \mu(N+1) - \mu(N) = E_C + \Delta E$$

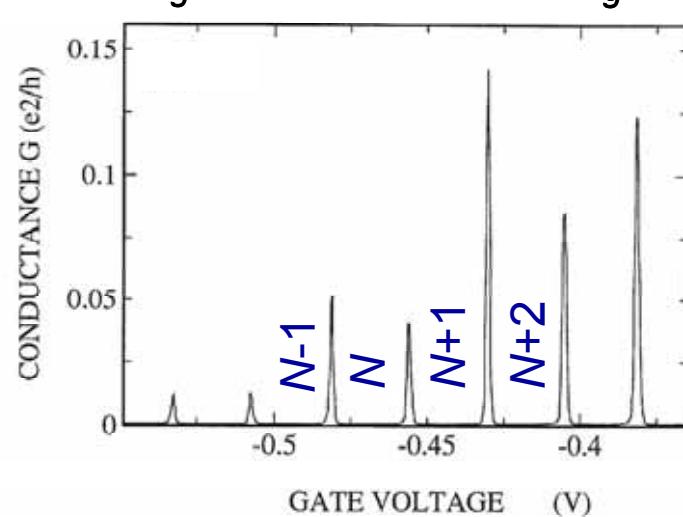
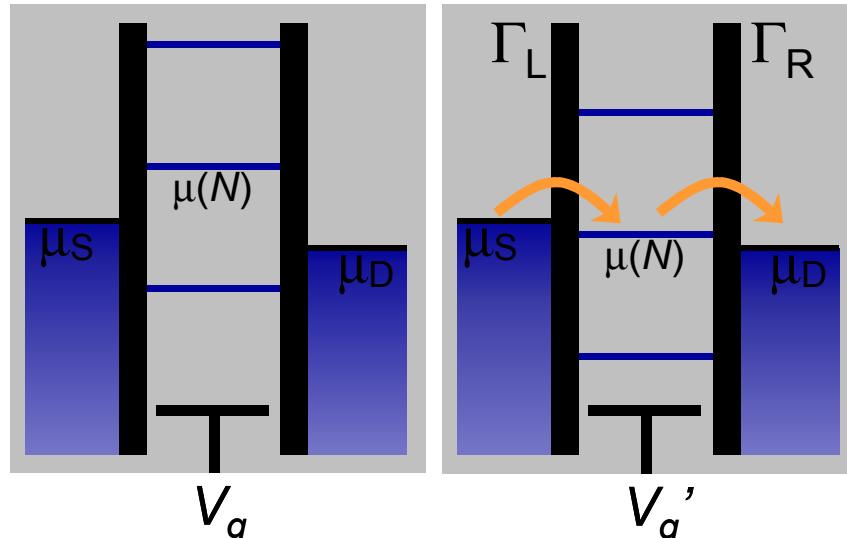
- Transport when $\mu(N) = \mu_S, \mu_D$
- Small C_Σ ♥ large addition energy
- At low T , small V_{sd} , energy not available



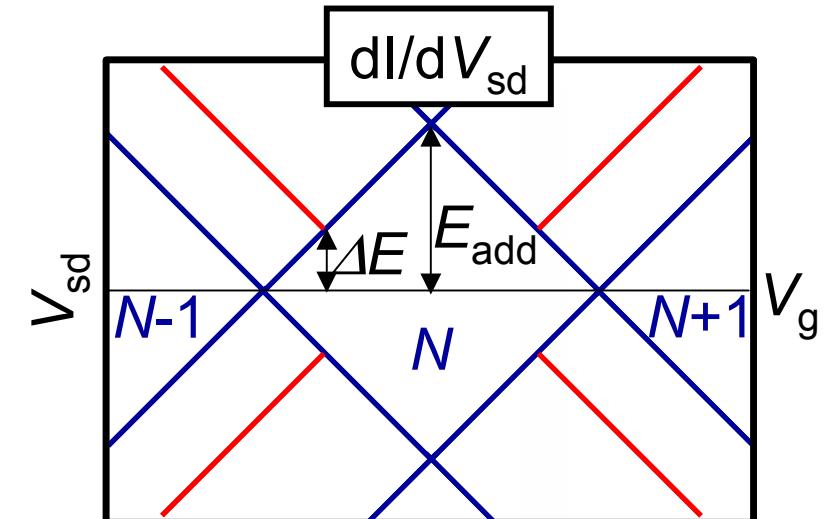
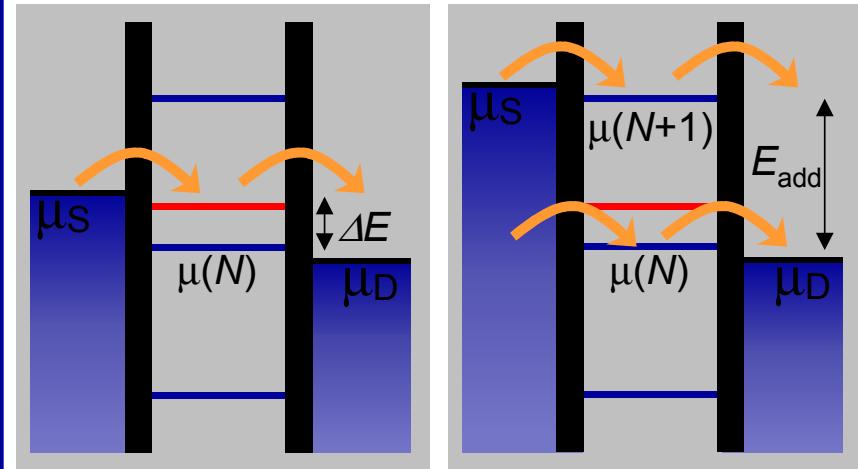
**CURRENT = 0
(COULOMB BLOCKADE)**

Single electron tunneling

Linear response

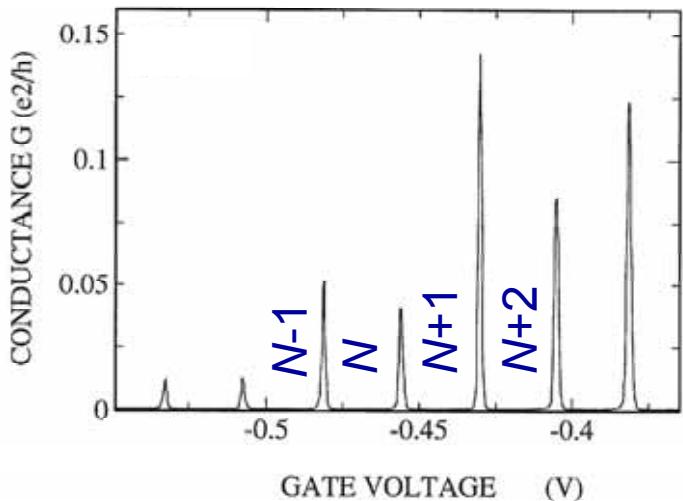


Non-linear transport



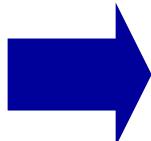
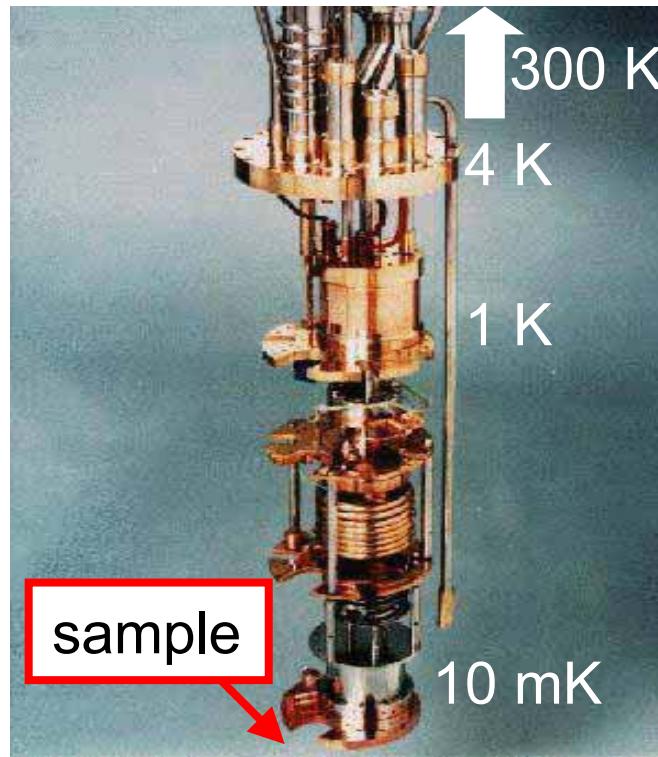
Large-bias spectroscopy!

Experimental set-up

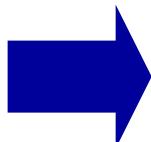


- Resonance width: lifetime ($\hbar\Gamma$), V_{sd} , temperature, noise
- Peak spacing: charging + single-particle spacing
- $E_C \sim$ few meV
- $\Delta E \sim$ few 0.1 meV
- $3 \text{ fA} / \text{Hz}^{1/2} \heartsuit 1 \text{ e}/50 \mu\text{s}$

Dilution refrigerator:
10 mK (\sim few μeV)

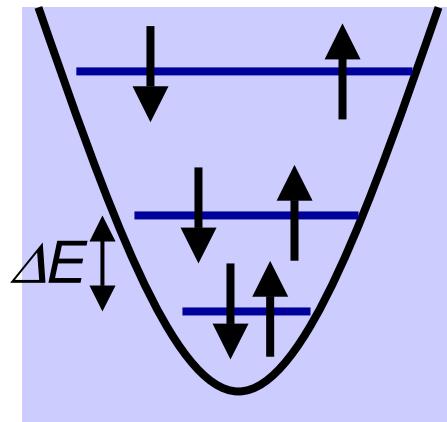


sample

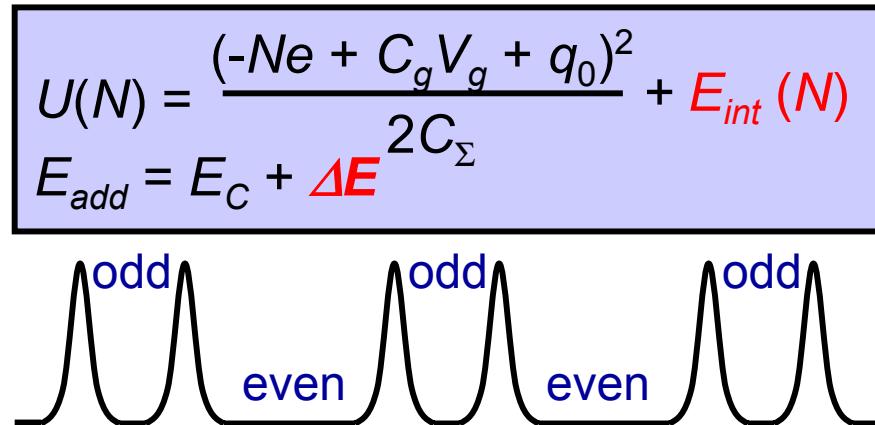


low-noise
amplifiers
+
filtering

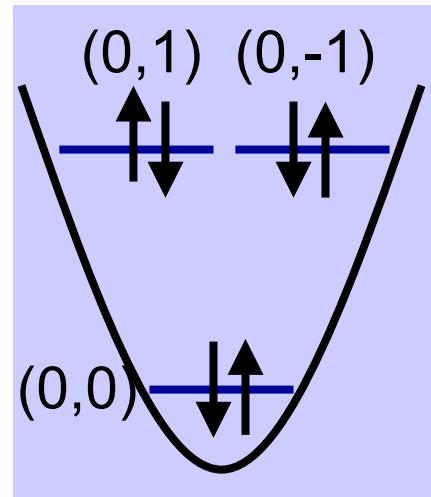
Single-particle levels & shell structure



orbitals 2-fold
degenerate ♥
expect
even/odd
effects in
peakspacing

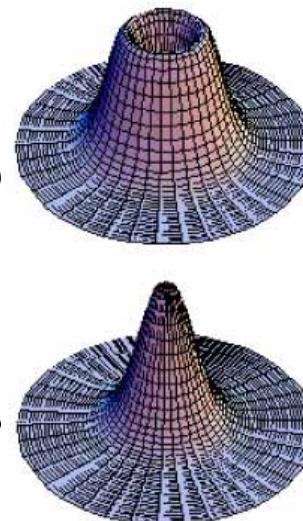


$$B = 0 \text{ T}: \\ E_{nl} = (2n+|l|+1)\hbar\omega_0$$



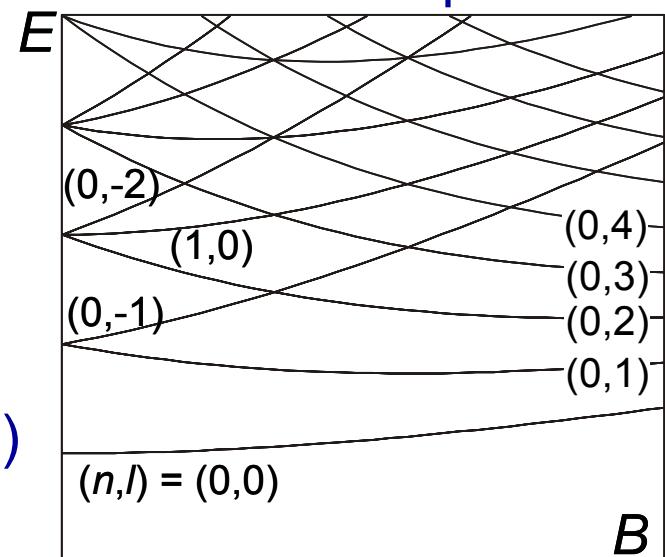
2p

1s



$N=2, 6, 12$
full shell
(noble
elements)

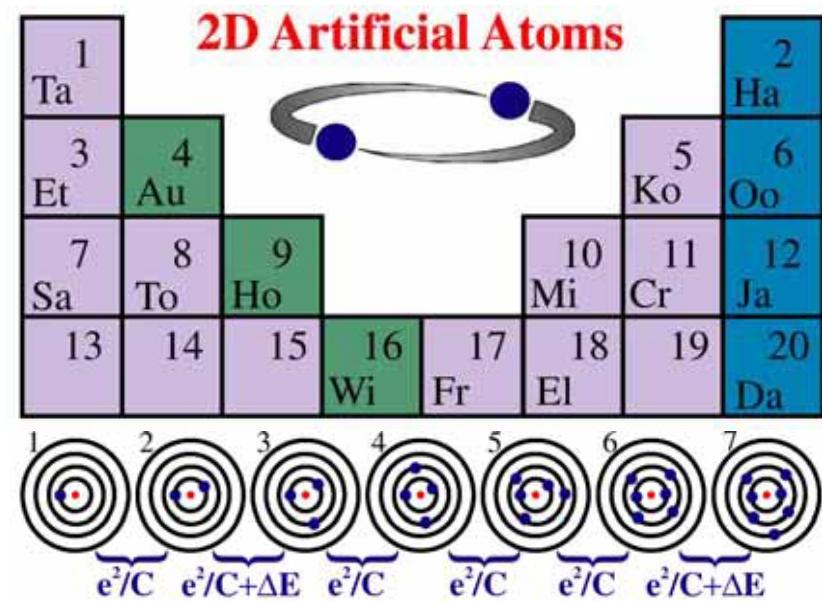
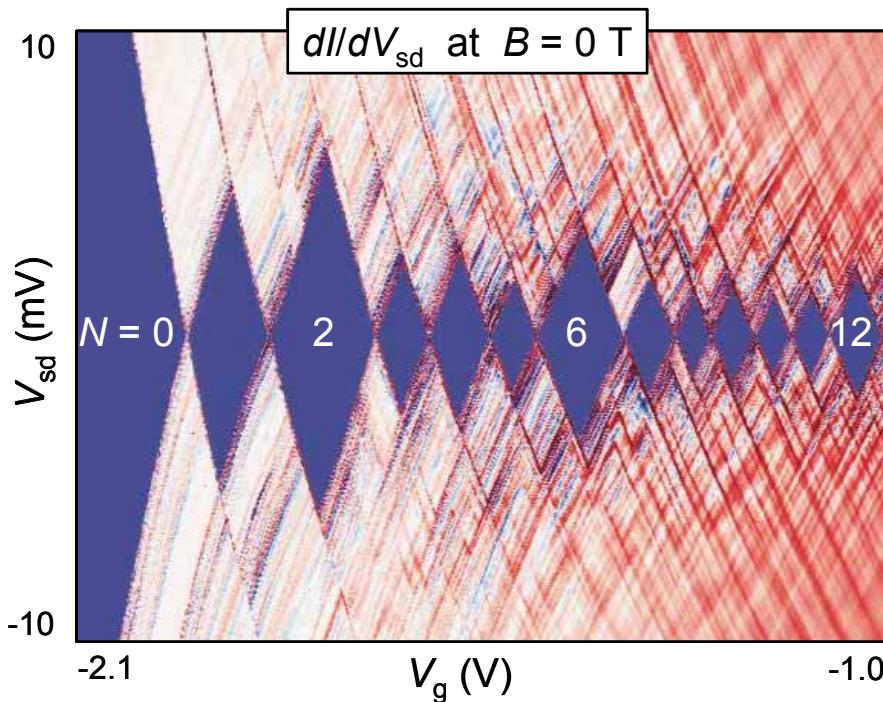
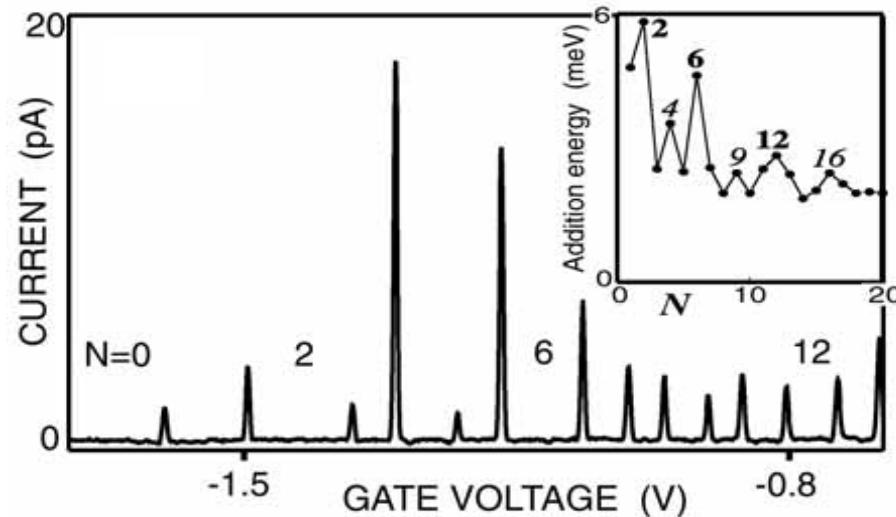
2D harmonic potential:
Darwin-Fock spectrum



B

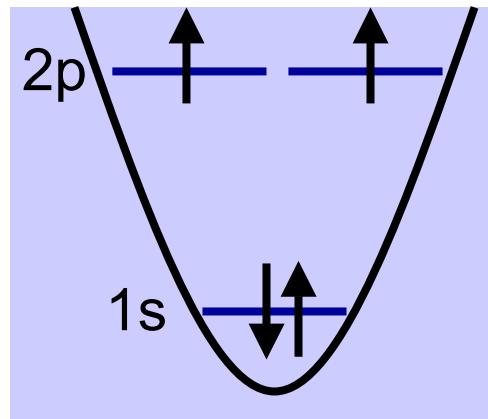
Artificial atoms

Kouwenhoven et al., Science 278, 1788 ('97)

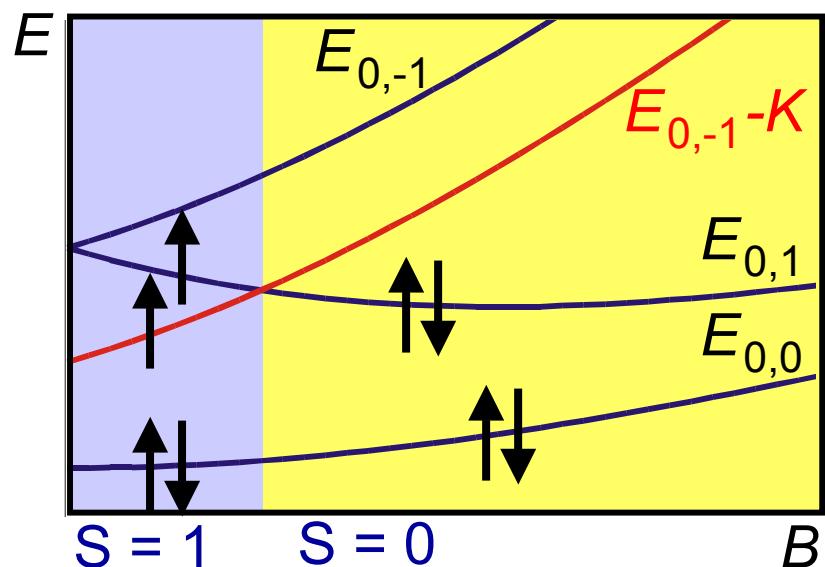
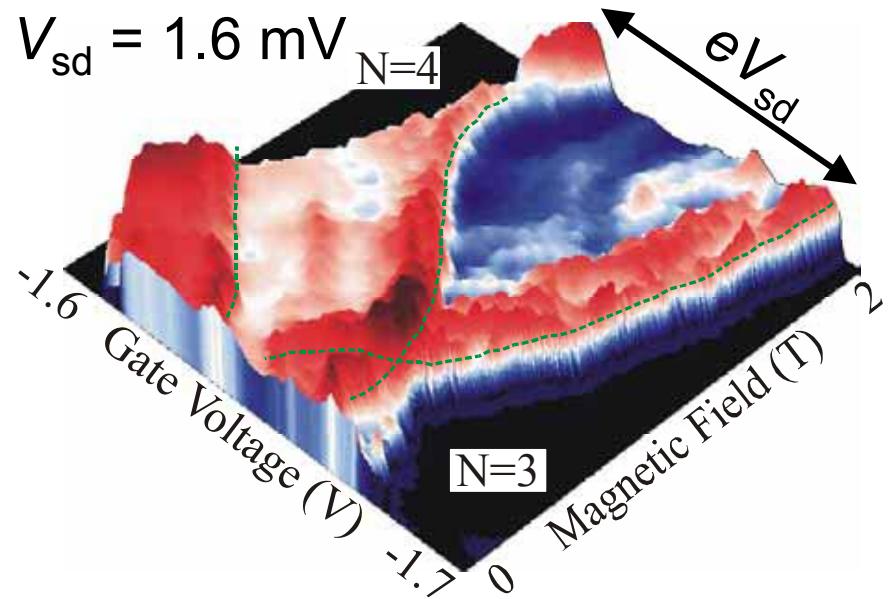


- Few-electron QDs are *artificial atoms*
- Change element by tuning V_g
- Full shell $N = 2, 6, 12, 20\dots$ (noble elements)
- Half-full shell $N = 4, 9, 16\dots$

Hund's rule for $N = 4$

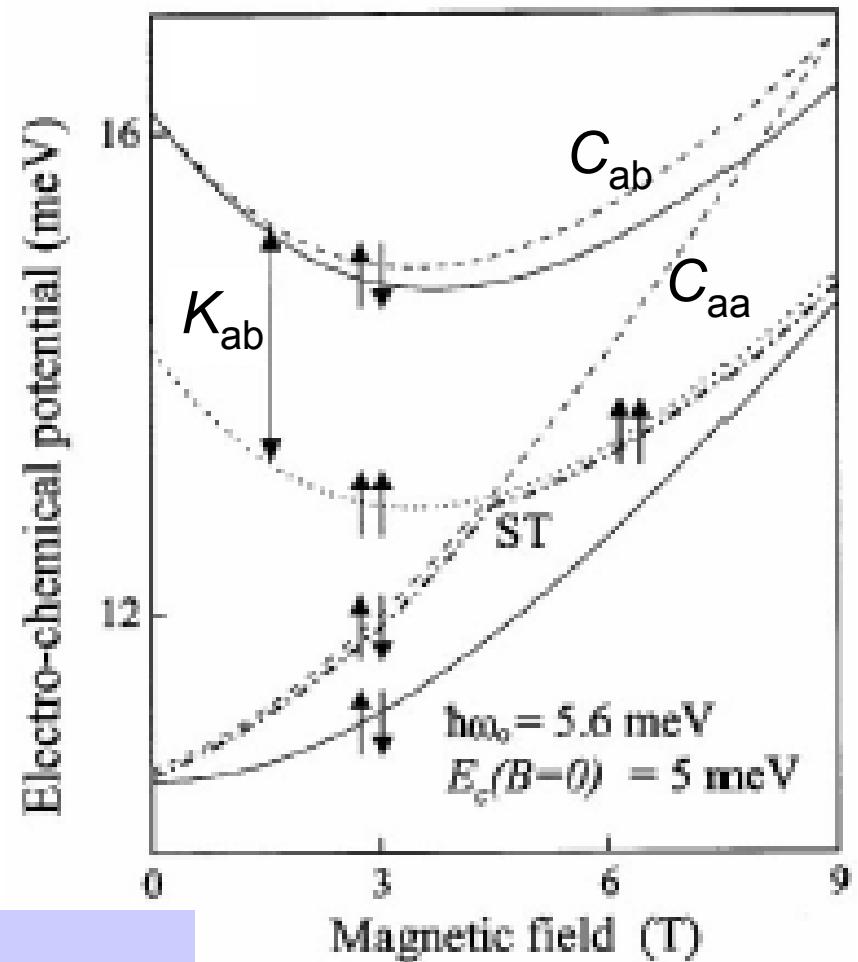
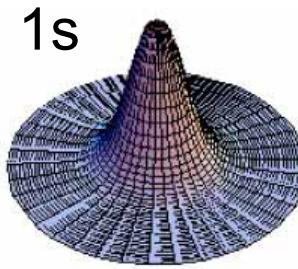
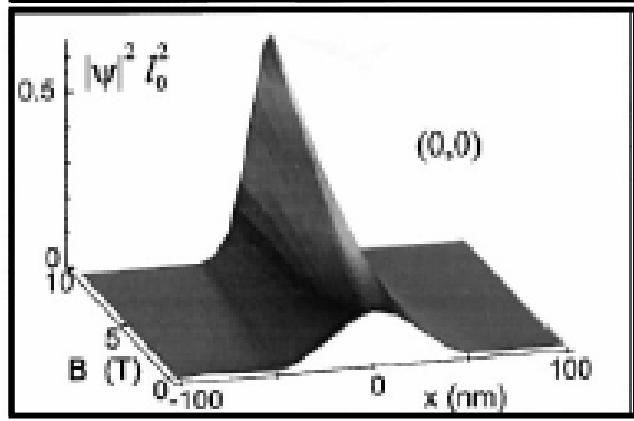
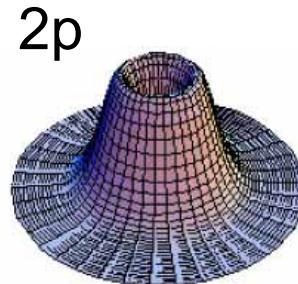
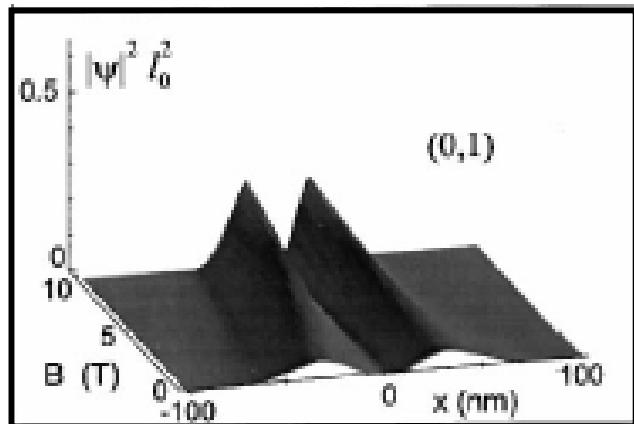


Half-filled shell for
 $N=4, 9, 16, \dots$
♥ total spin
maximized
(exchange)



- Transition in angular momentum (l) accompanied by spin transition
- Exchange energy ~ 0.8 meV
- Spin states well described by 2-electron states (S and T)

Singlet-triplet transition for $N = 2$

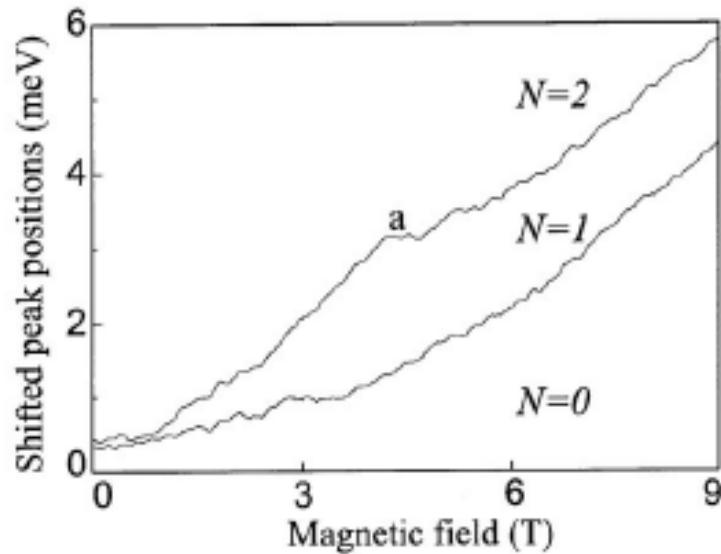


- B_{\perp} shrinks wavefunction ♥ increases Coulomb interaction
- C_{aa} grows faster than C_{ab}
- Exchange interaction favors parallel spins

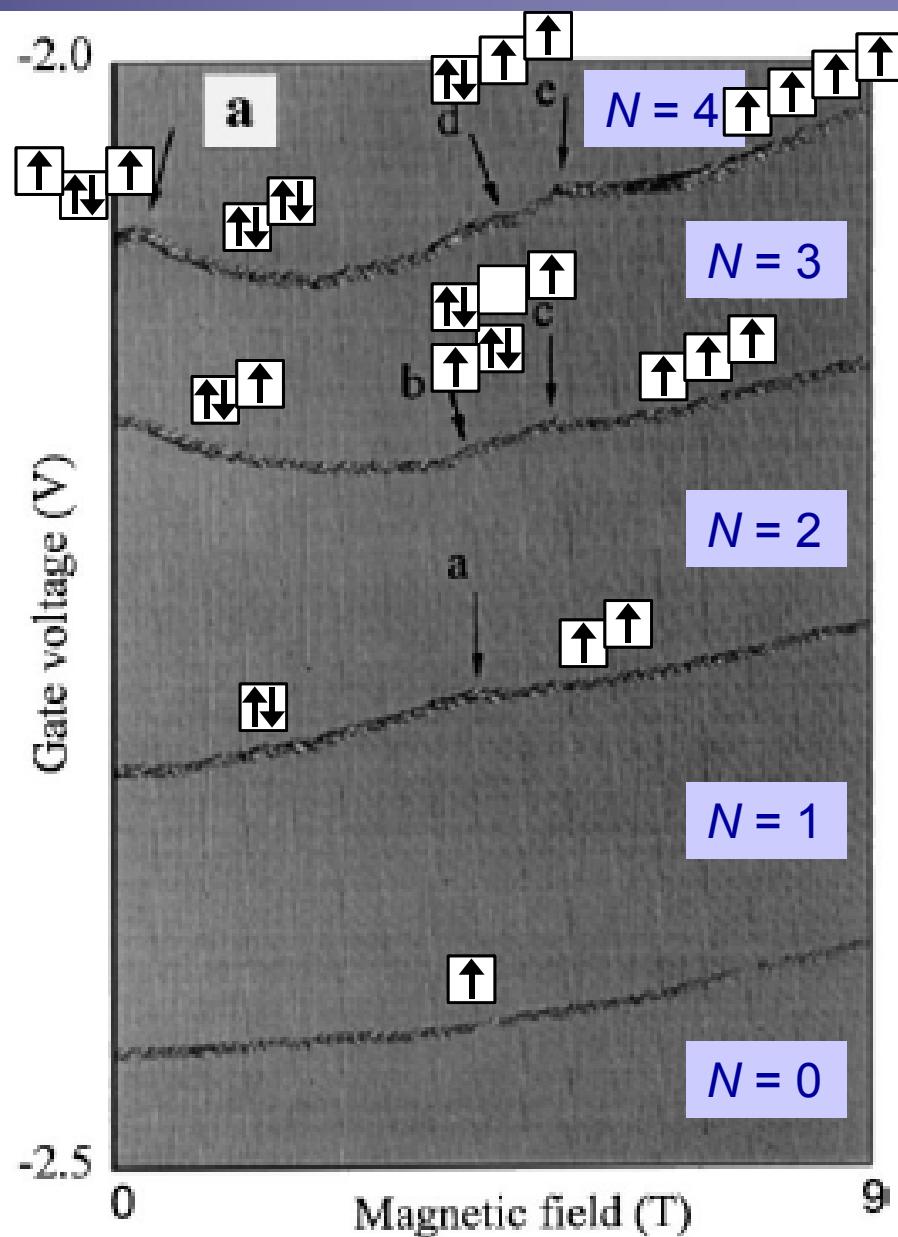


2nd electron in
2p orbital to
form triplet

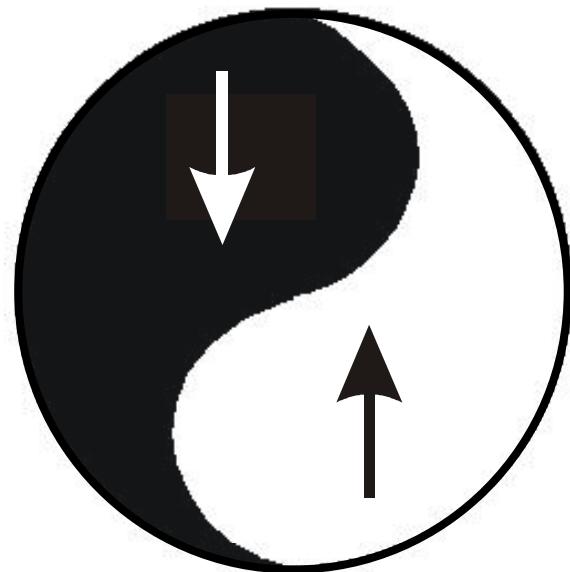
Spin transitions vs. B



- Spins gradually polarized as B increases
- Spin states more complicated as N increases
- ***Single-spin state ($N = 1$) always simple!***

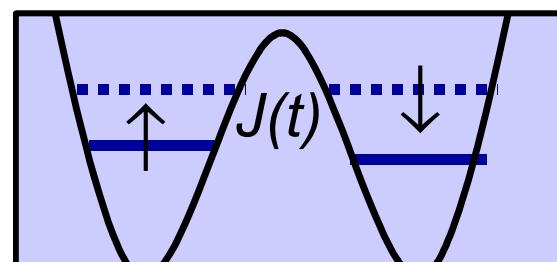
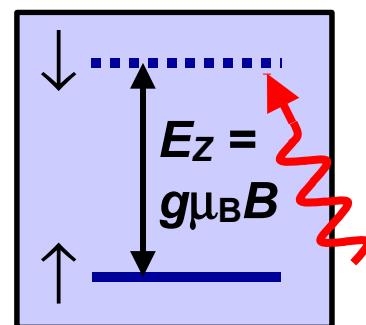
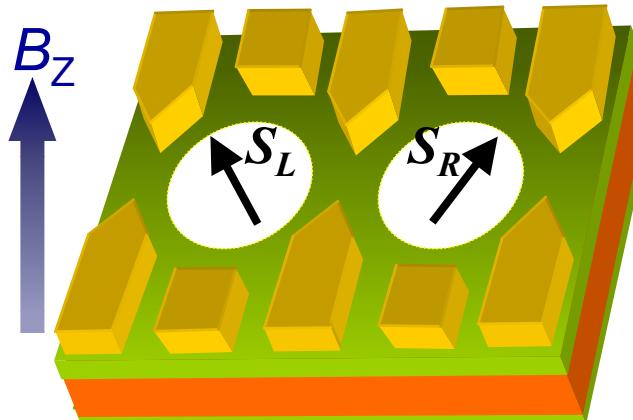


Part III: One-electron quantum dots as spin qubits



Spin qubits

Loss & DiVincenzo, PRA 57, 120 (1998)

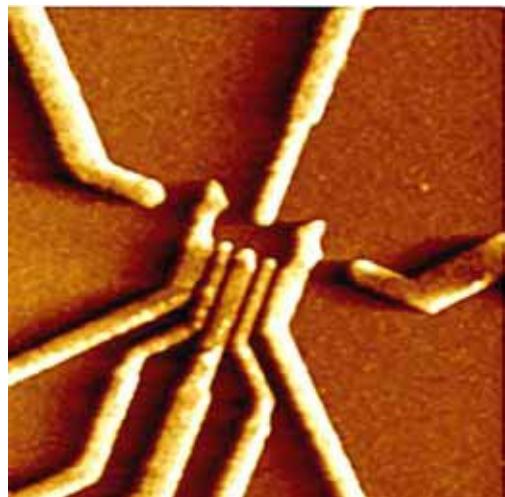


$$H_s(t) = J(t) \mathbf{S}_L \cdot \mathbf{S}_R$$

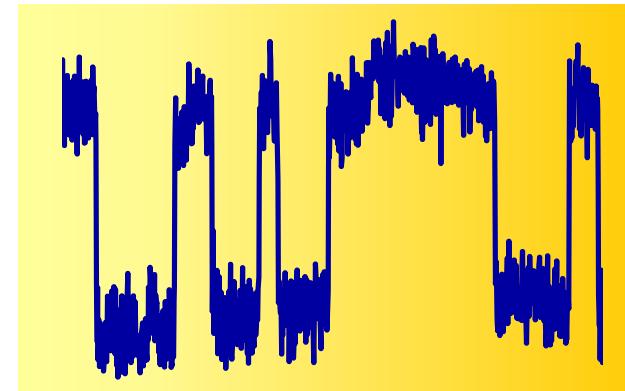
- Qubit defined by Zeeman-split levels of *single electron* in quantum dot
- 1-qubit control:
 - magnetic (ESR-field)
 - electric (modulate effective g-factor)
- 2-qubit coupling: electric (exchange interaction between dots)
- ***Read-out of the qubit state!***

We need...

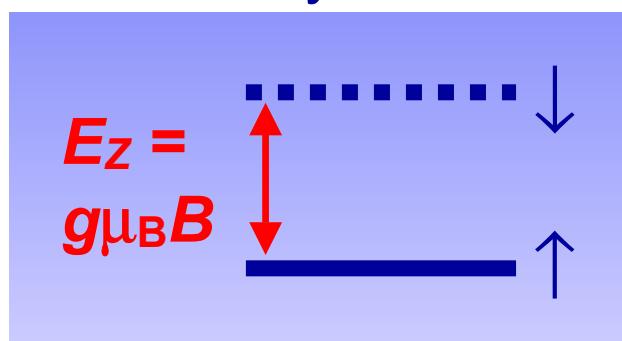
one-electron
double dots...



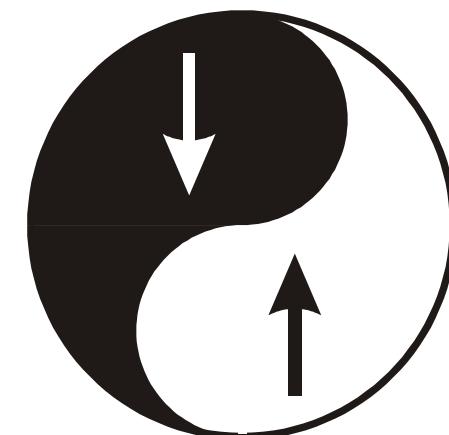
...fast charge
detection...



...two-level
system...



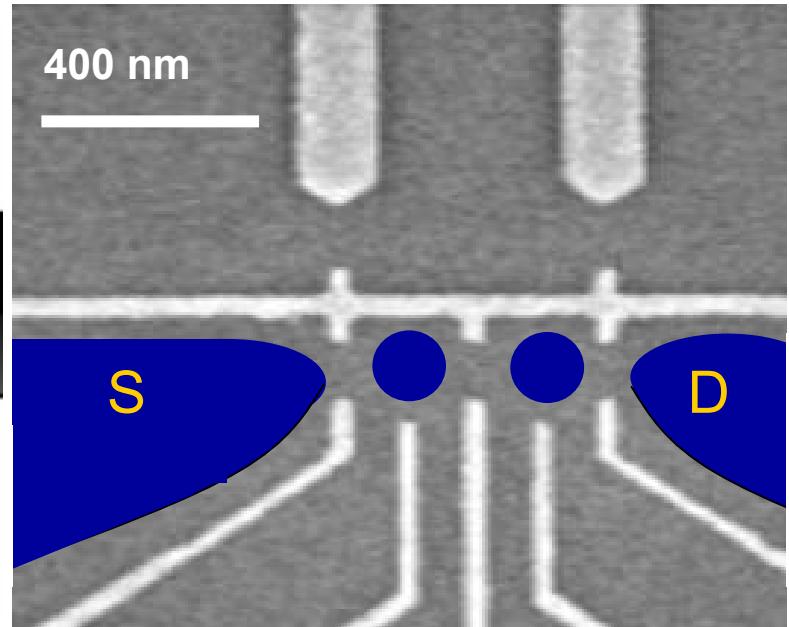
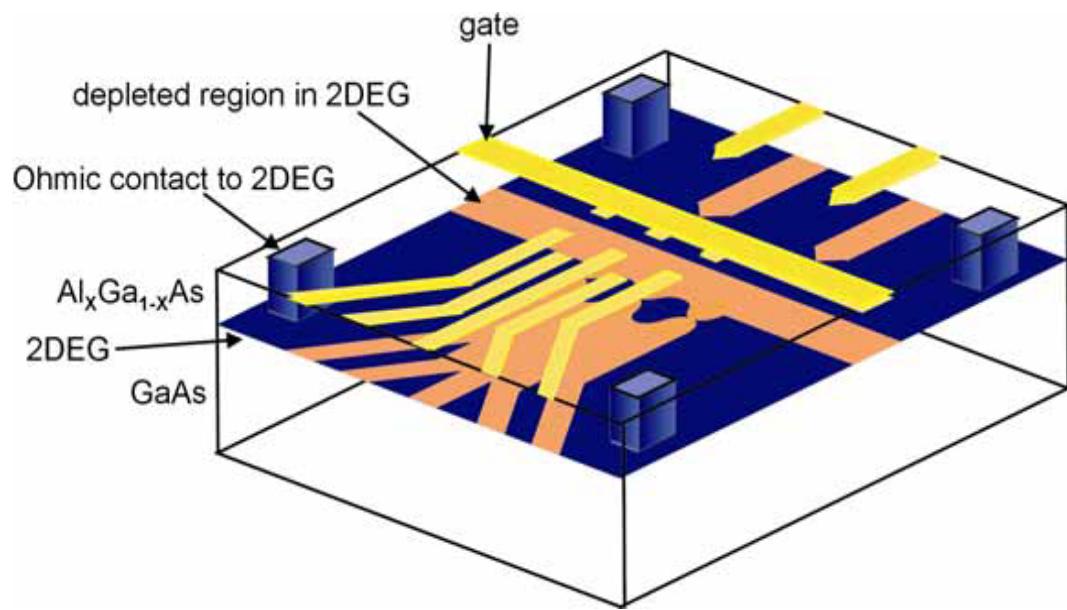
....single spin
measurement!



Few-electron quantum dot circuit with integrated charge detector



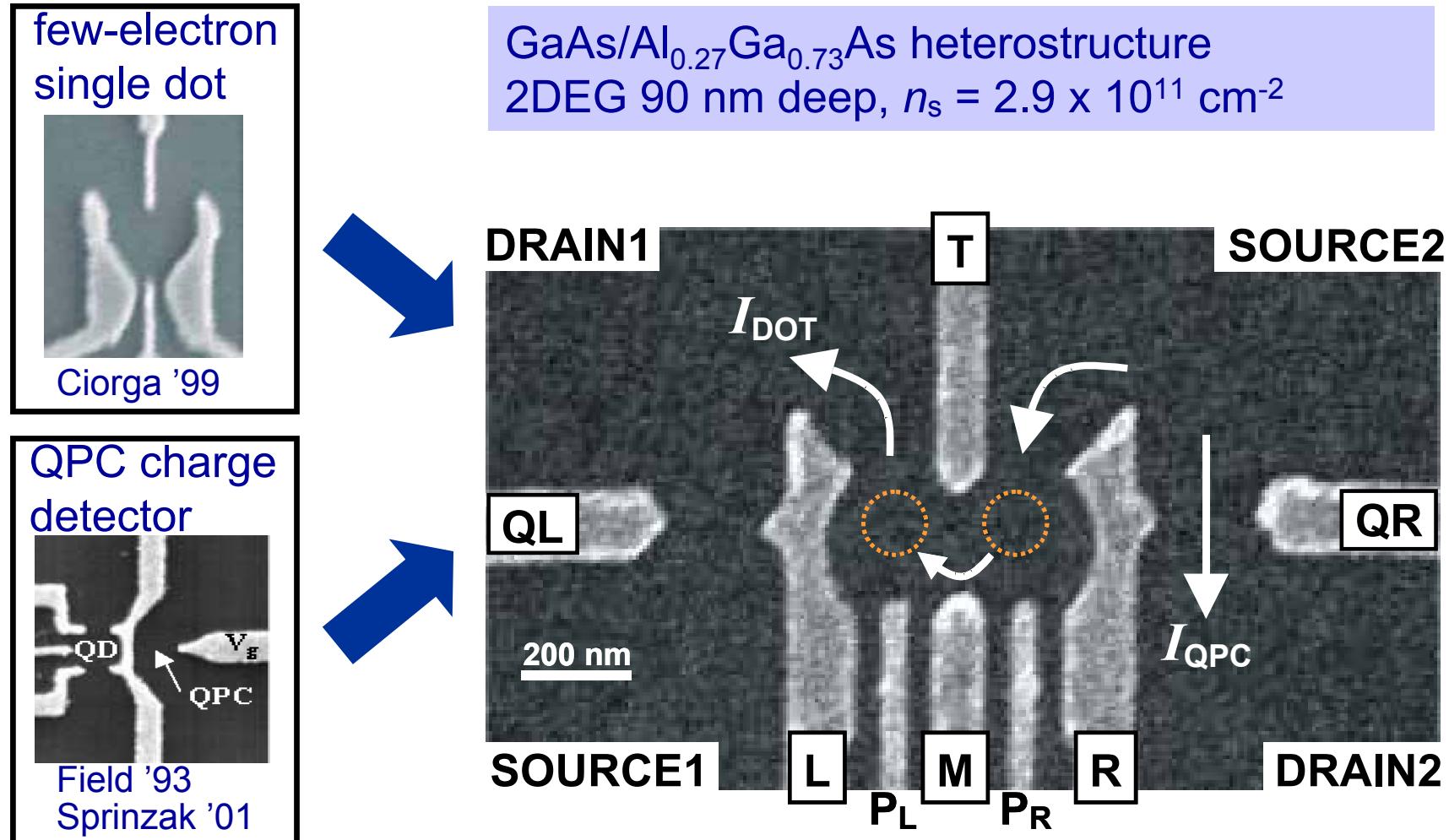
N = 1 in lateral quantum dots?



- High-mobility 2DEG ($\sim 10^6 \text{ cm}^2/\text{Vs}$)
- Density $\sim 10^{11} \text{ cm}^{-2}$ ♥ $\lambda_F \sim 30 \text{ nm}$
- Smallest gate structure $\sim 40 \text{ nm}$
- Dot size $\sim 200 \text{ nm}$
 ♥ discrete energy spectrum

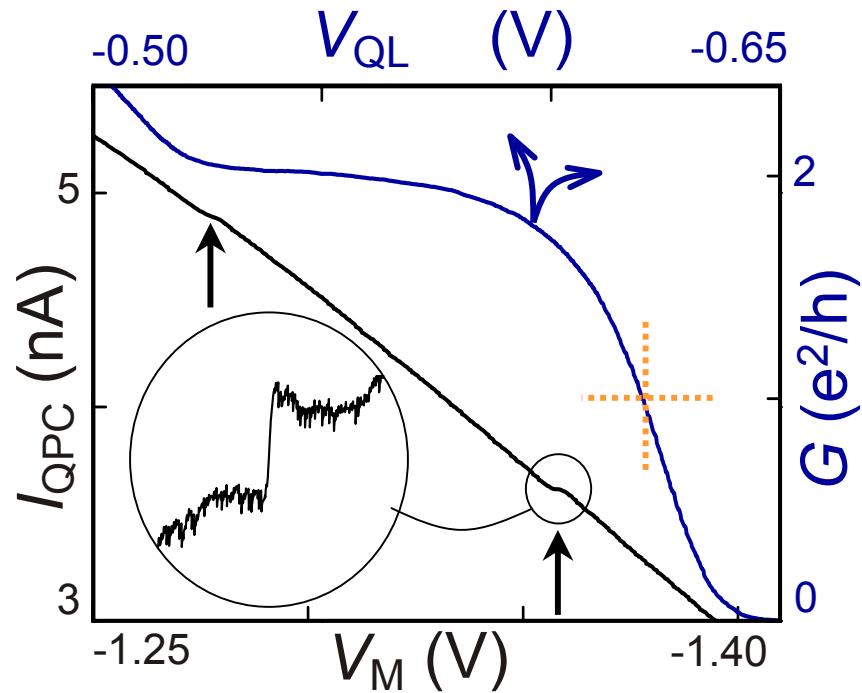
BUT: Barriers close as dot is depleted ♥ current too small to measure!

Lateral few-electron double dot

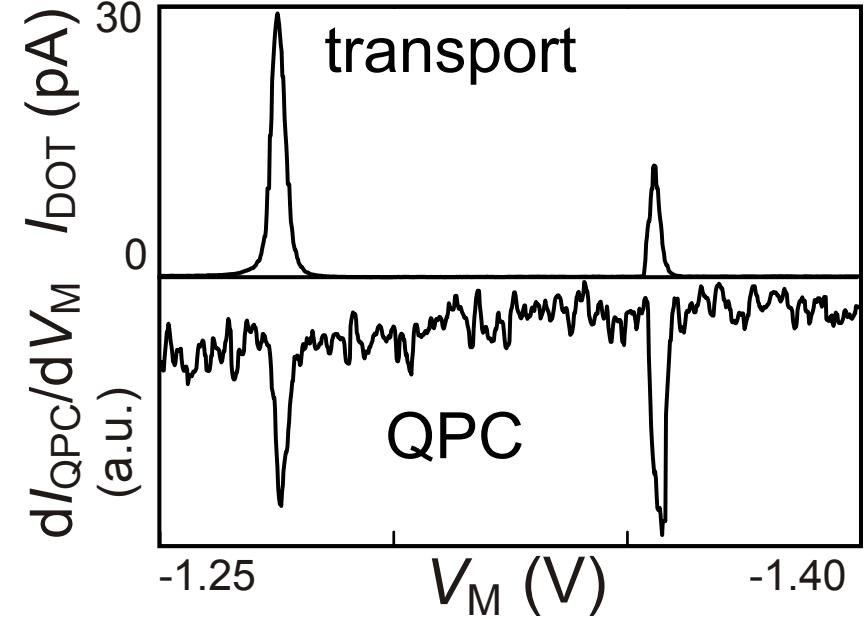


Design: barriers don't close as DQD is depleted

QPC - detector for single dot

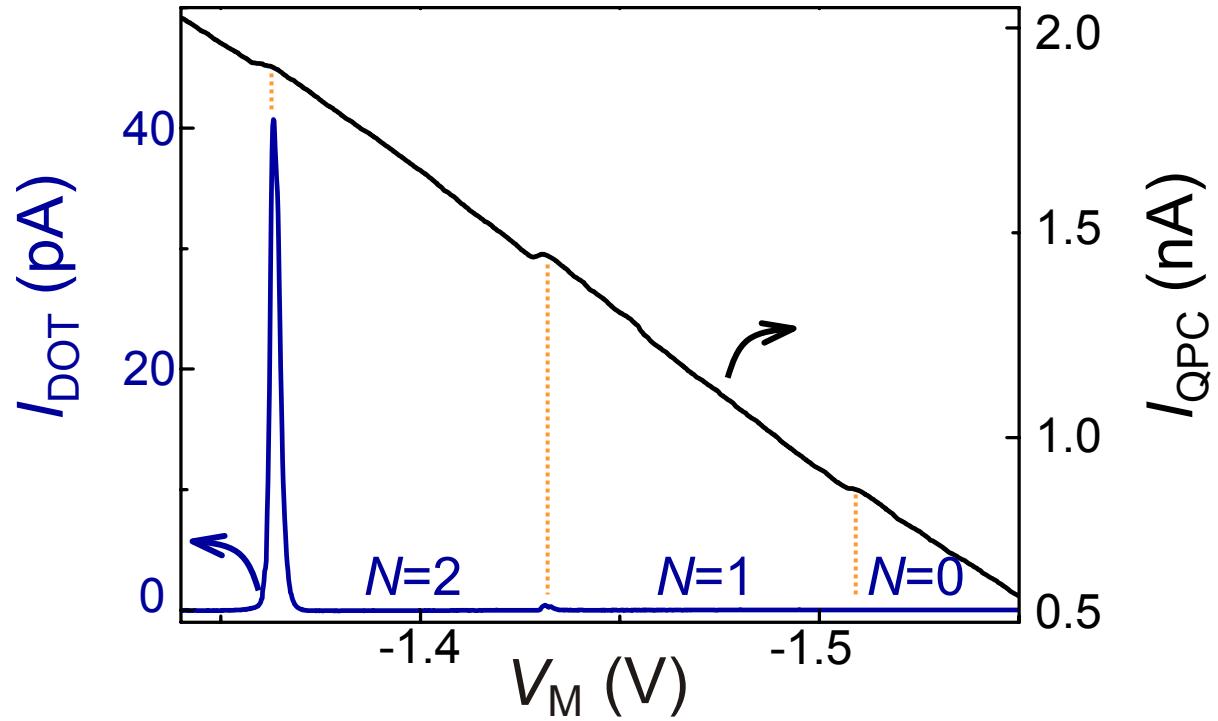


- Tune to steepest point ($G \sim e^2/h$)
- Jumps in I_{QPC} ♥ change in electron number



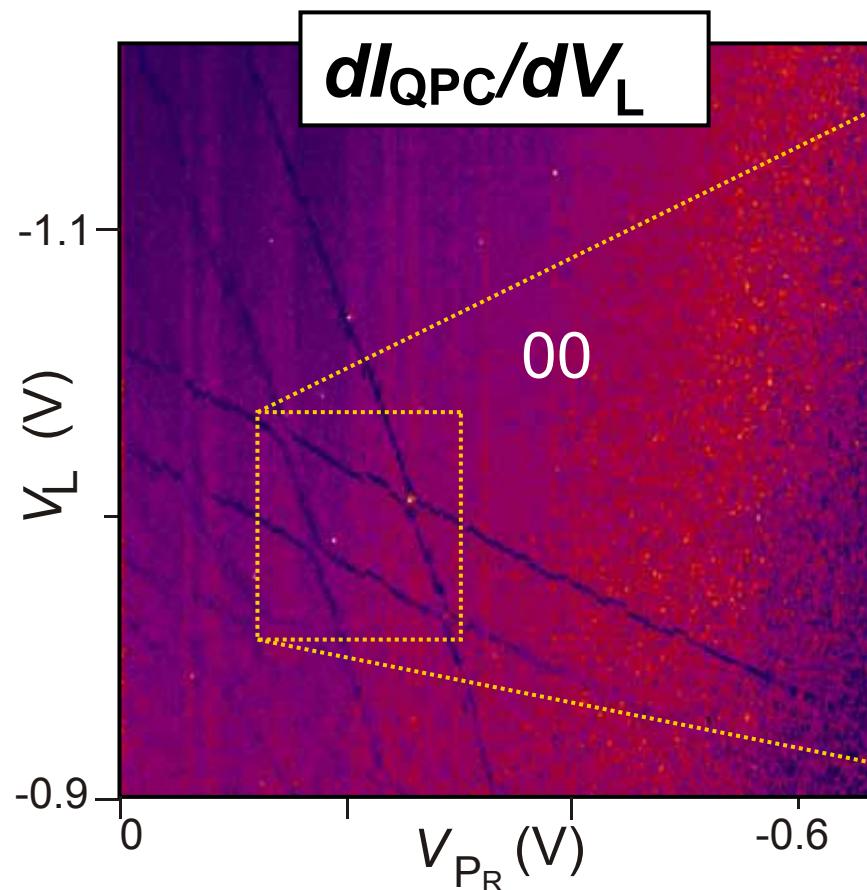
- Modulate V_M with lock-in
♥ measure dI_{QPC}/dV_M
- **Dips in QPC-signal coincide with Coulomb peaks in transport!**
- Sensitivity $\sim 0.1e^-$ (at 17 Hz)

QPC - detector for isolated single dot

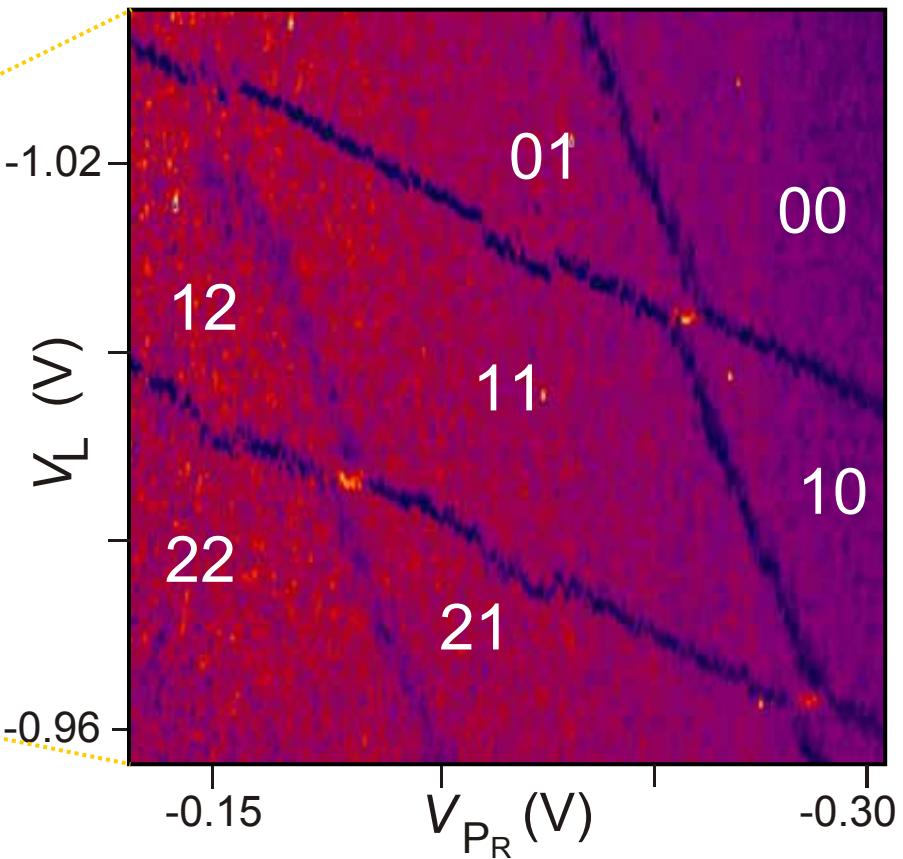


QPC can detect charge
transitions in *isolated* QD!

QPC - detector for double dot

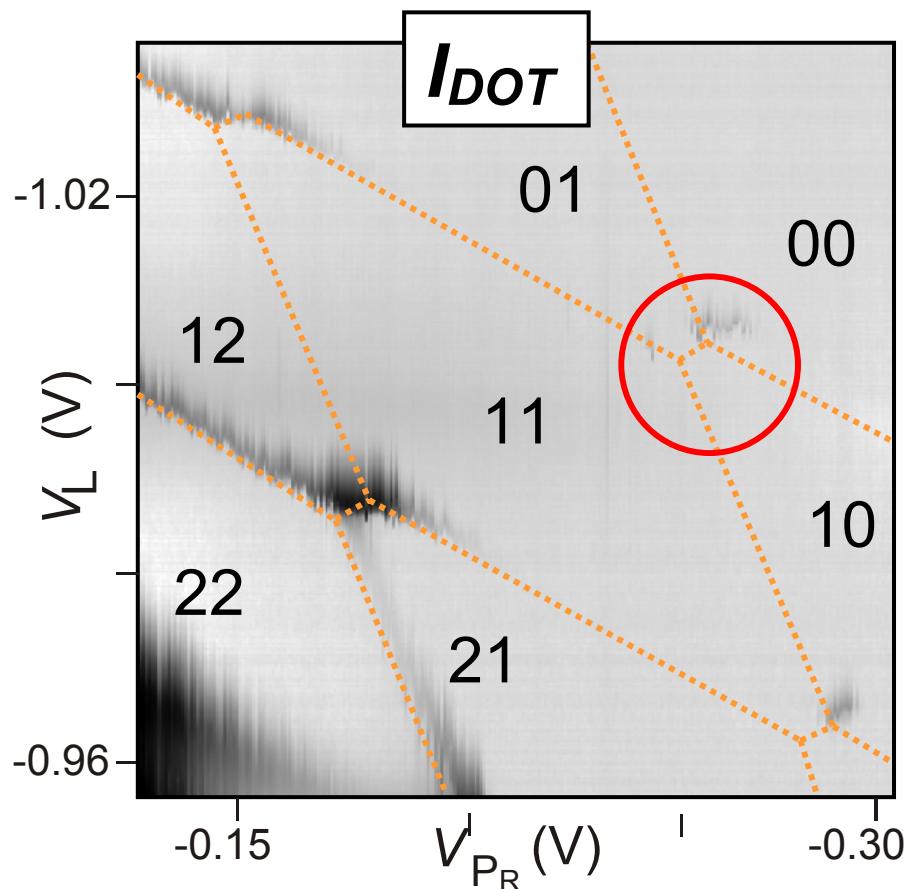


- Double dot emptied completely !
- $E_C = 4.5 \text{ meV}$

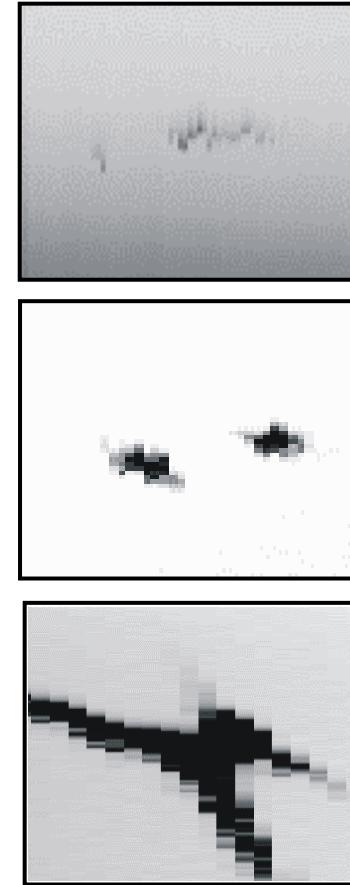


QPC detects *all* charge transitions, also between dots

Tunable few-electron double dot

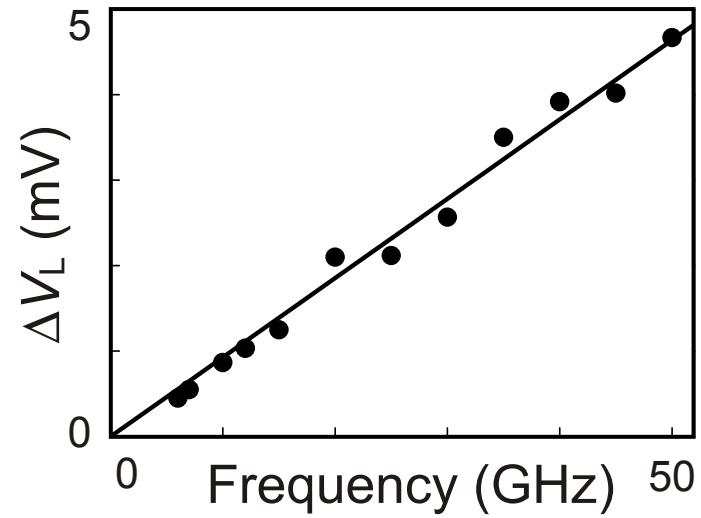
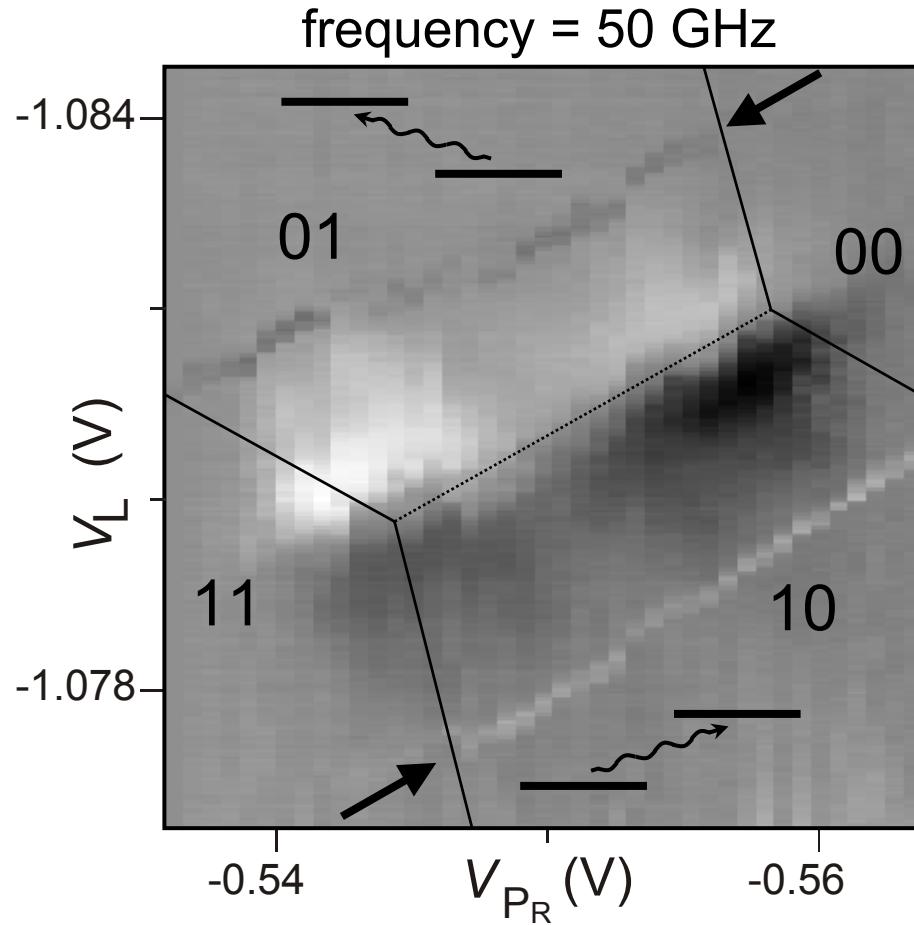


All triple points visible
♥ **barriers still open**



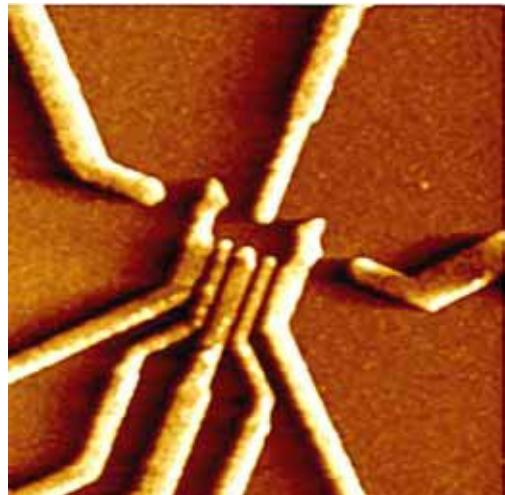
Double dot
fully tunable!

Photon-assisted tunneling



- Photon-assisted current visible for $f > 6$ GHz
- Always linear ♥ coherent coupling between two dots not visible

Summary...



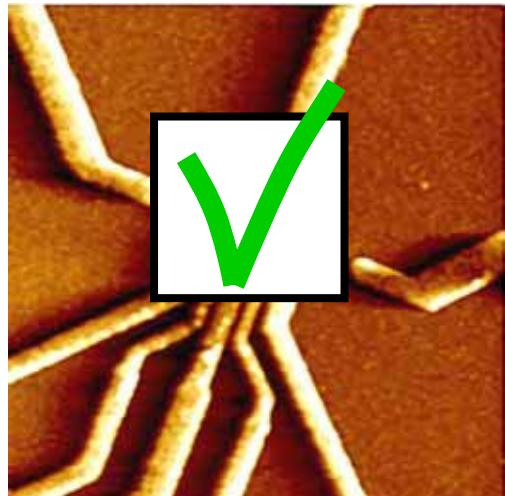
We can:

- **isolate *single electron spin* in (double) quantum dot**
- **study it using transport or charge detection**

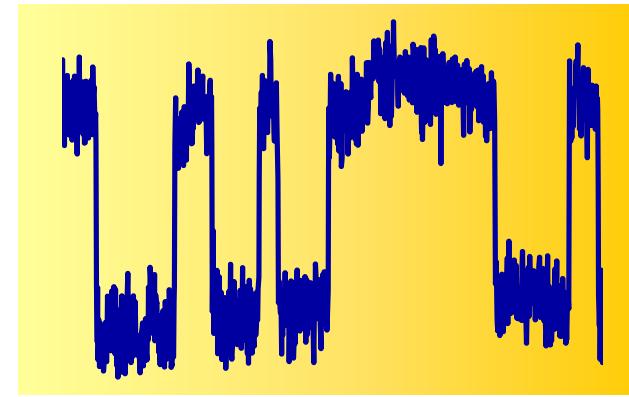
PRB 67, 161308(R) (2003)

We need...

one-electron
double dots...

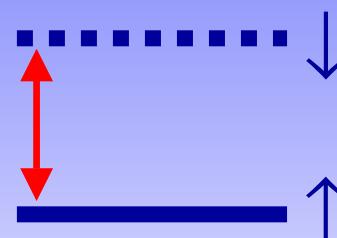


...fast charge
detection...

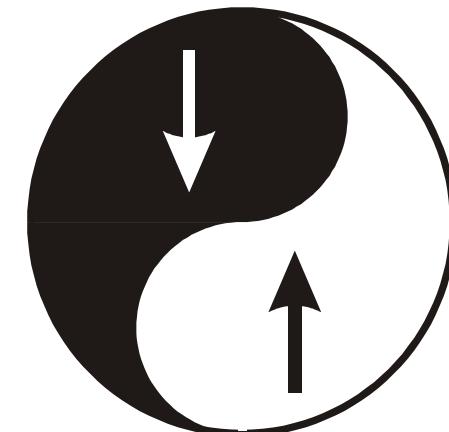


...two-level
system...

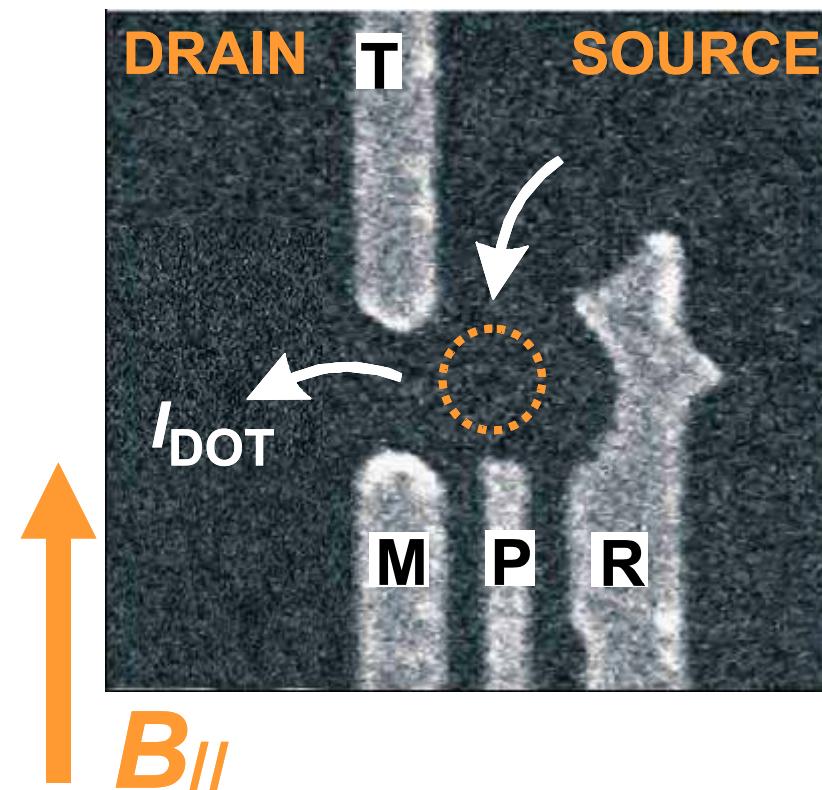
$$E_z = g\mu_B B$$



....single spin
measurement!

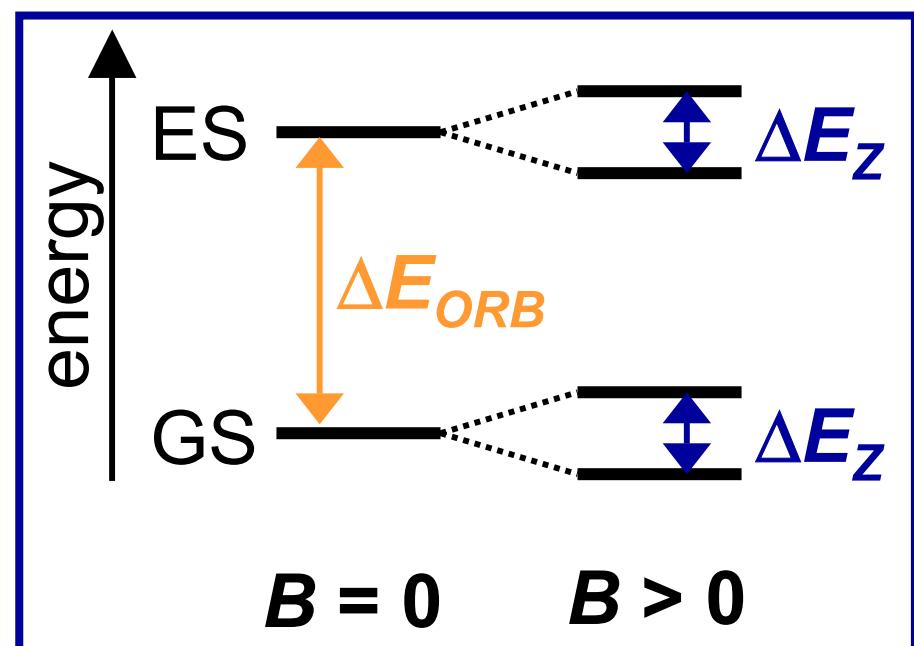


Zeeman splitting in an artificial Hydrogen atom

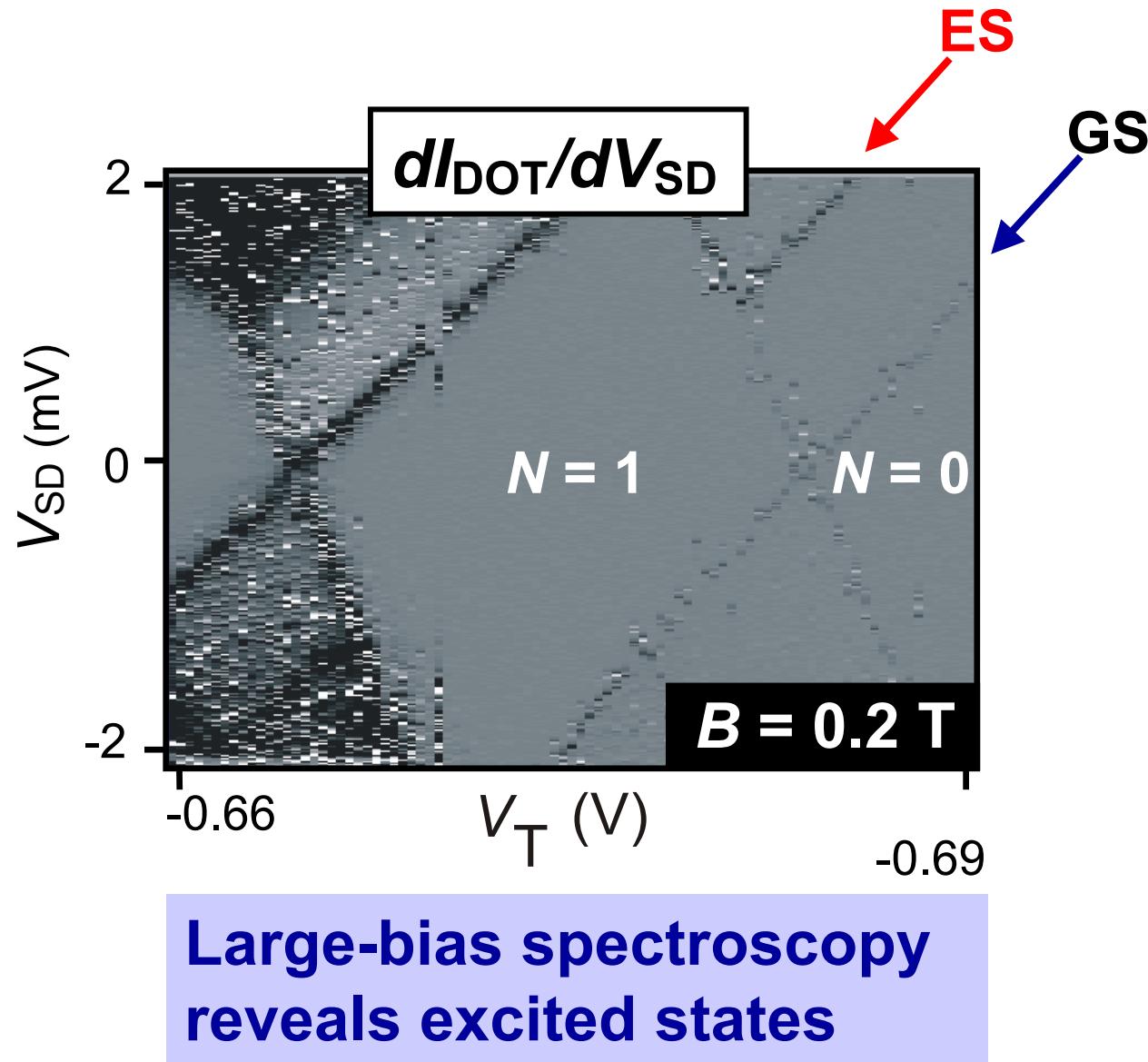


magnetic field parallel
to 2DEG

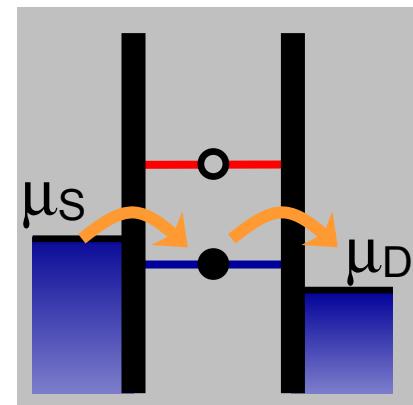
Similar double dot device:
other gates grounded



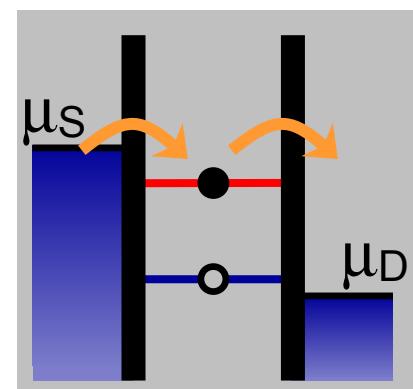
$N = 1$ Coulomb diamond



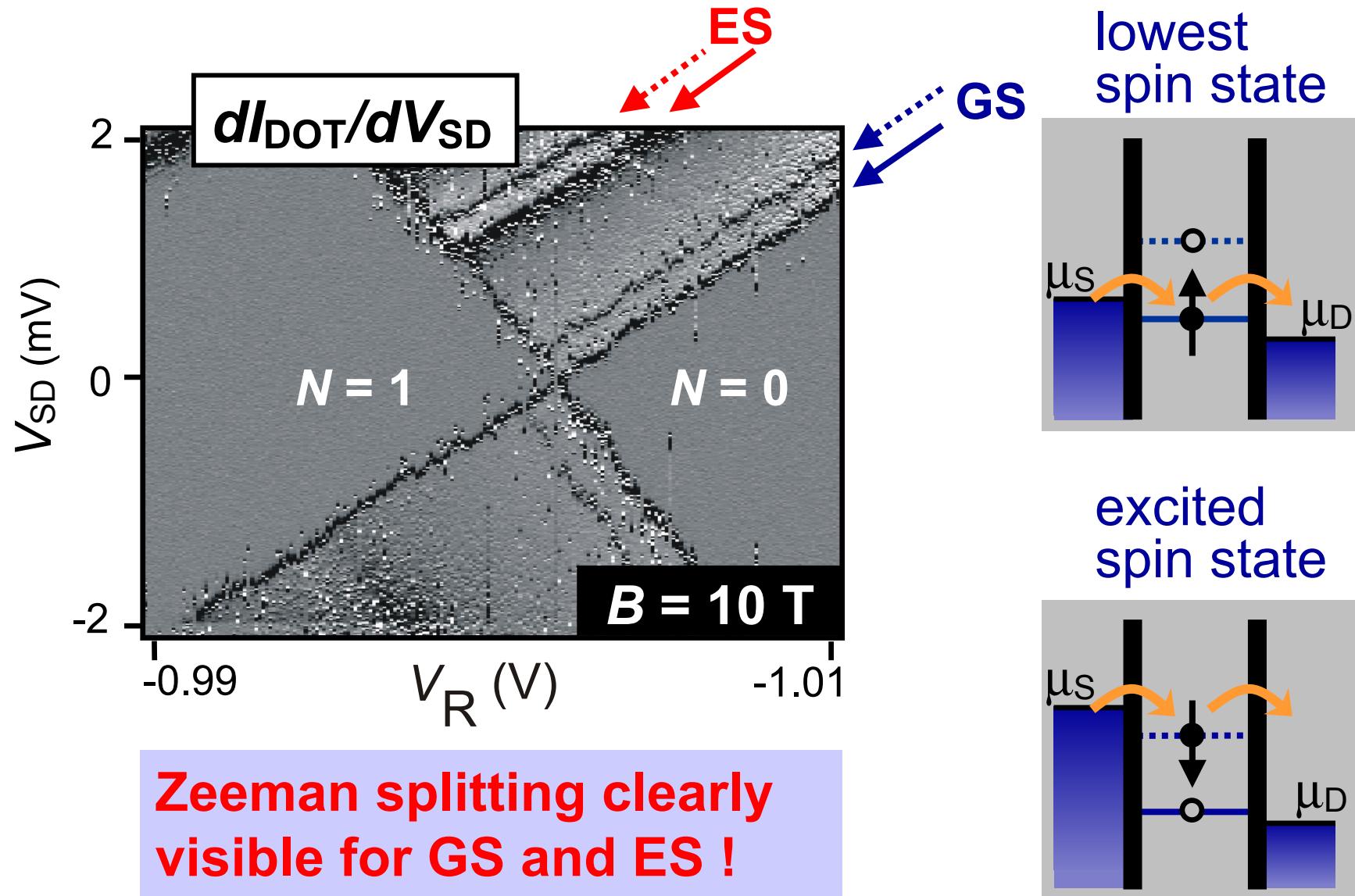
$N = 1$ orbital ground state



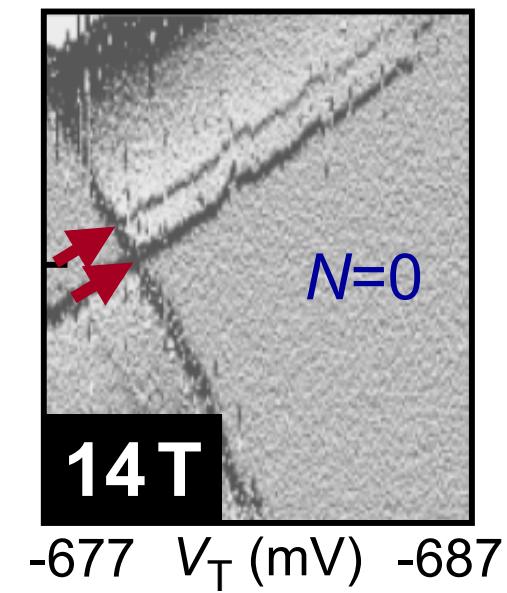
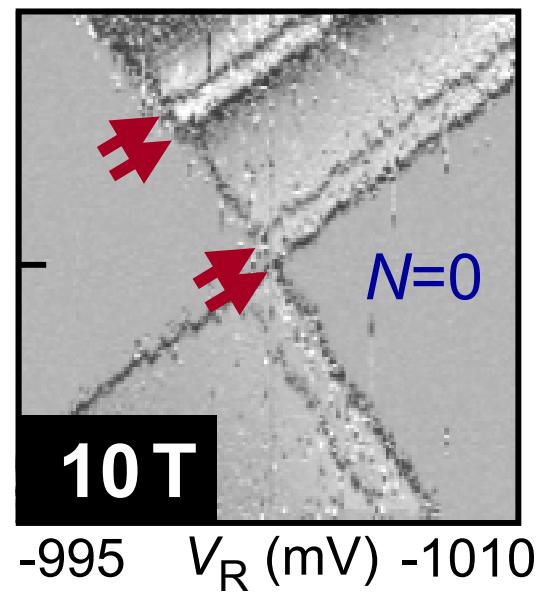
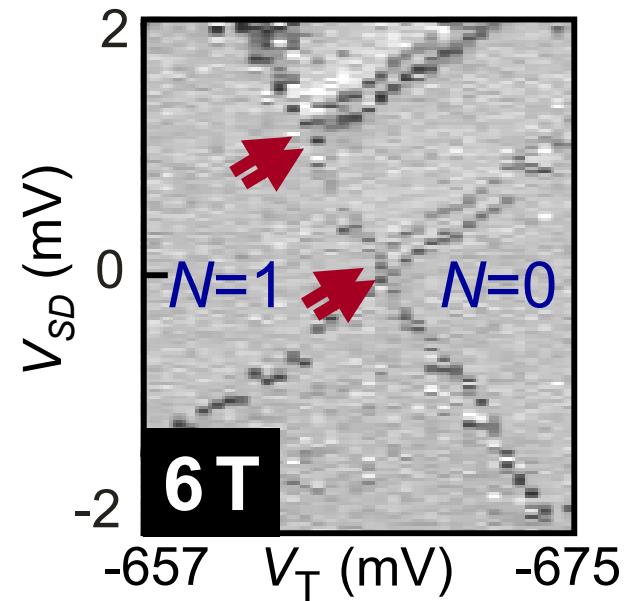
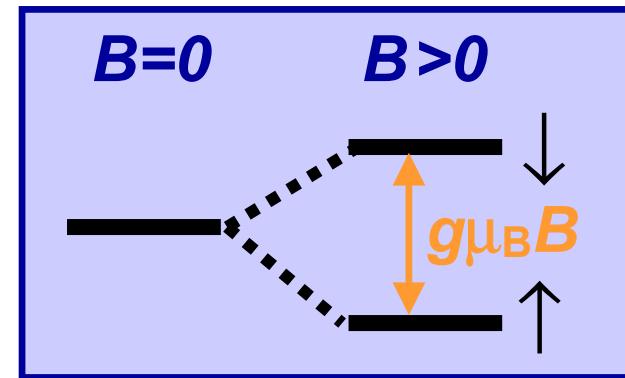
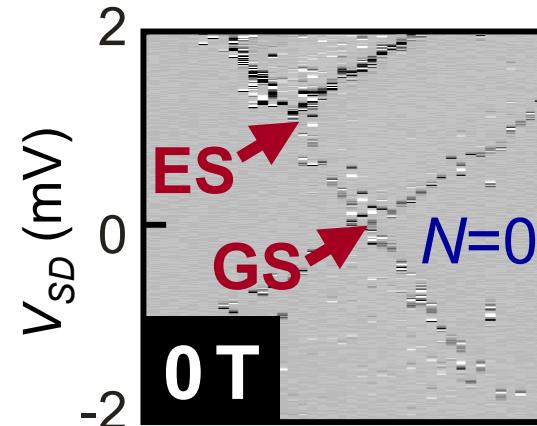
$N = 1$ orbital excited state



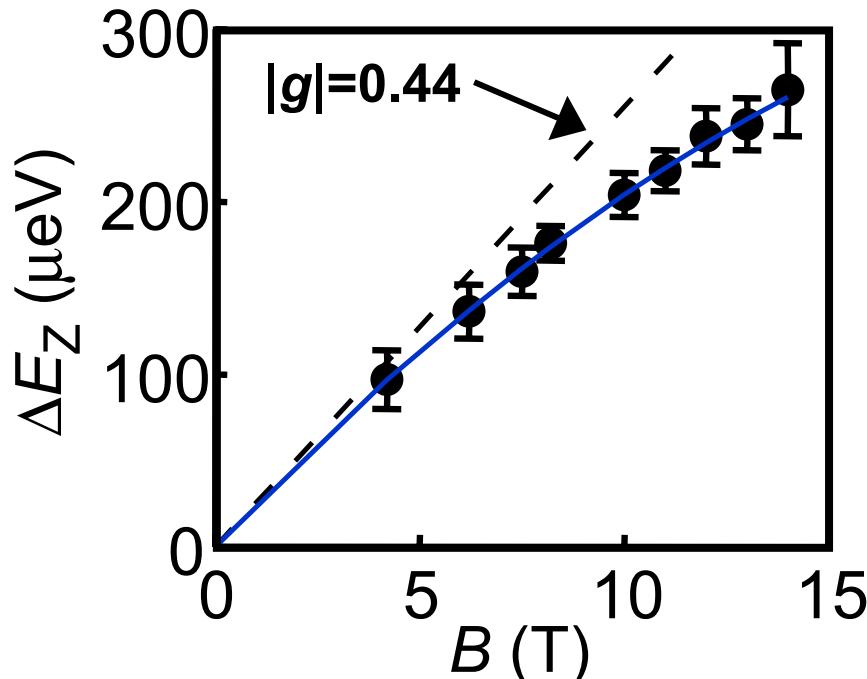
$N = 1$ Zeeman splitting in B_{\parallel}



$N = 1$ Zeeman splitting in B_{\parallel}



Spectroscopy of qubit 2-level system

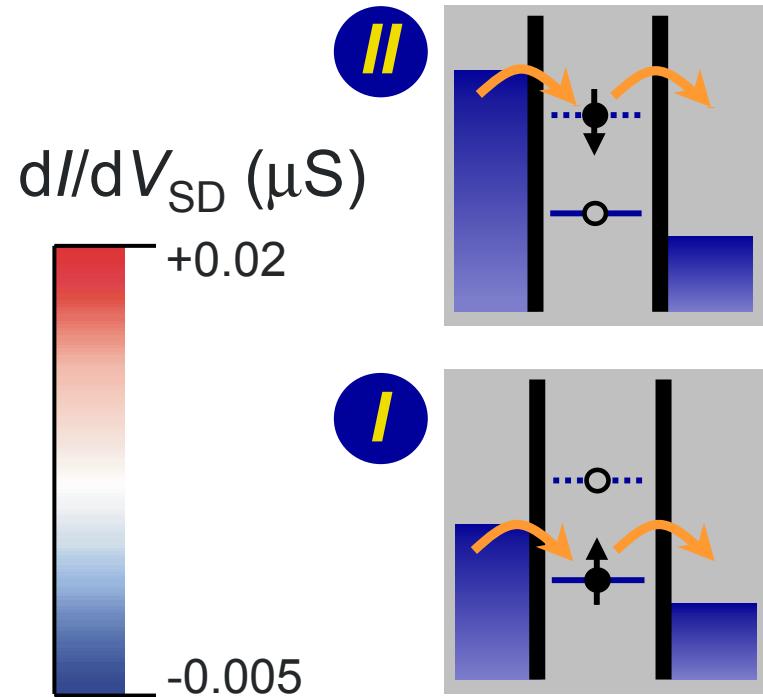
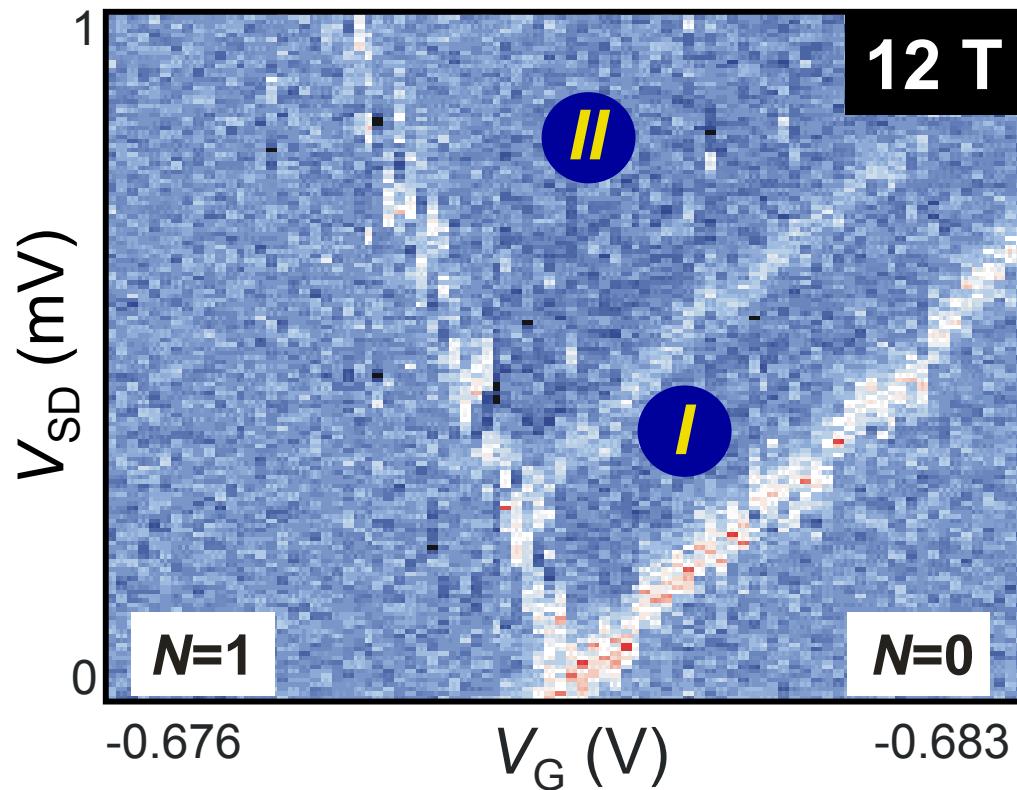


$$E_z = g \mu_B B_{||}$$

$$|g| = 0.43 - 0.008B$$

N = 1 dot as filter for spin-up

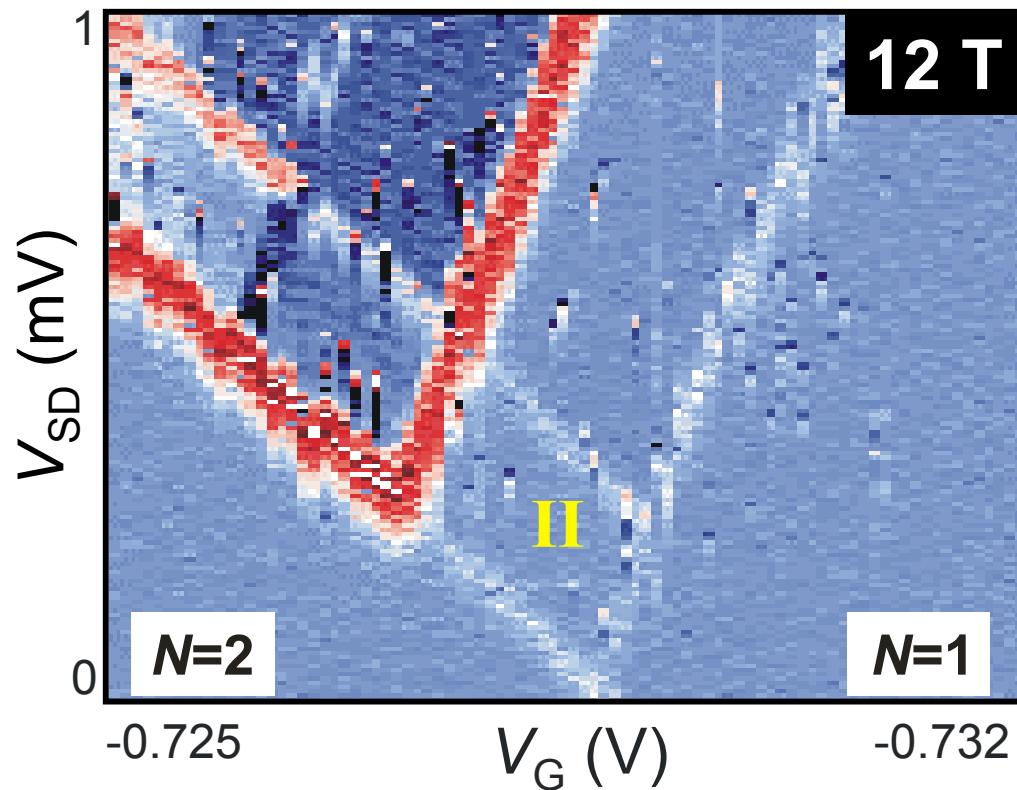
theory: Recher et al., PRL 85,1962 (2000)



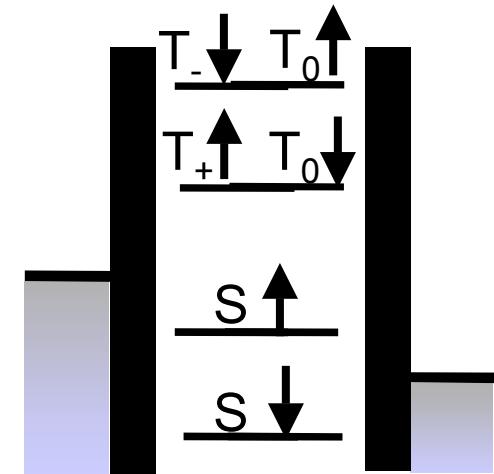
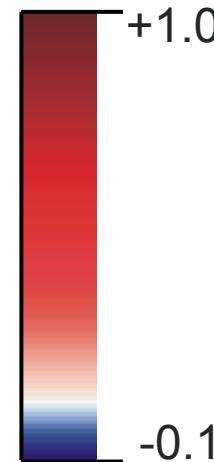
Quantum dot in region I acts as *spin filter*; only spin-up electrons can pass through

N = 2 dot as filter for spin-down

theory: Recher et al., PRL 85,1962 (2000)



d//dV_{SD} (μ S)



Quantum dot in region II acts as
spin filter; only **spin-down**
electrons can pass through

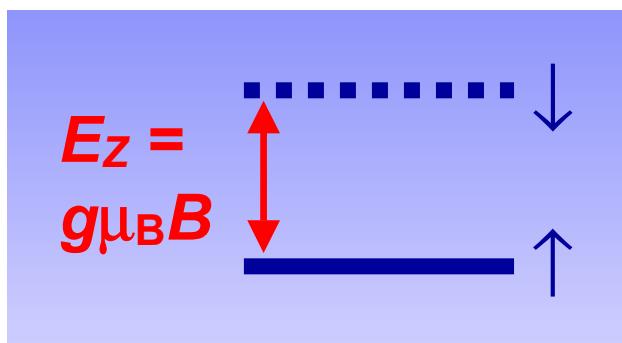
Summary...

We have:

- identified (stable) two-level system
- ($T_1 > 50 \mu\text{s}$)
- bipolar spin filter

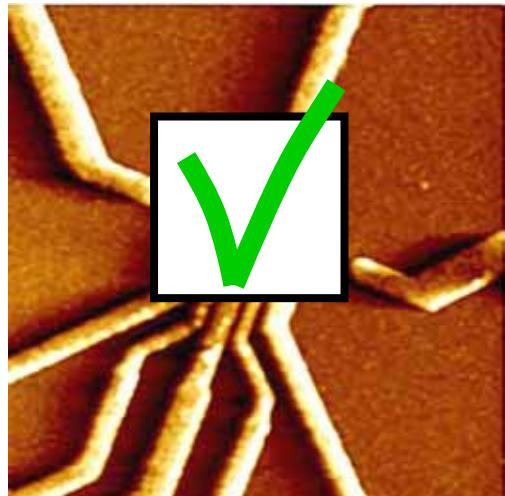
Hanson *et al.*, PRL 91, 196802 (2003)

Cond-mat/0311414

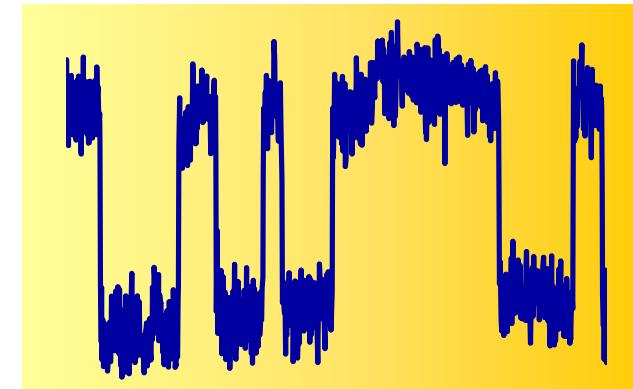


We need...

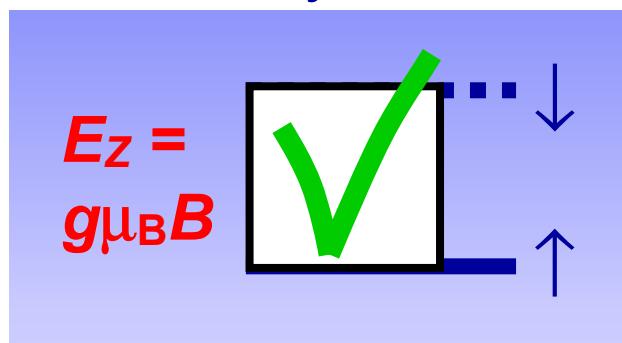
one-electron
double dots...



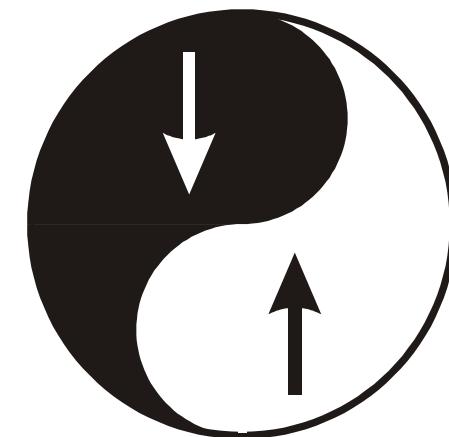
...fast charge
detection...



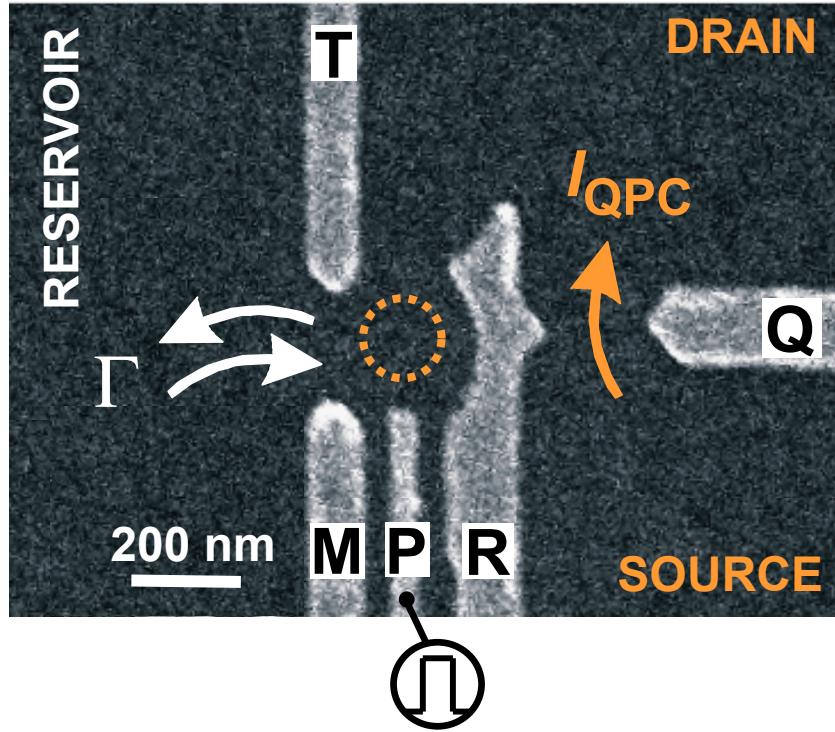
...two-level
system...



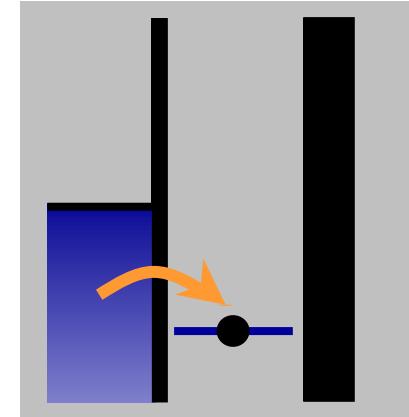
....single spin
measurement!



Quantum Point Contact as a fast charge detector

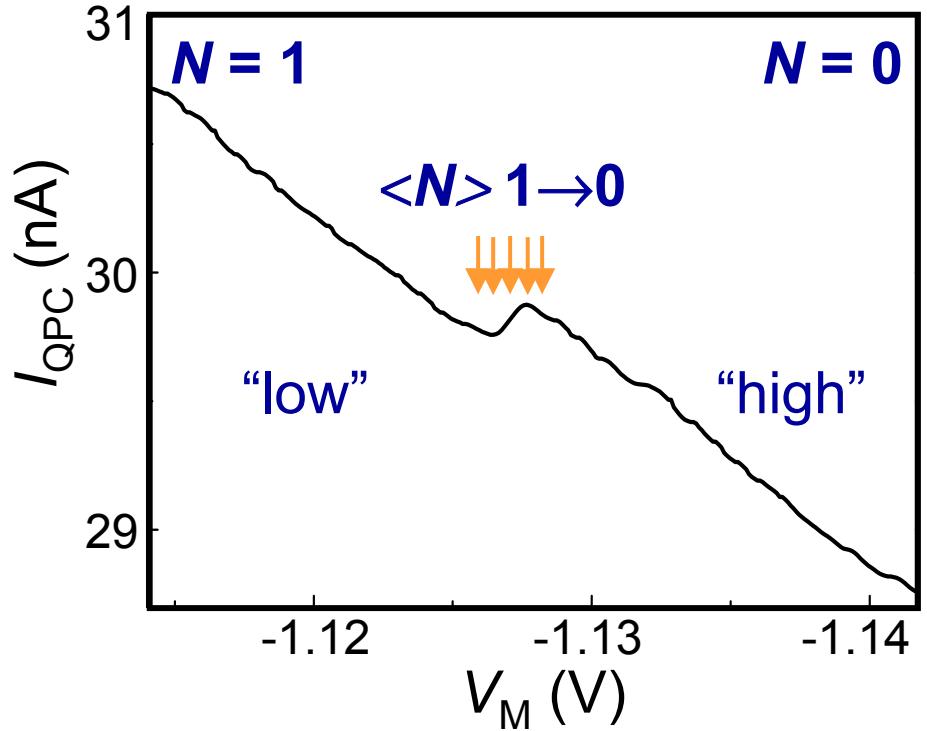
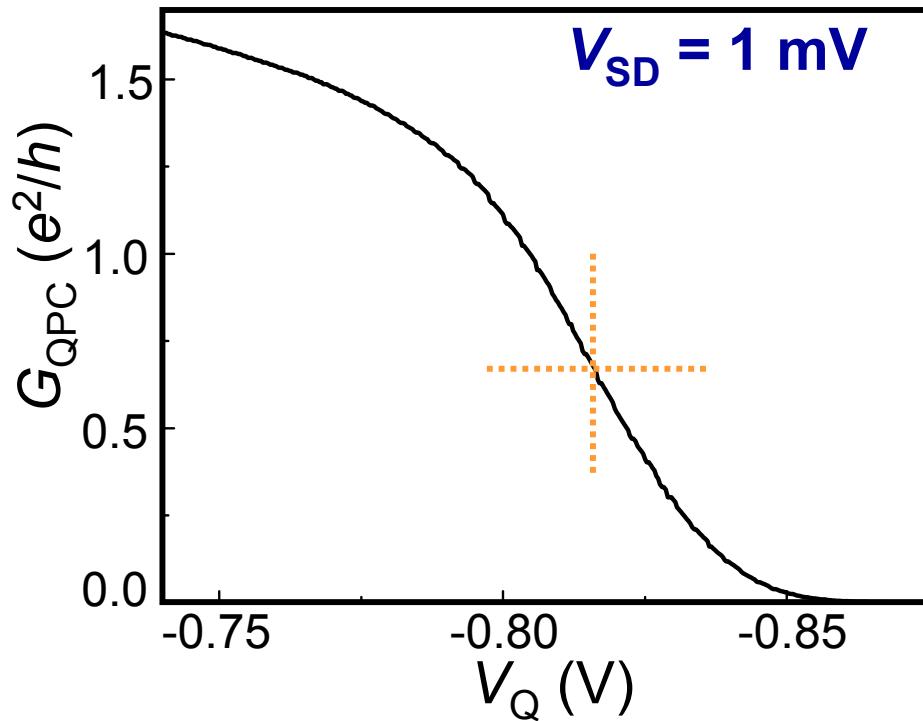


- Tunnel barrier to QPC-channel closed completely
- QD weakly connected to reservoir
- Detect individual tunnel events



- Fast IV-converter:
100 kHz, $0.8\text{nV}/\text{Hz}^{1/2}$
- Fast ISO-amp:
300 kHz
- Operating bandwidth:
40 kHz
- Shot noise limit:
100 MHz

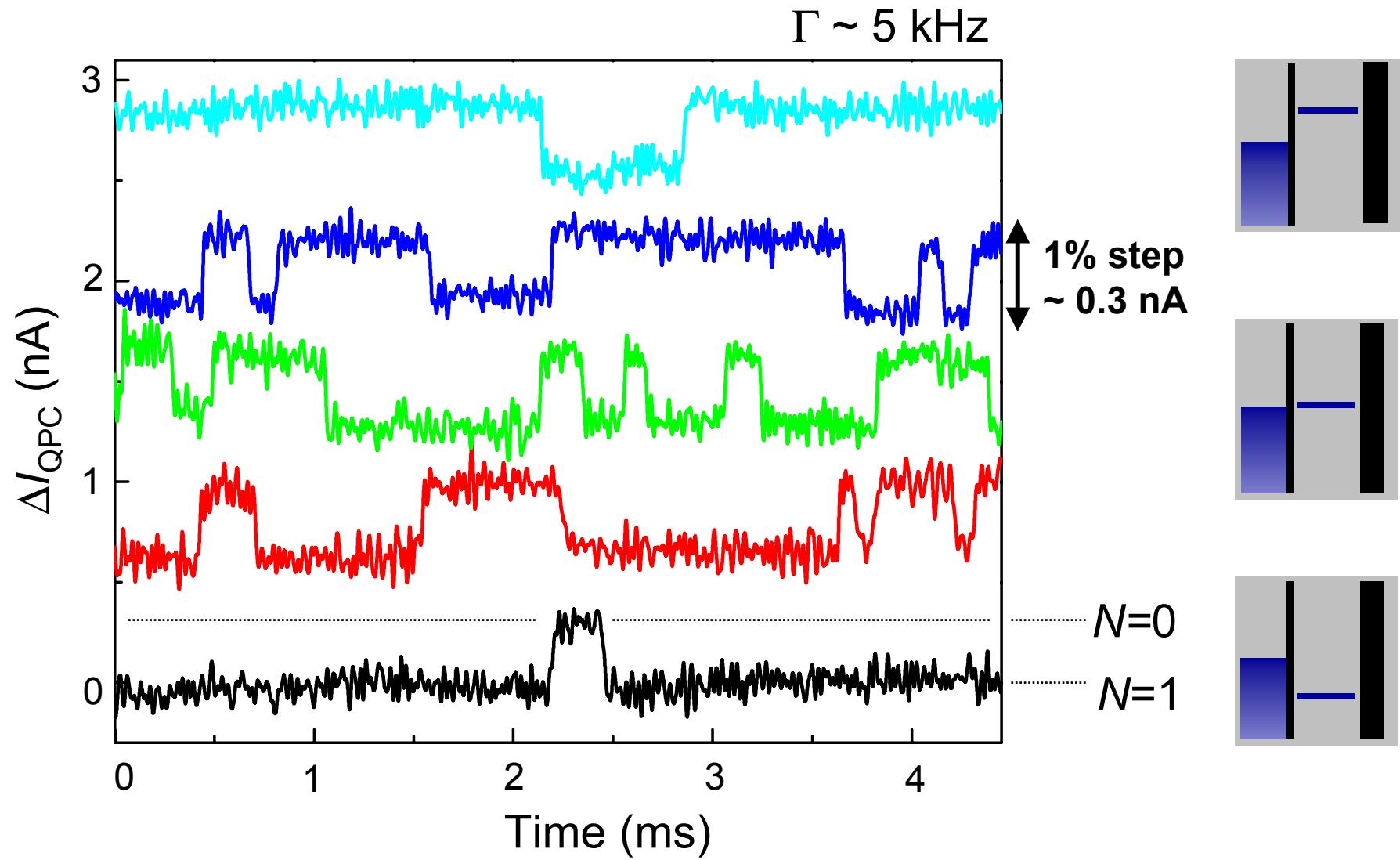
QPC average charge detection (dc)



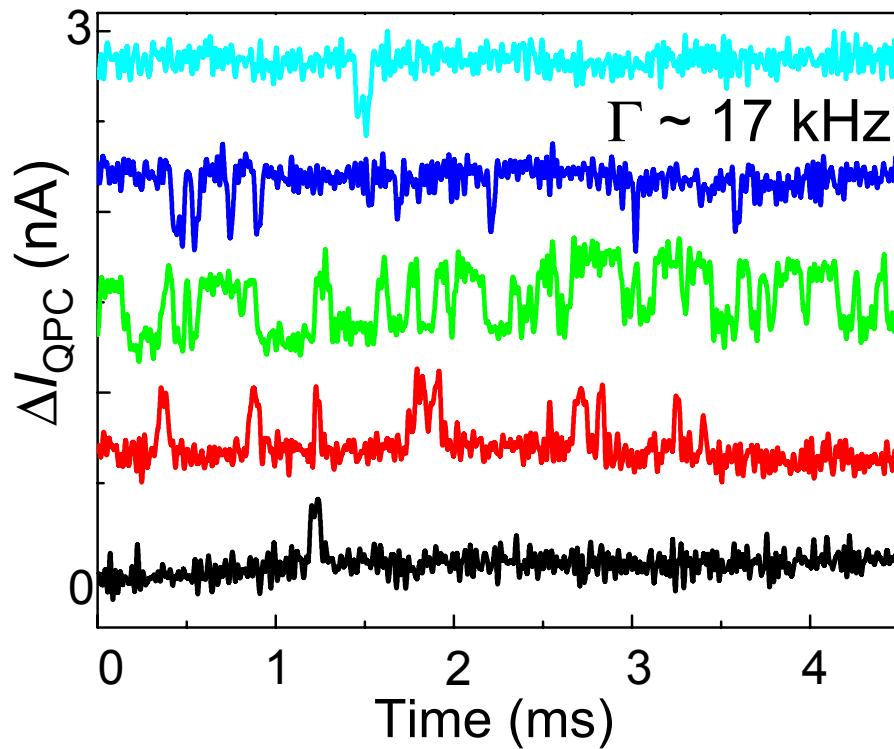
- $V_{\text{SD}} = 1 \text{ mV}$
- $G_{\text{QPC}} \sim 0.5 - 1.0 \text{ } e^2/h$
- $I_{\text{QPC}} \approx 30 \text{ nA}$

- Sweep dot-gate (V_M)
- I_{QPC} increases ($\sim 1\%$) when $\langle N \rangle$ from 1 to 0

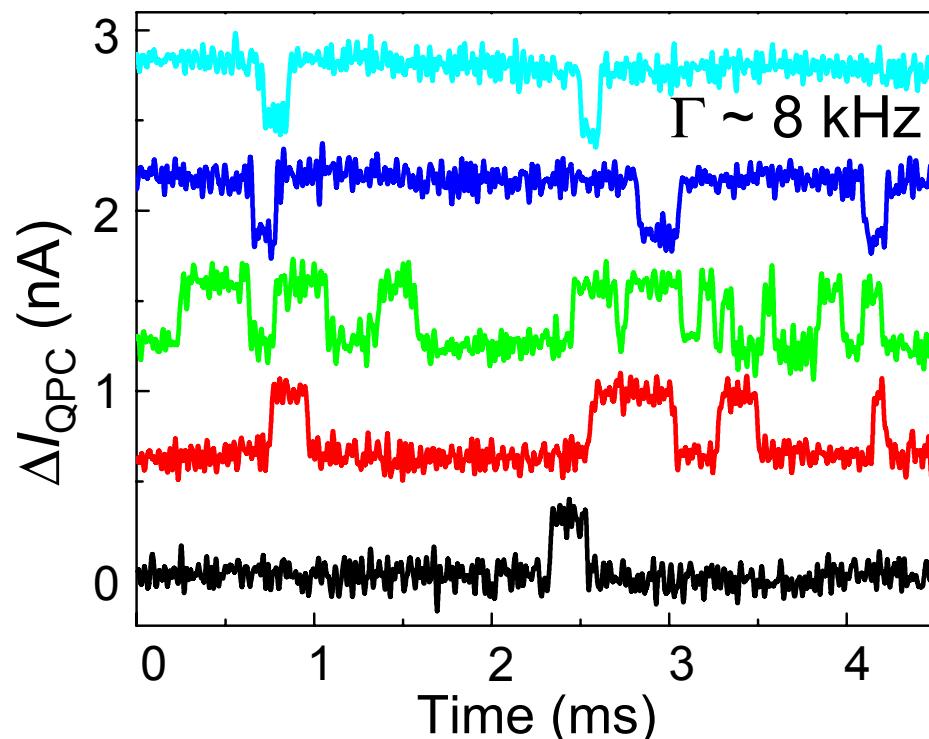
Real-time single-electron tunneling



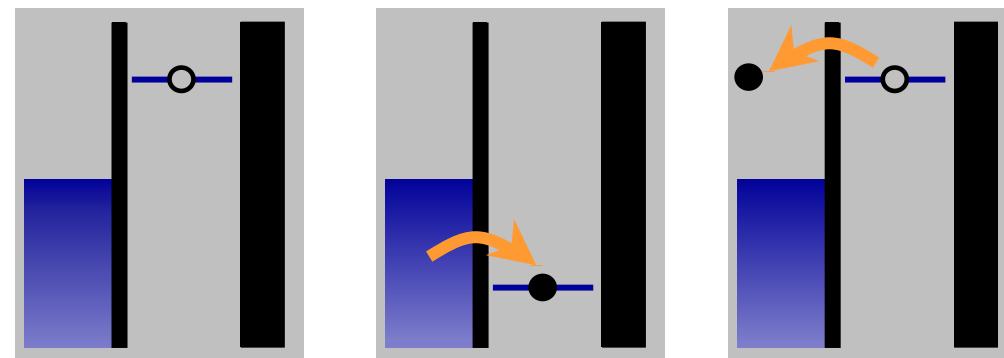
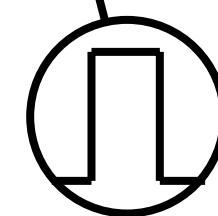
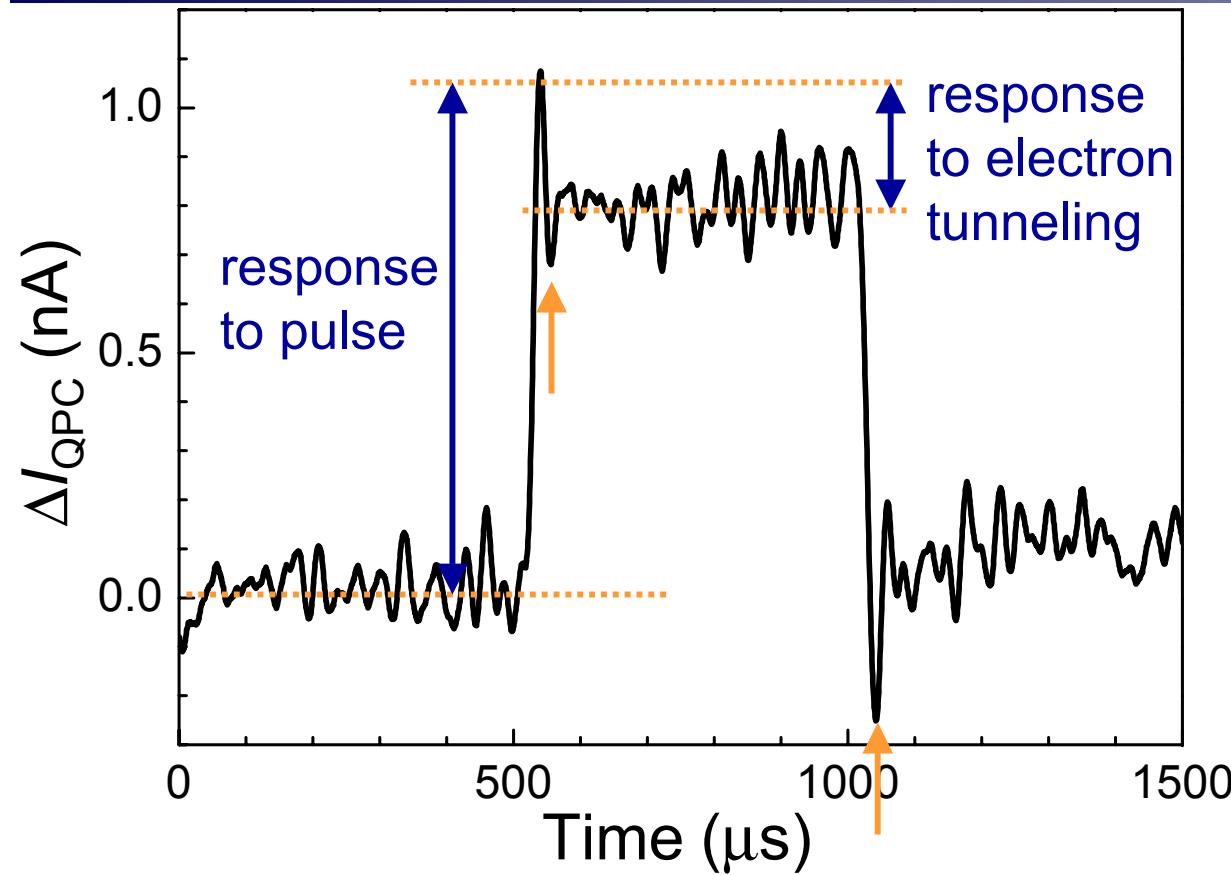
Real-time single-electron tunneling



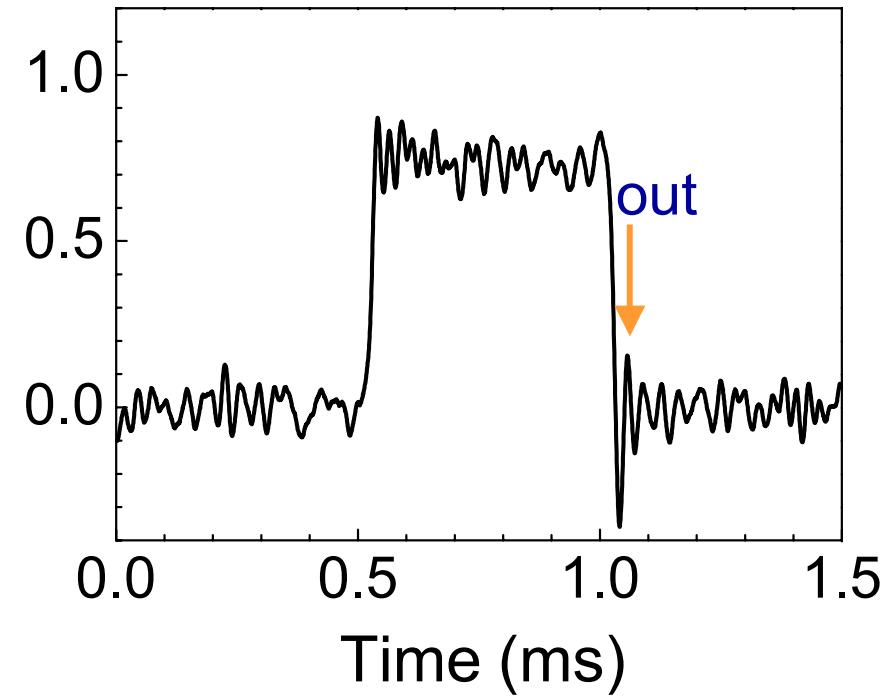
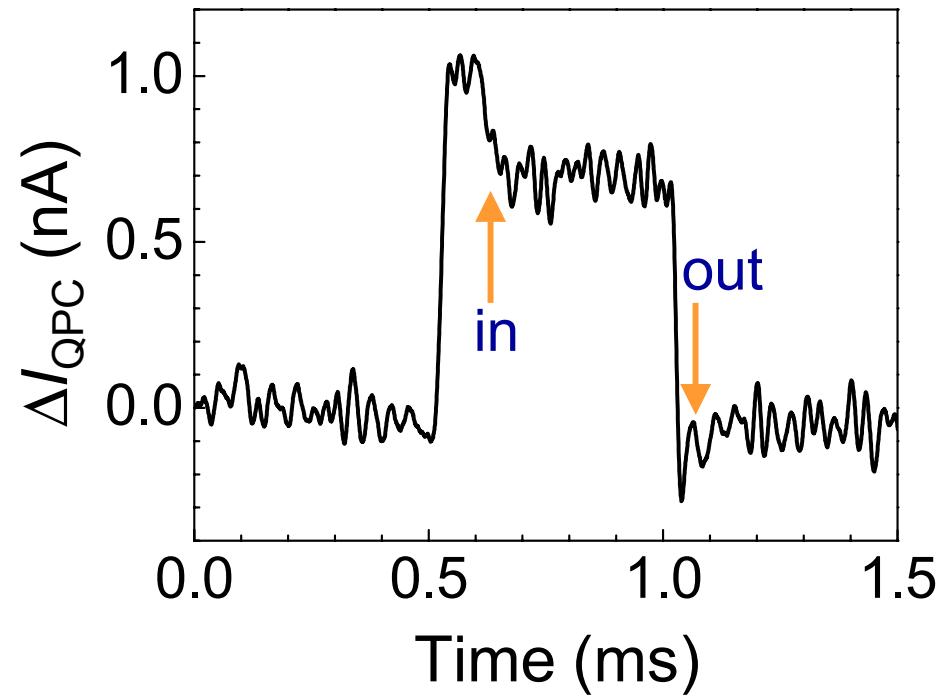
- $V_{\text{SD}} = 1 \text{ mV}$
- $I_{\text{QPC}} \sim 30 \text{ nA}$
- $\Delta I_{\text{QPC}} \sim 1\%$
- speed $\sim 10 \mu\text{s}$
- sensitivity
 $\sim 10^{-3} e (\text{Hz})^{-\frac{1}{2}}$



Tunneling induced by pulse



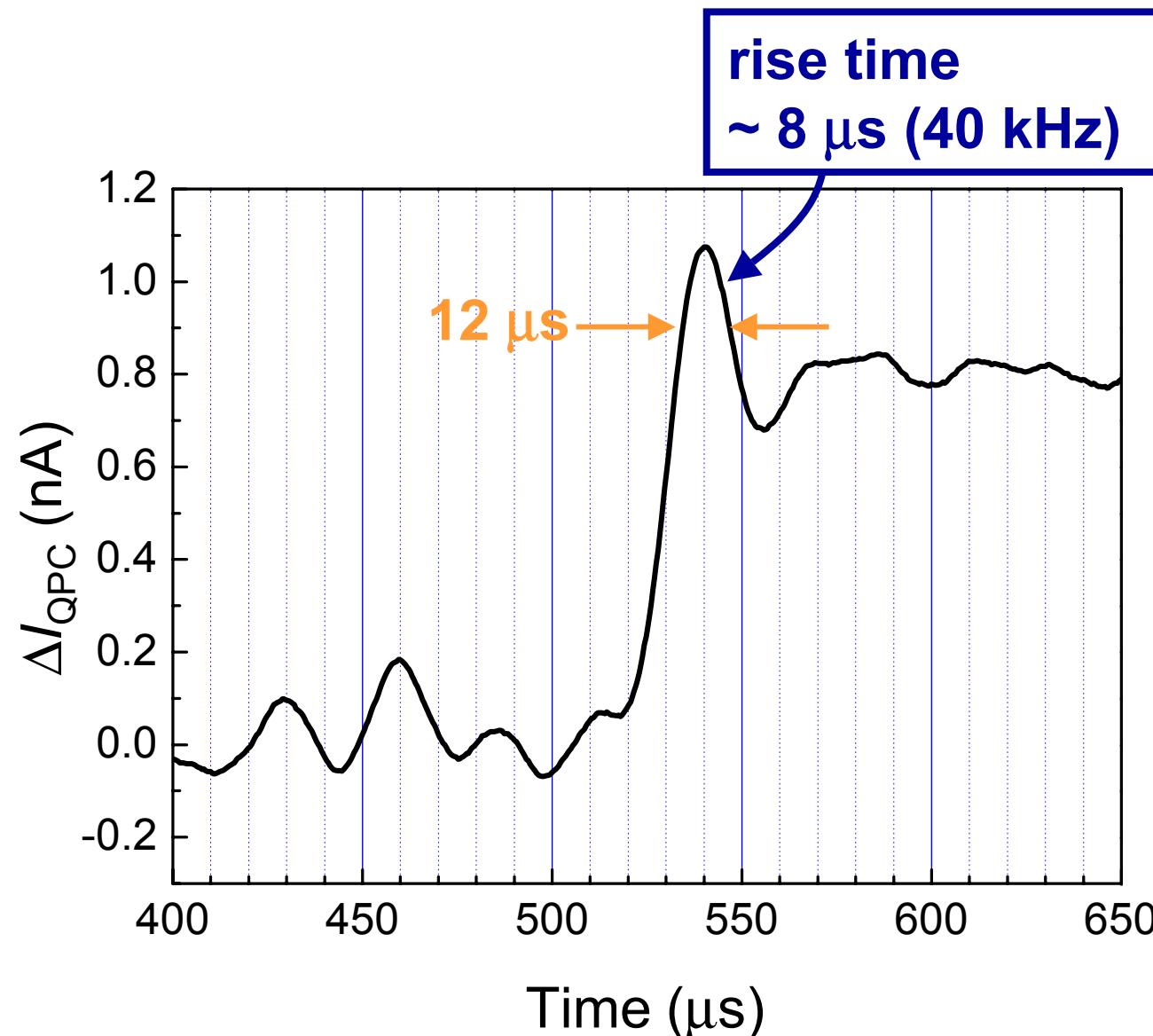
Tunnel-time is stochastic



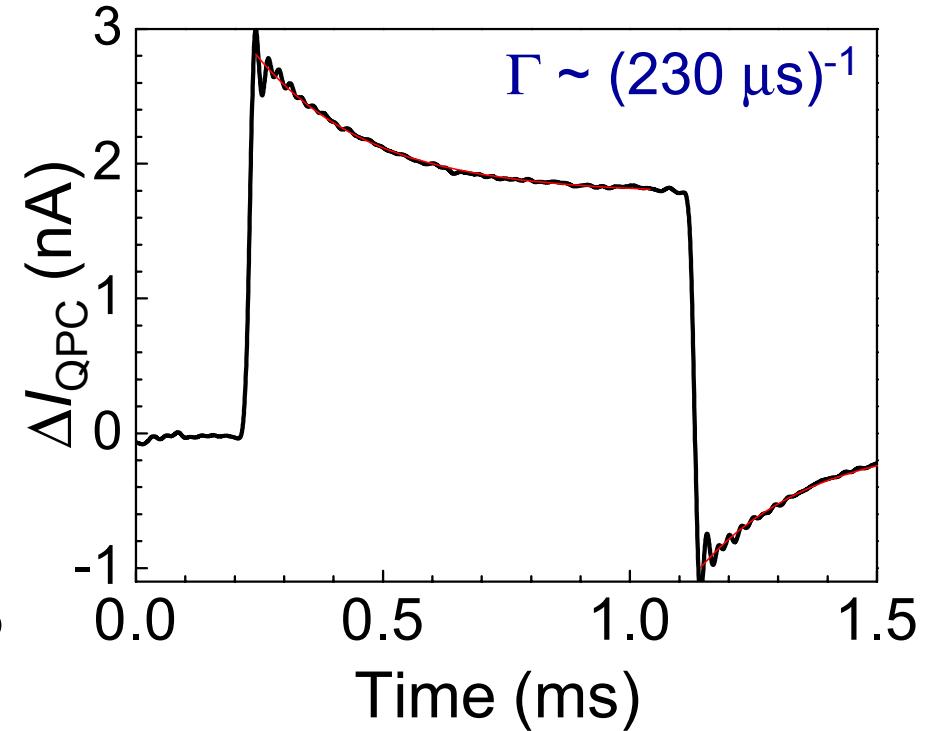
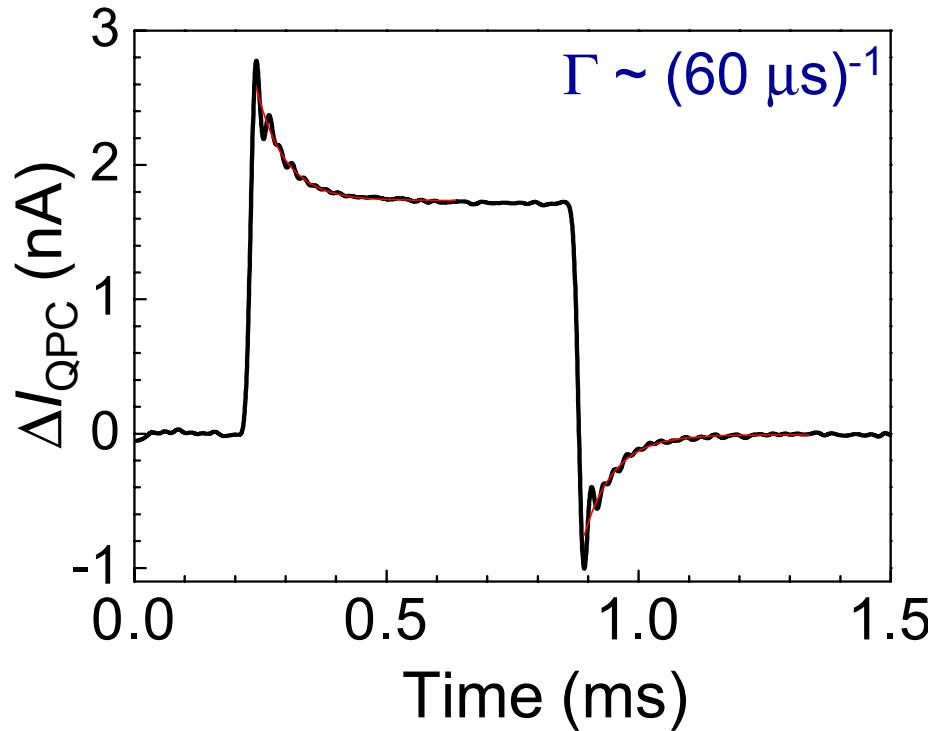
- Tunnel-in event visible
- Tunnel-out event very fast

- Tunnel-in event too fast
- Tunnel-out event visible

Fastest tunnel events



Histograms tunnel time

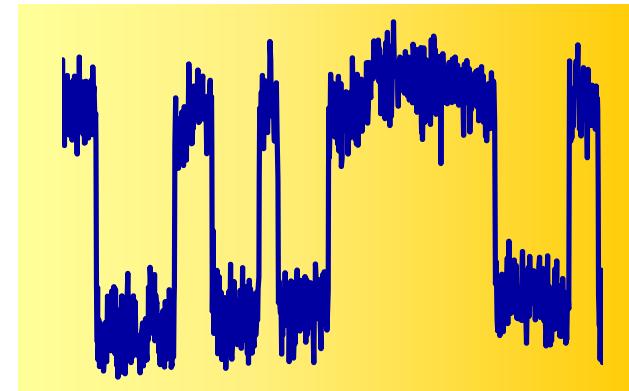


increase tunnel barrier

Summary...

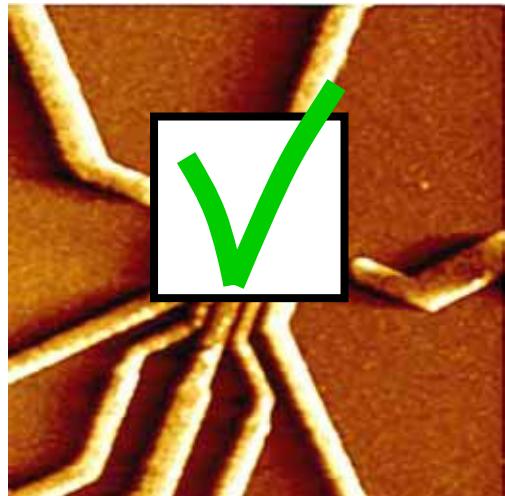
We can:

**measure single-electron
tunneling in *real-time*
(~10 μ s)**

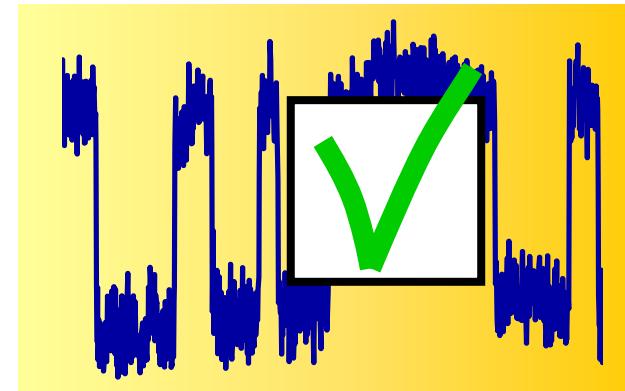


We need...

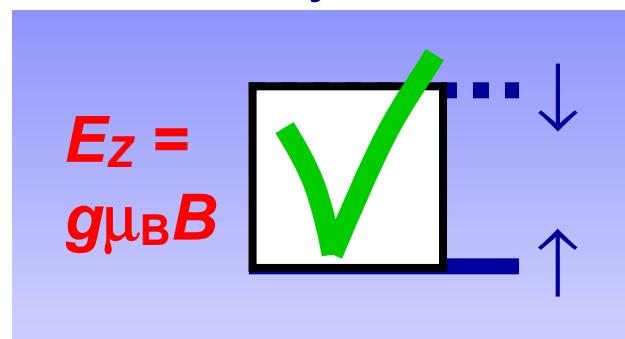
one-electron
double dots...



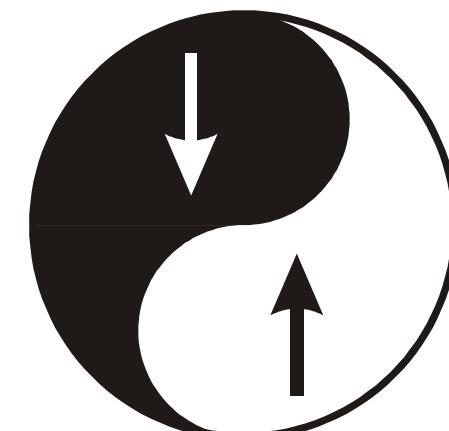
...fast charge
detection...



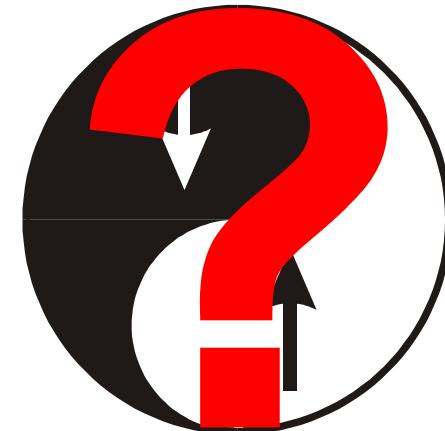
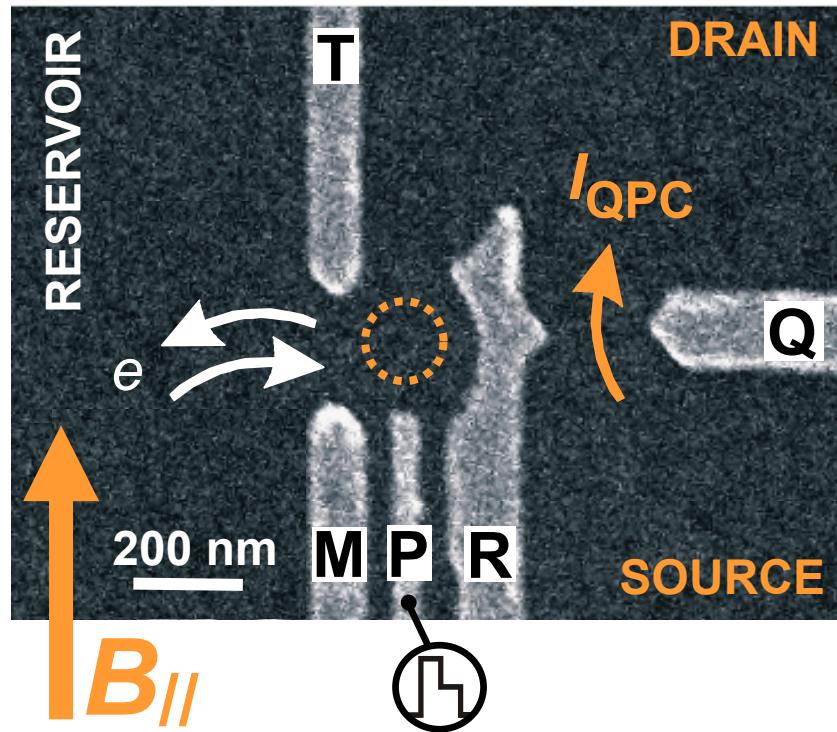
...two-level
system...



....single spin
measurement!

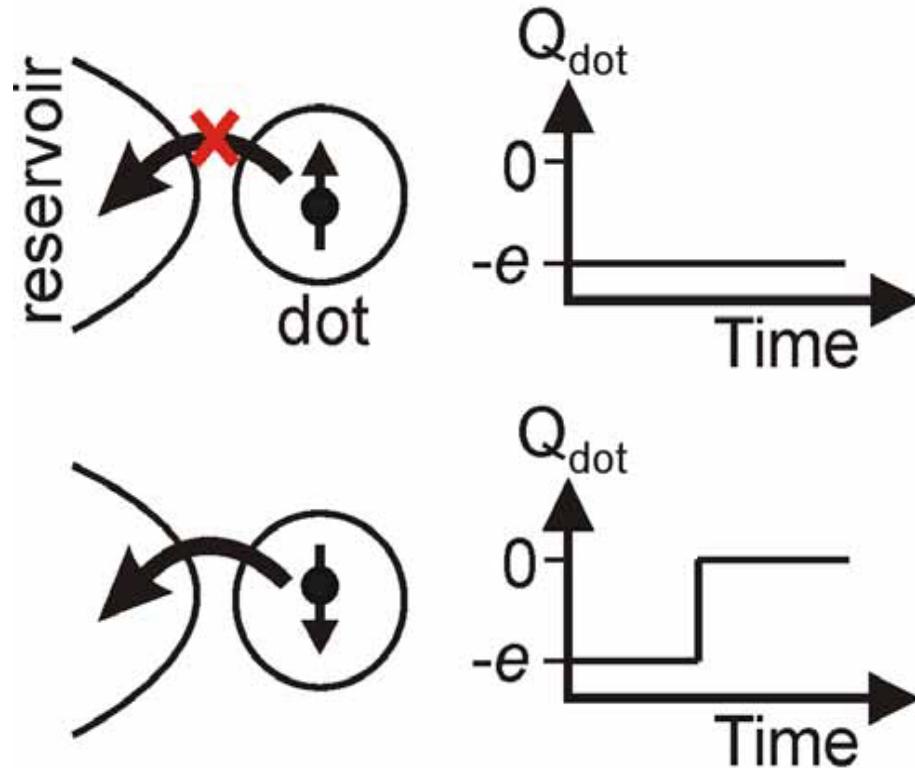


Single-shot measurement of a single electron-spin

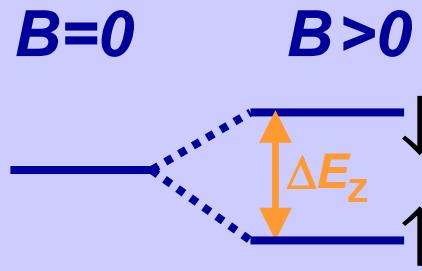


Single-spin measurement concept

- Spin magnetic moment: $\mu_B = 9.27 \times 10^{-23} \text{ A m}^2$ very small!
- But: spin attached to electron (which has charge)
- So: correlate spin orientation to electron's position
- Then measure charge

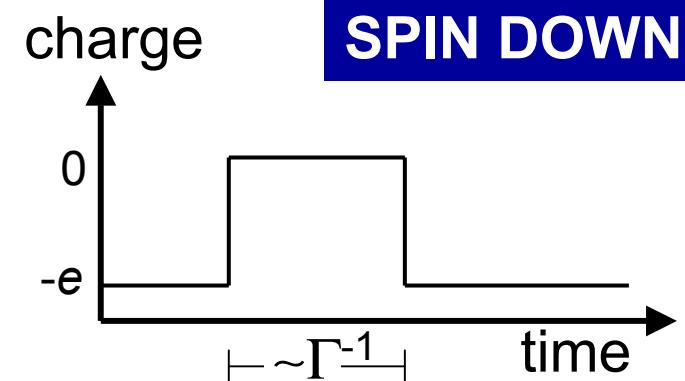
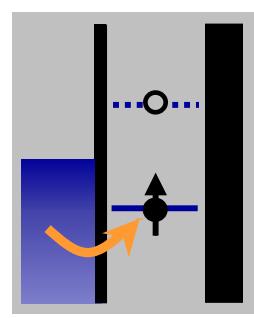
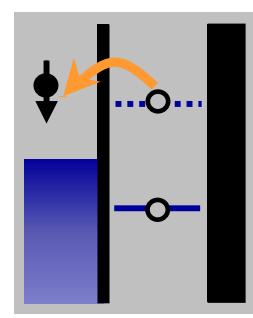
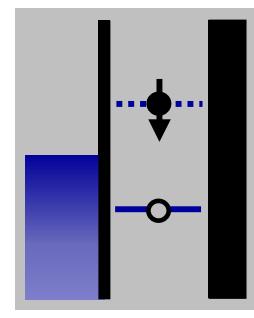
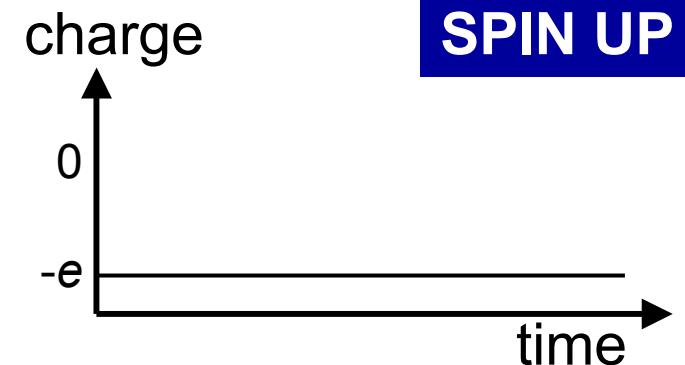
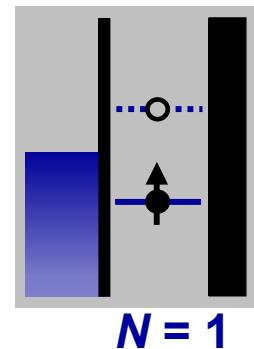


Spin-to-charge conversion

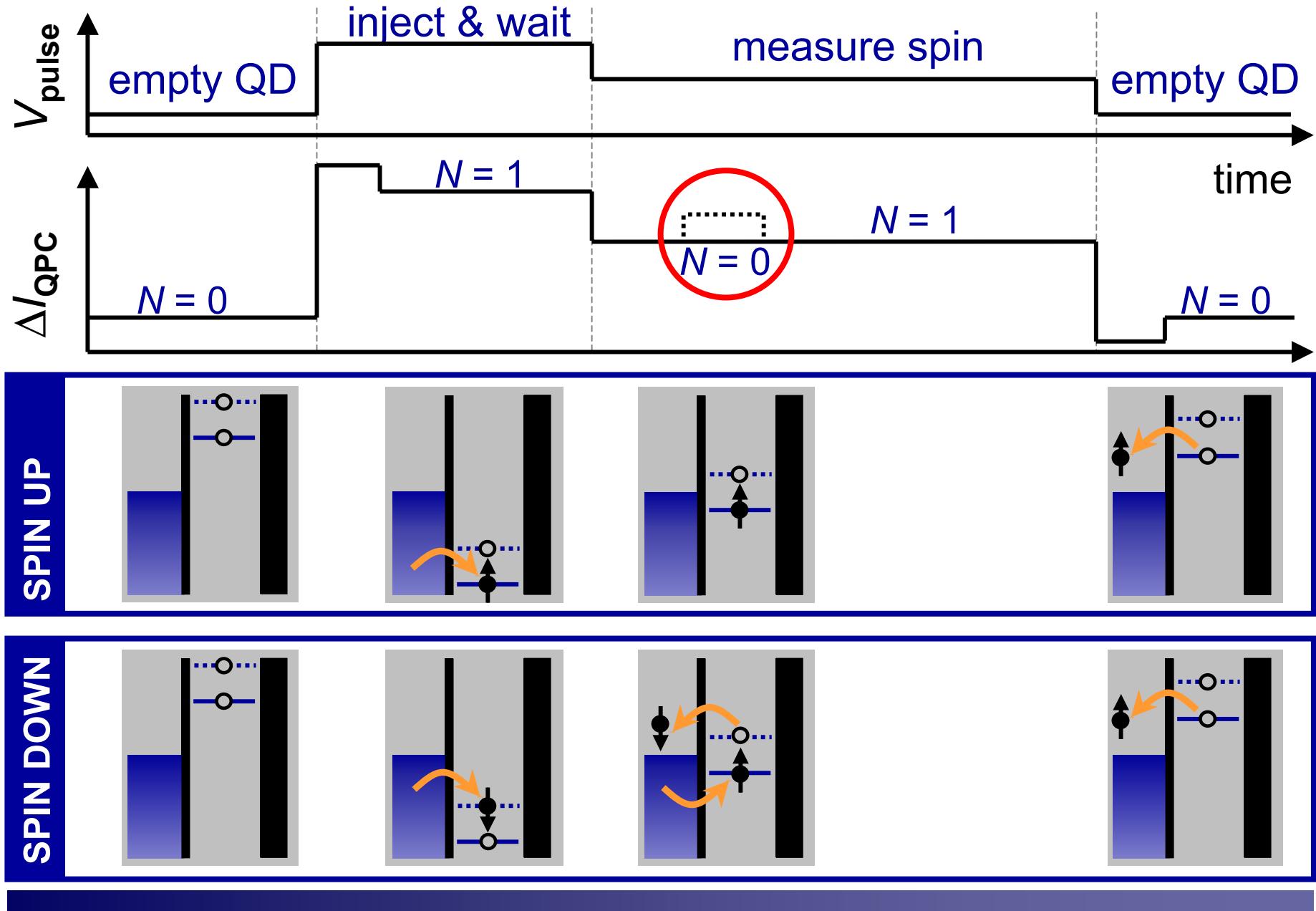


Use Zeeman splitting $\Delta E_z = g\mu_B B$

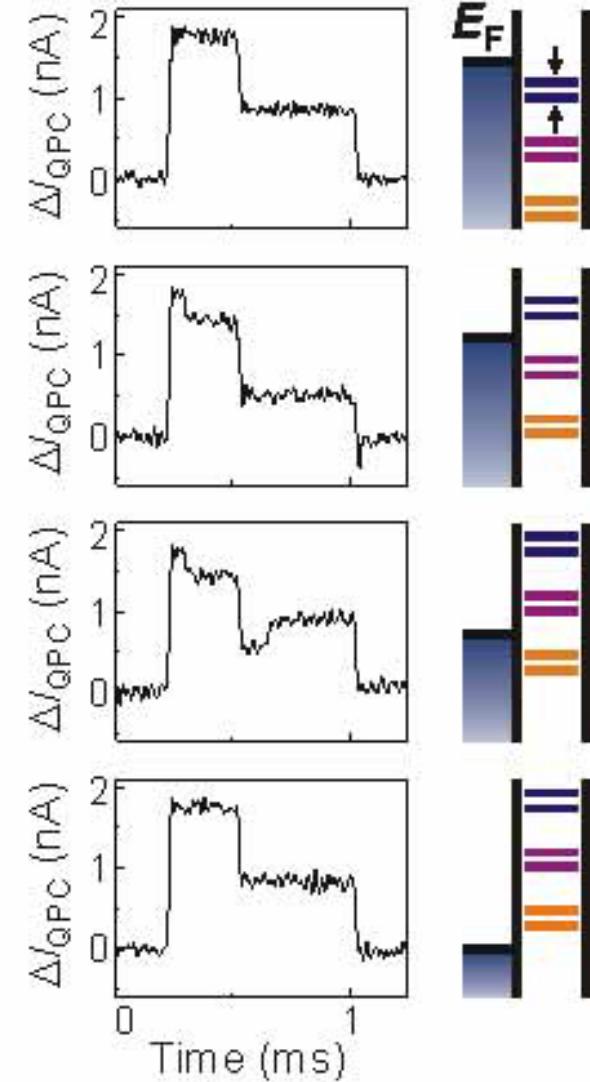
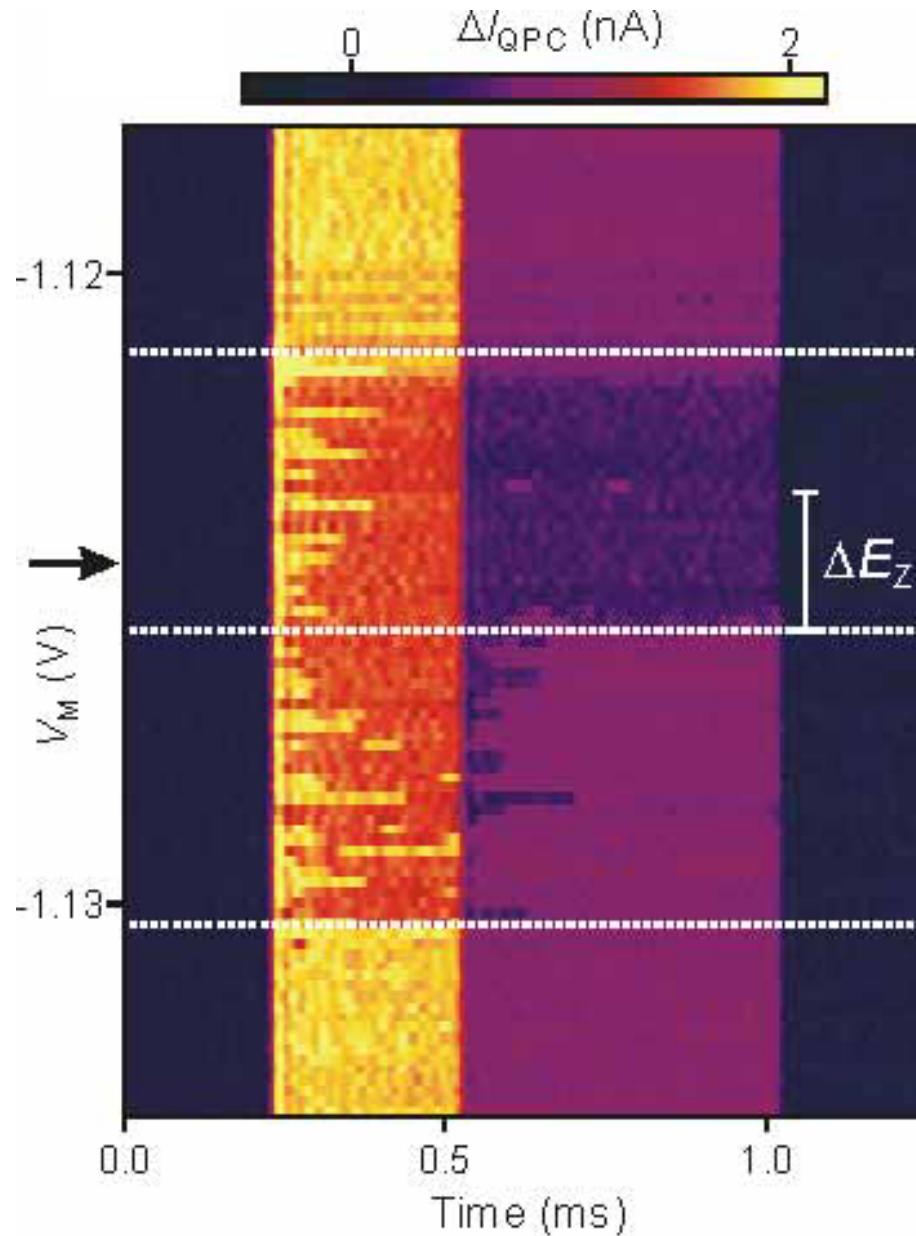
- Spin-to-charge conversion (within $T_1 > 50 \mu\text{s}$)
- Fast charge read-out



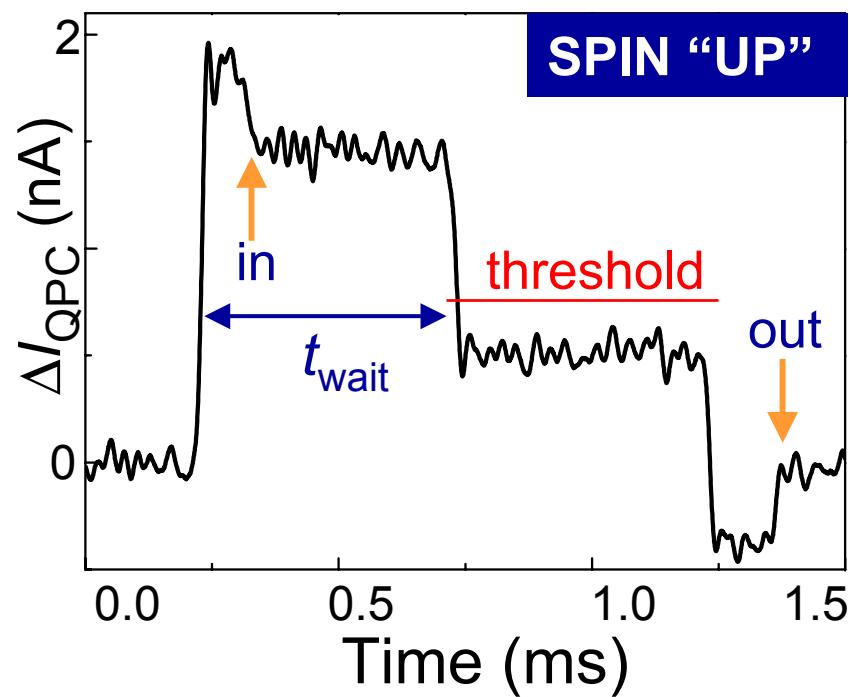
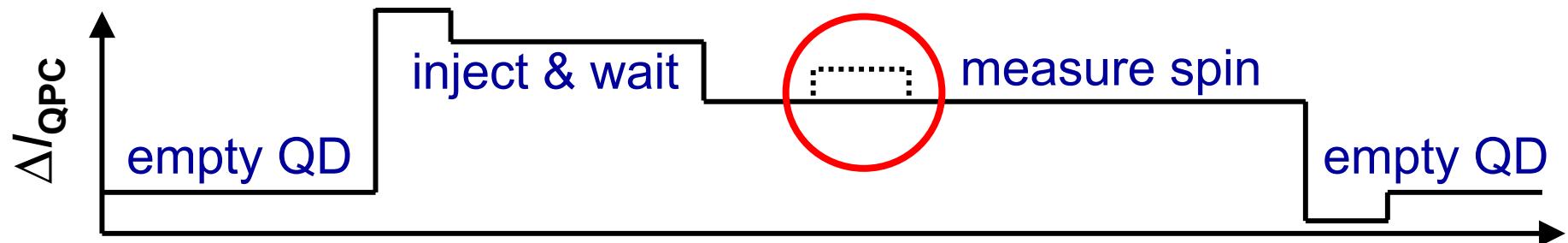
Spin read-out procedure



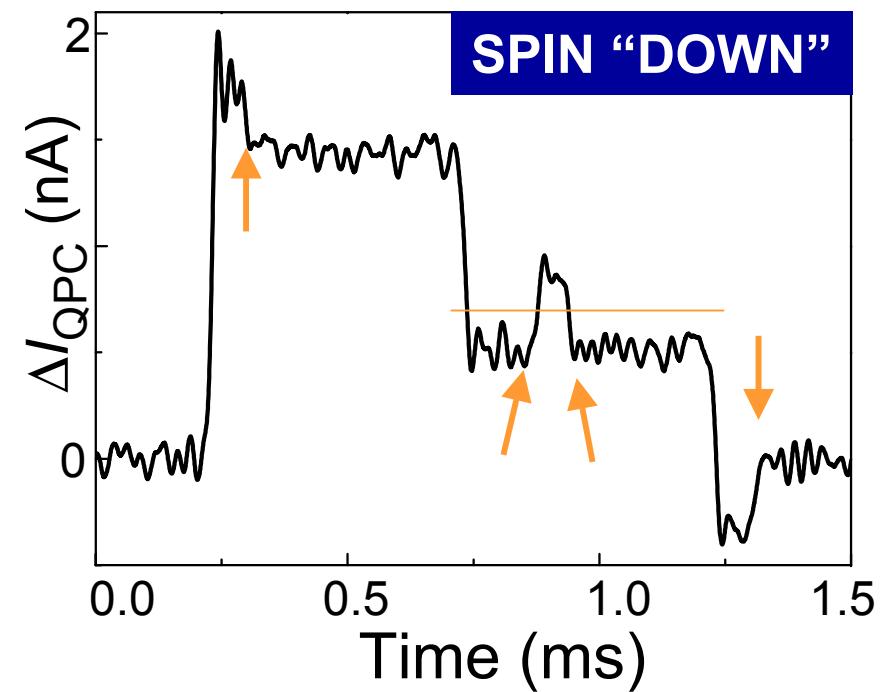
Finding the spin read-out regime



Single-shot spin read-out results

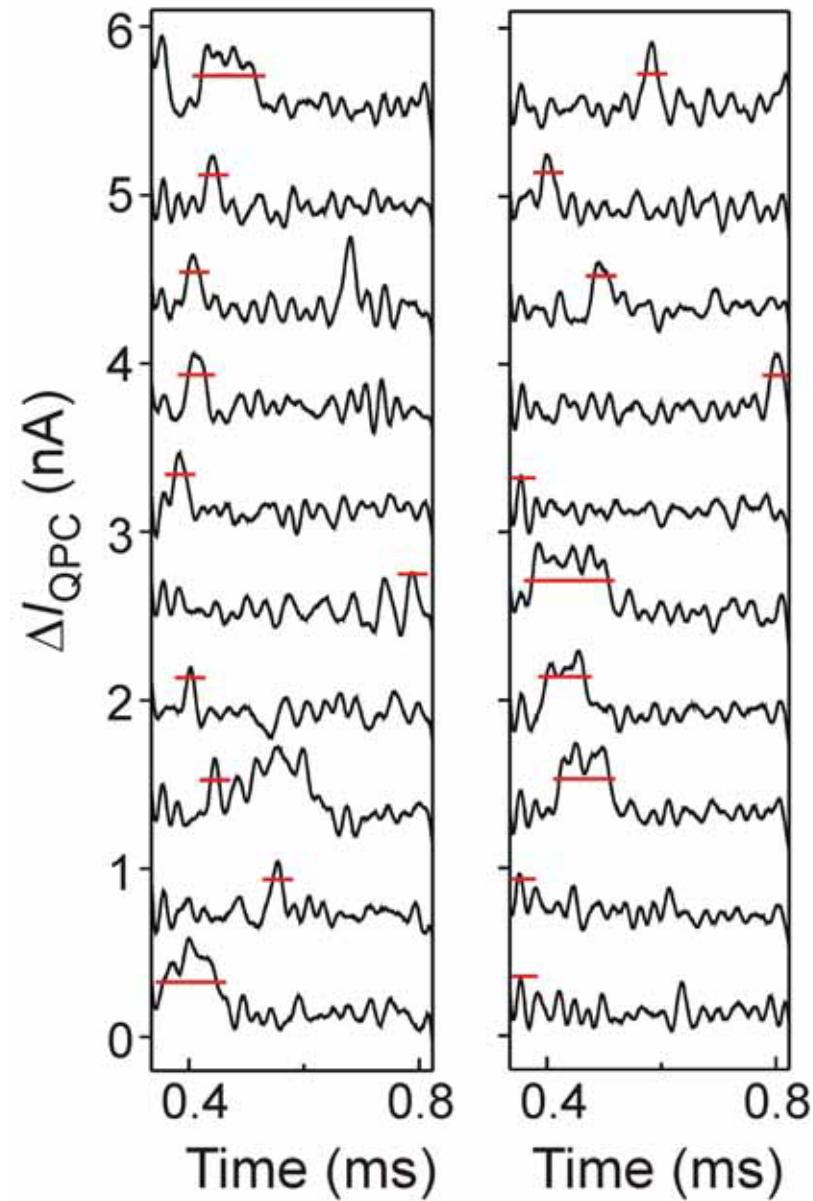
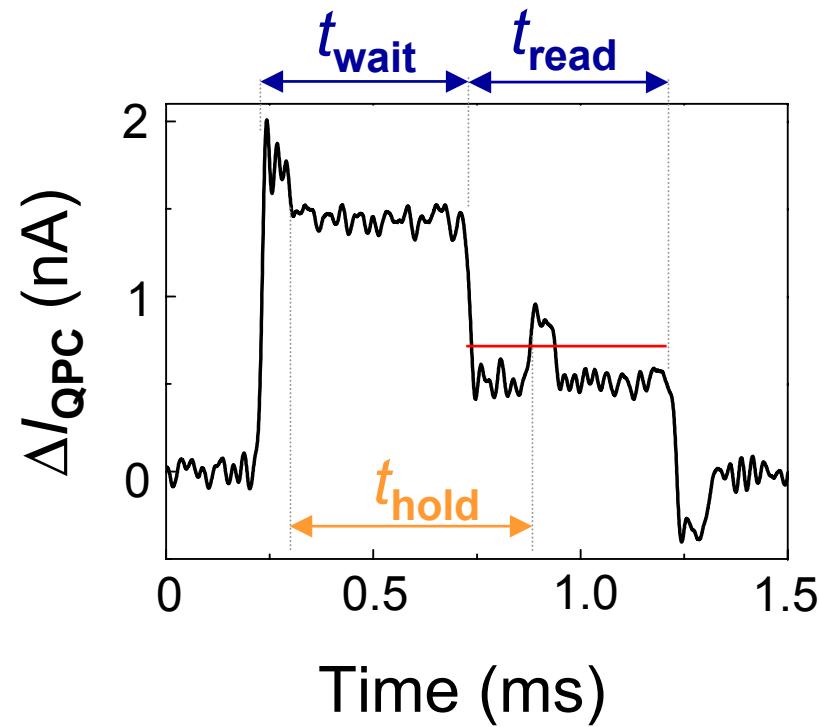


Flat signal in read-out region
♥ spin is “up”

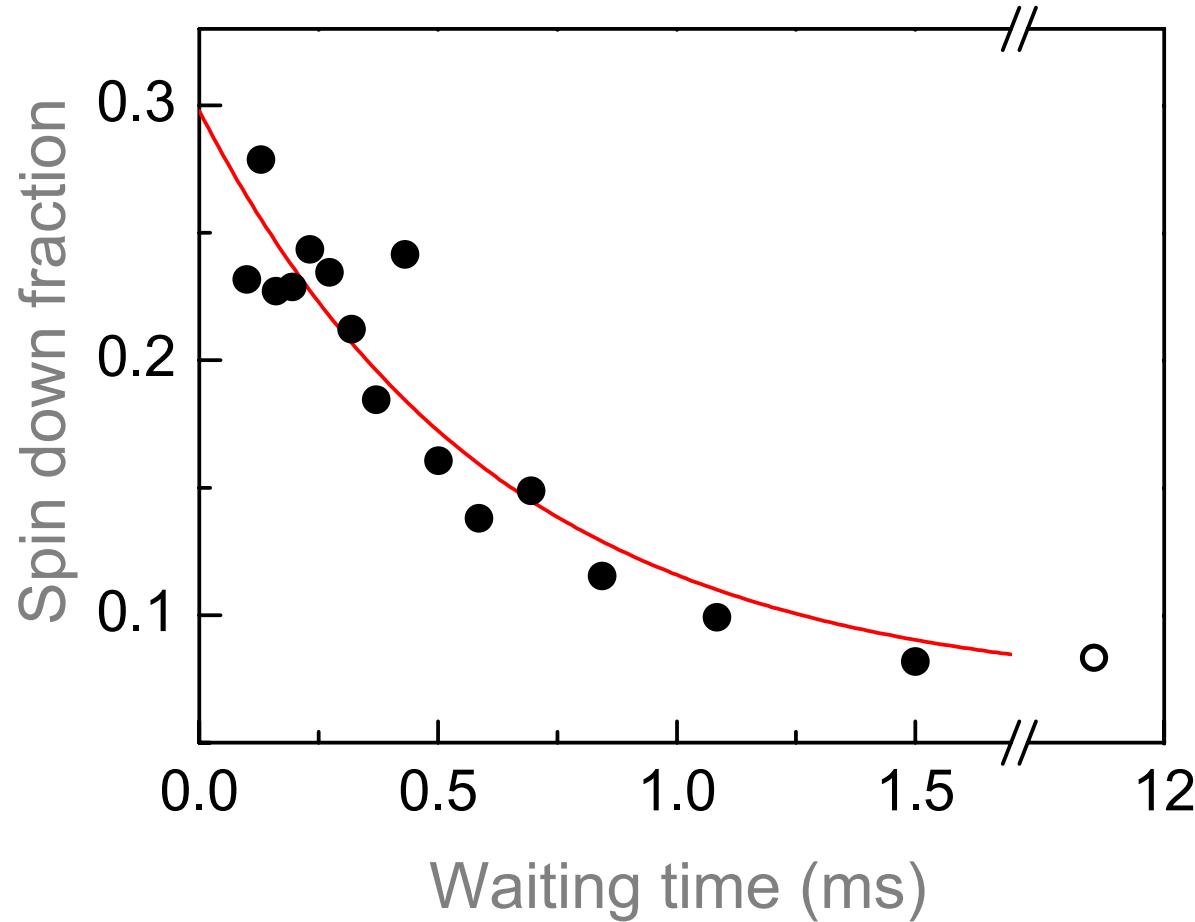


Step during read-out region
♥ spin is “down”

More spin-down traces

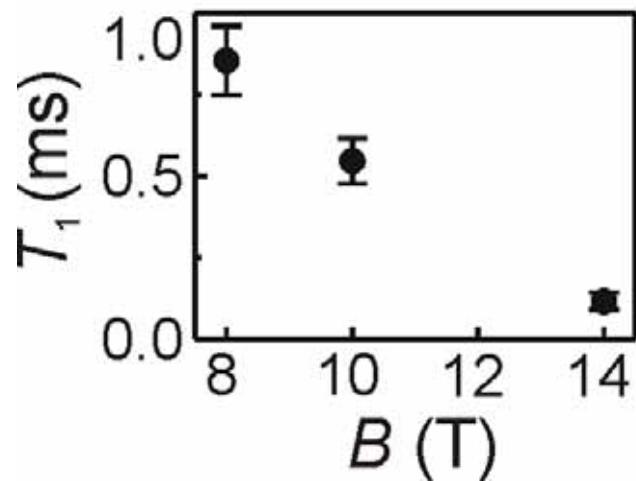


Verification spin read-out

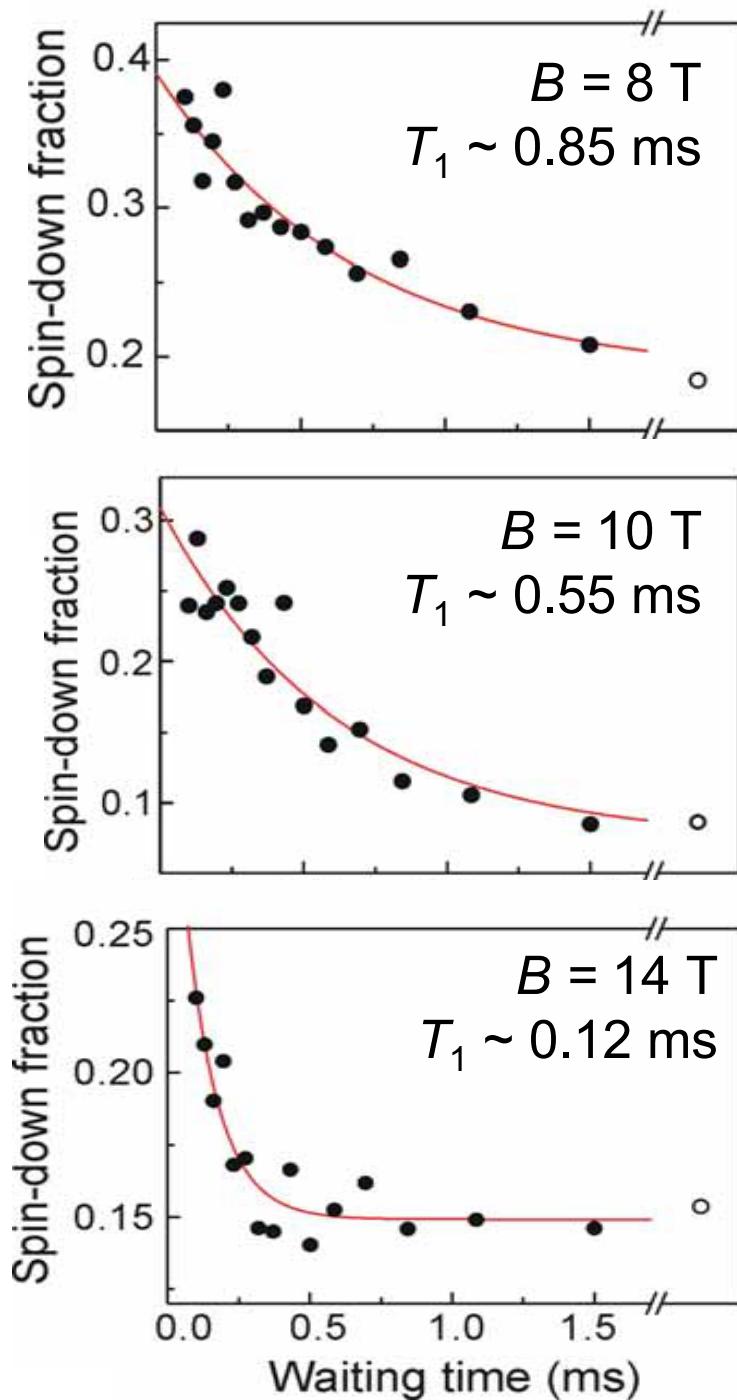


Measurement of T_1

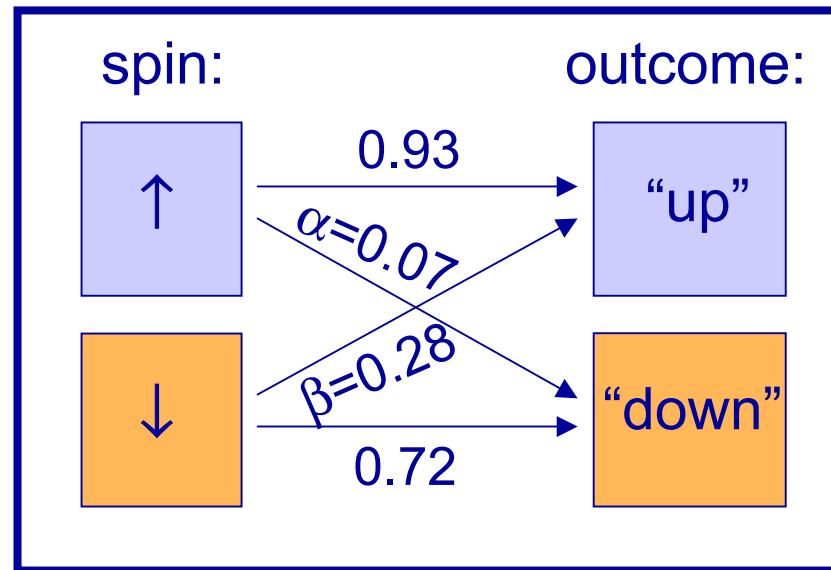
Elzerman et al., to appear in Nature



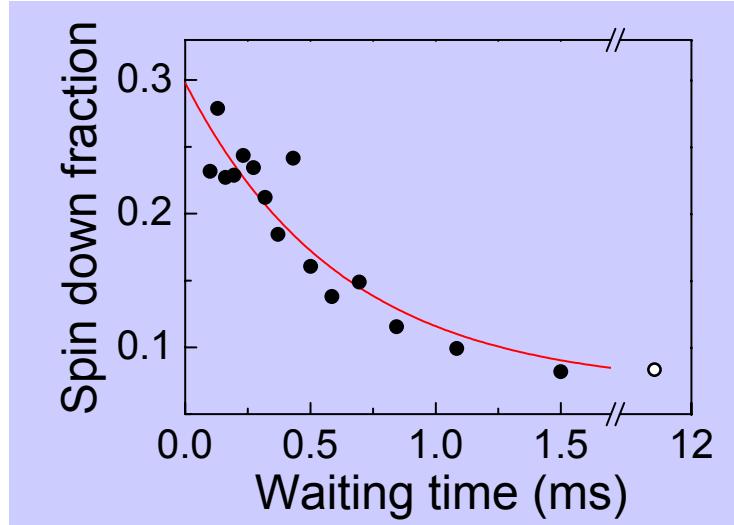
- Surprisingly long T_1
- T_1 goes up at low B



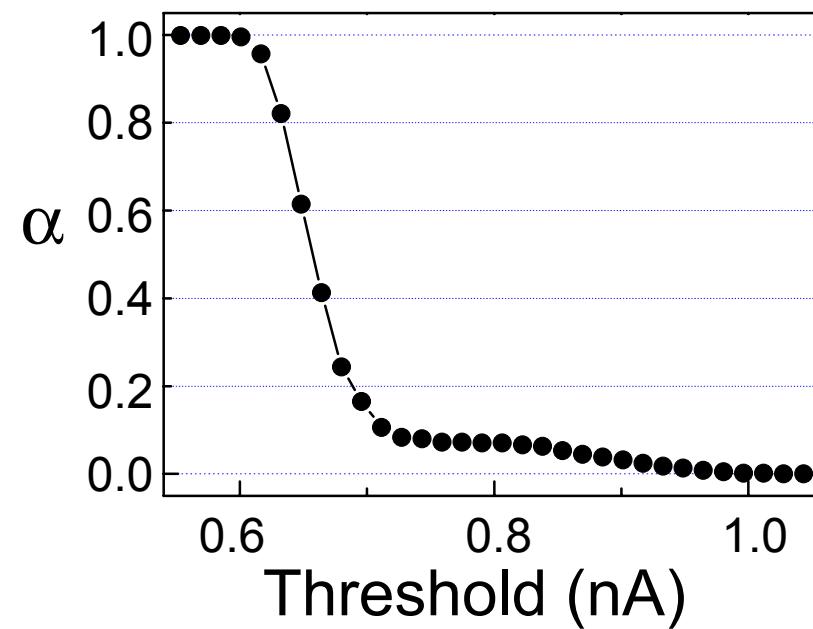
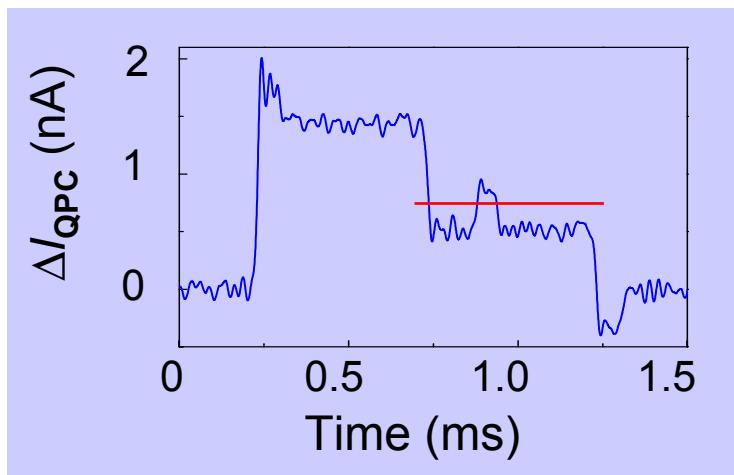
Read-out characterization



Characterization $\alpha = P$ (“down” if \uparrow)



$$p (1 - \beta - \alpha) \exp(-t / T_1) + \alpha$$

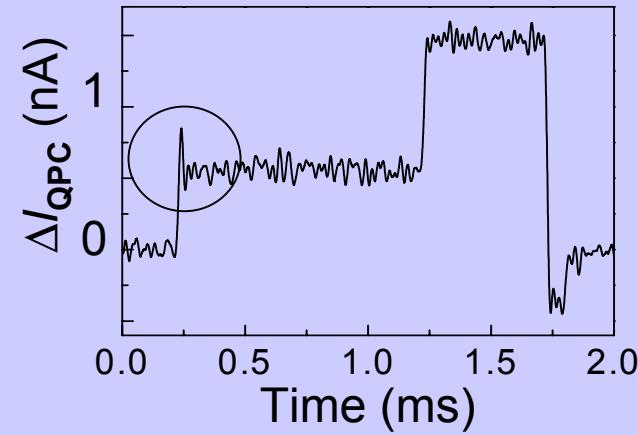
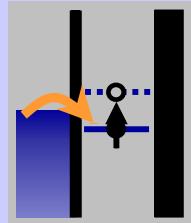


Characterization $\beta = P$ (“up” if \downarrow)

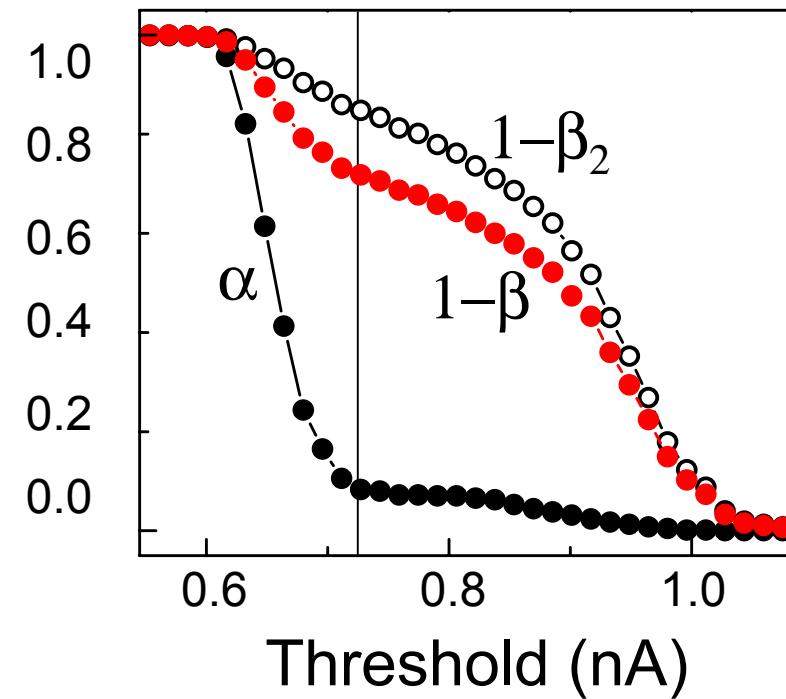
$\beta_1 = P$ (flips before tunnelling)

$$\frac{1/T_1}{1/T_1 + \Gamma_{\downarrow}} = \frac{1}{1 + \Gamma_{\downarrow} T_1}$$

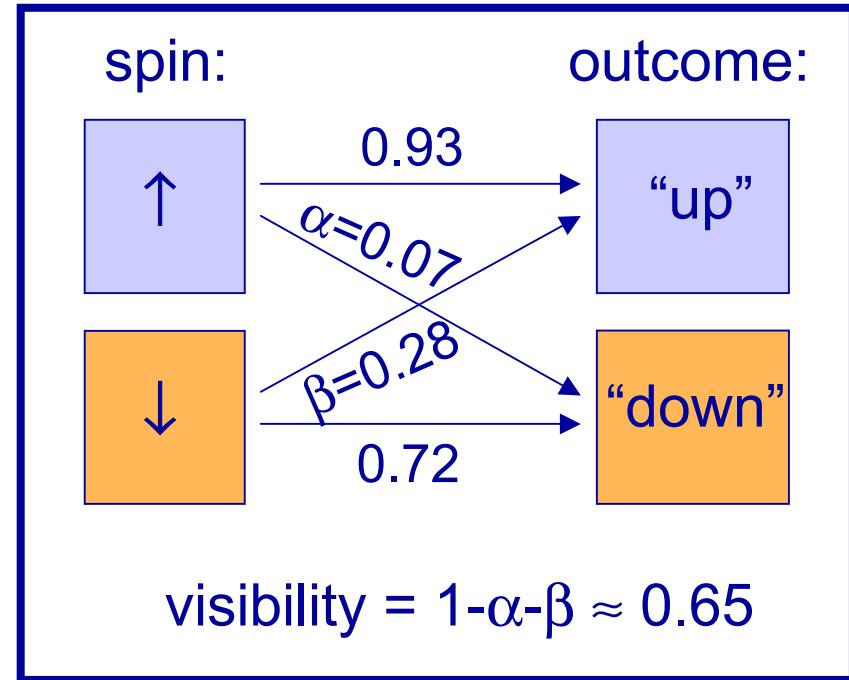
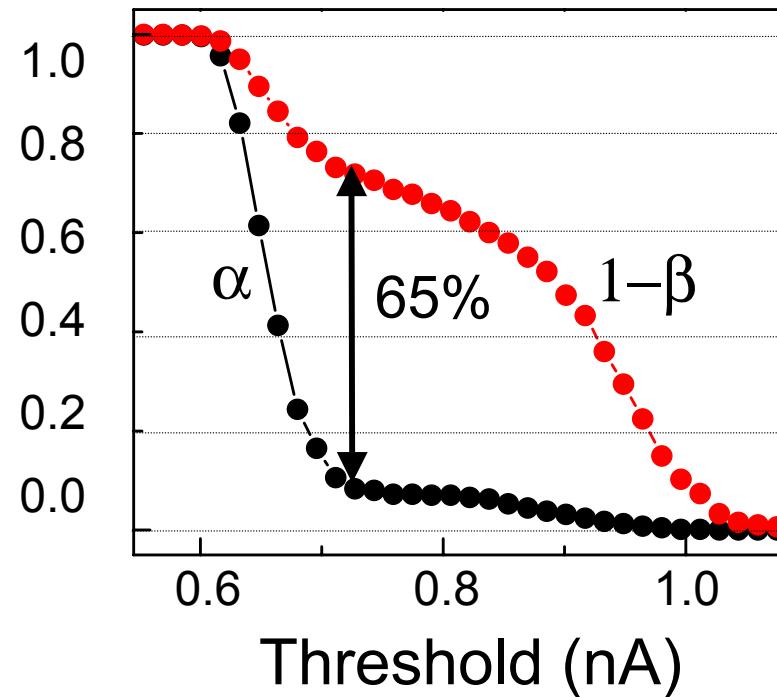
$\beta_2 = P$ (miss step)



$$1-\beta = (1-\beta_1)(1-\beta_2) + \alpha\beta_1$$



Summary single-shot read-out fidelity



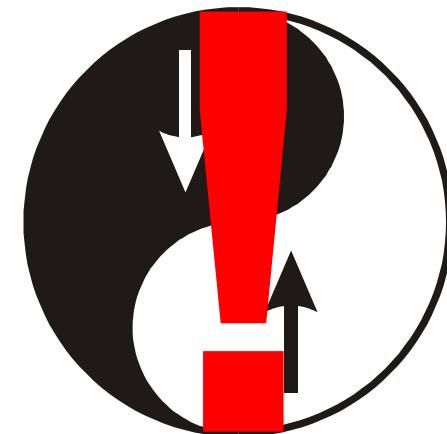
Future improvements:

- α : lower T_{el}
- β : faster charge detection

Summary...

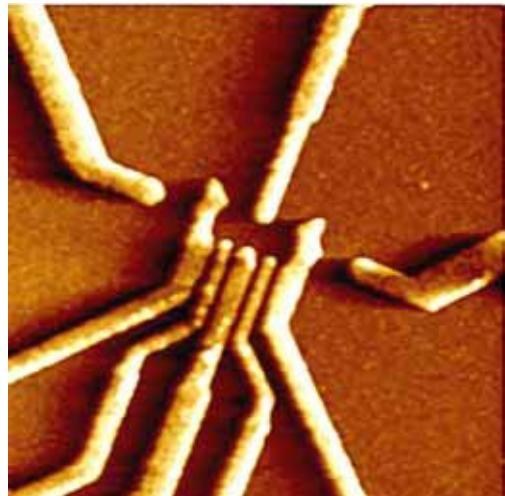
We can:

**perform single-shot
measurement of one
electron-spin**

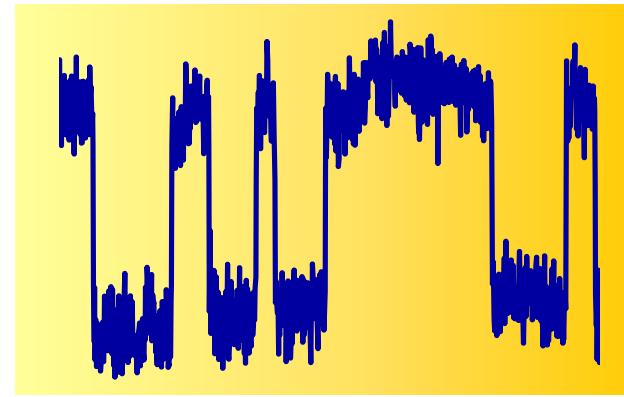


We have...

one-electron
double dots...



...fast charge
detection...

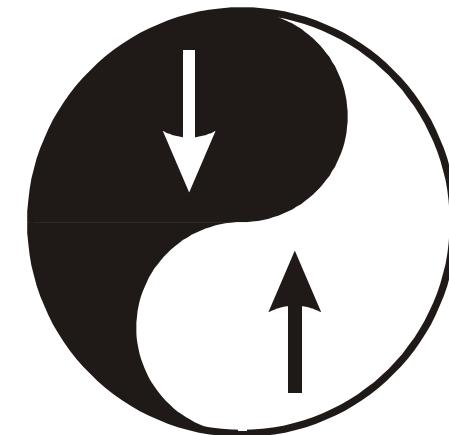


...two-level
system...

$$E_z = g\mu_B B$$

A diagram illustrating a two-level system. It shows two horizontal dashed lines representing energy levels. A red arrow labeled $E_z = g\mu_B B$ points downwards between the levels. Below each level is a horizontal bar with a black arrow pointing upwards, representing spin states.

....single spin
measurement!



Part IV: Outlook

Coherent spin manipulation (ESR)

Two-qubit operations ($\text{SWAP}^{\frac{1}{2}}$)

Bell's inequalities for massive particles

Quantum simulation?

Quantum computation??