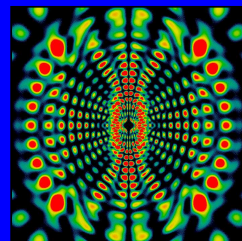


Non-monotonic quantum-to-classical transition in multi-particle interference

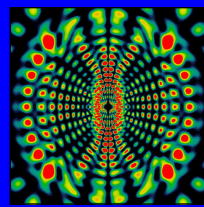
Andreas Buchleitner

Quantum optics and statistics

Institute of Physics, Albert Ludwigs University of Freiburg



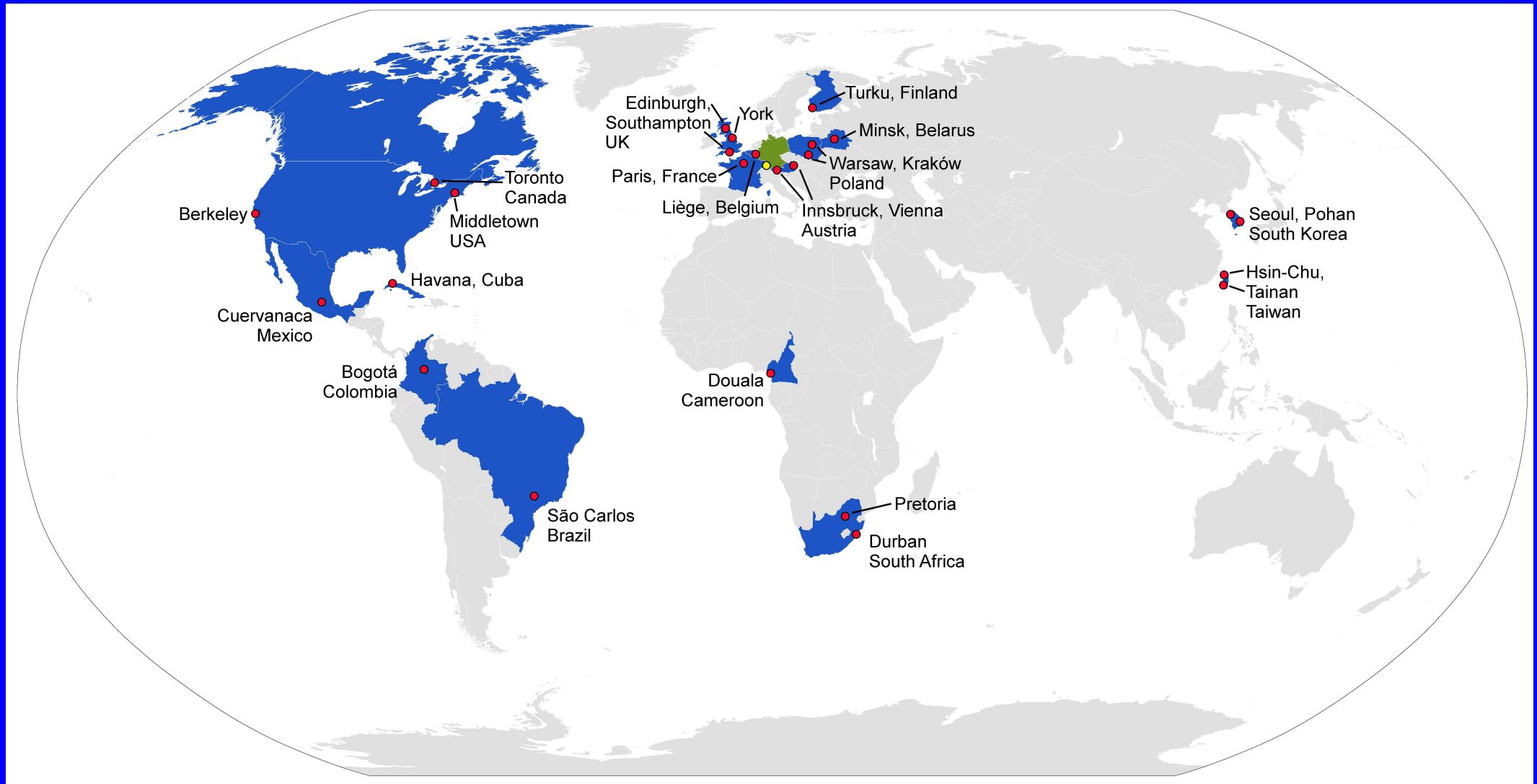
**Nonlinear Dynamics
in Quantum Systems**



Nonlinear Dynamics in Quantum Systems



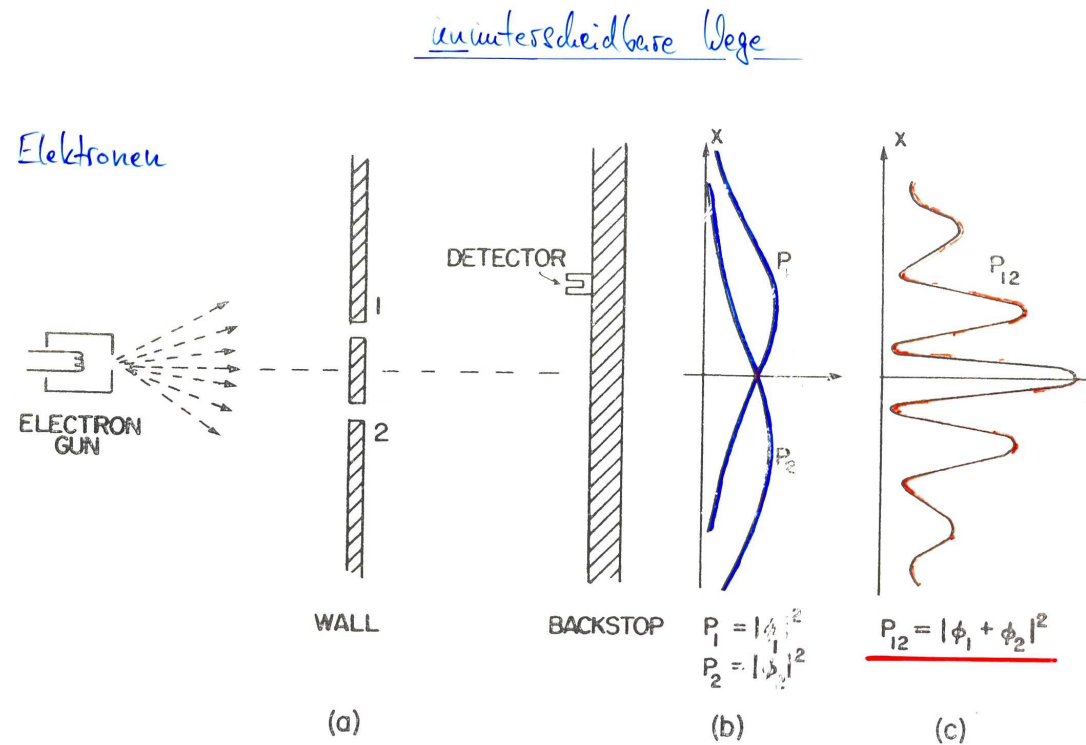
In Collaboration with . . .



. . . M. Tichy, F. Mintert, Y.-S. Ra, H.-T. Lim, O. Kwon, Y.-H. Kim

Single particle quantum interference

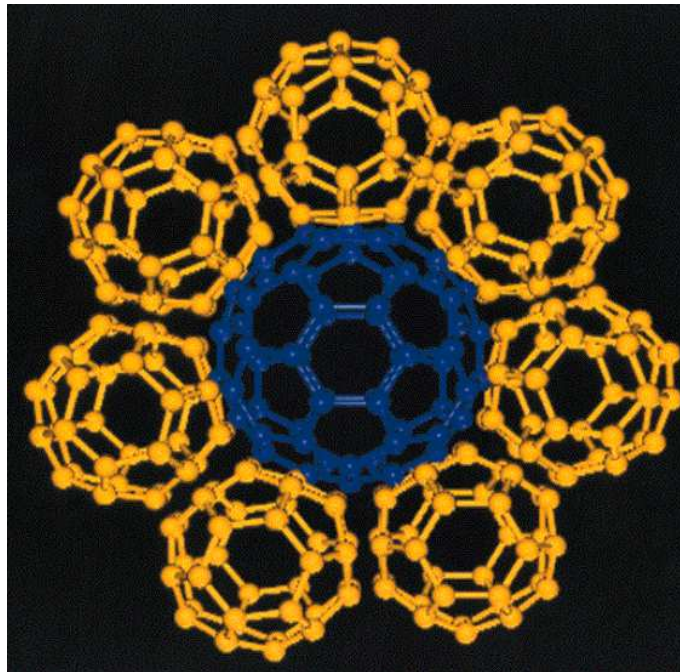
- Coherent superposition of single particle amplitudes



[Feynman, Lecture Notes on Quantum Mechanics]

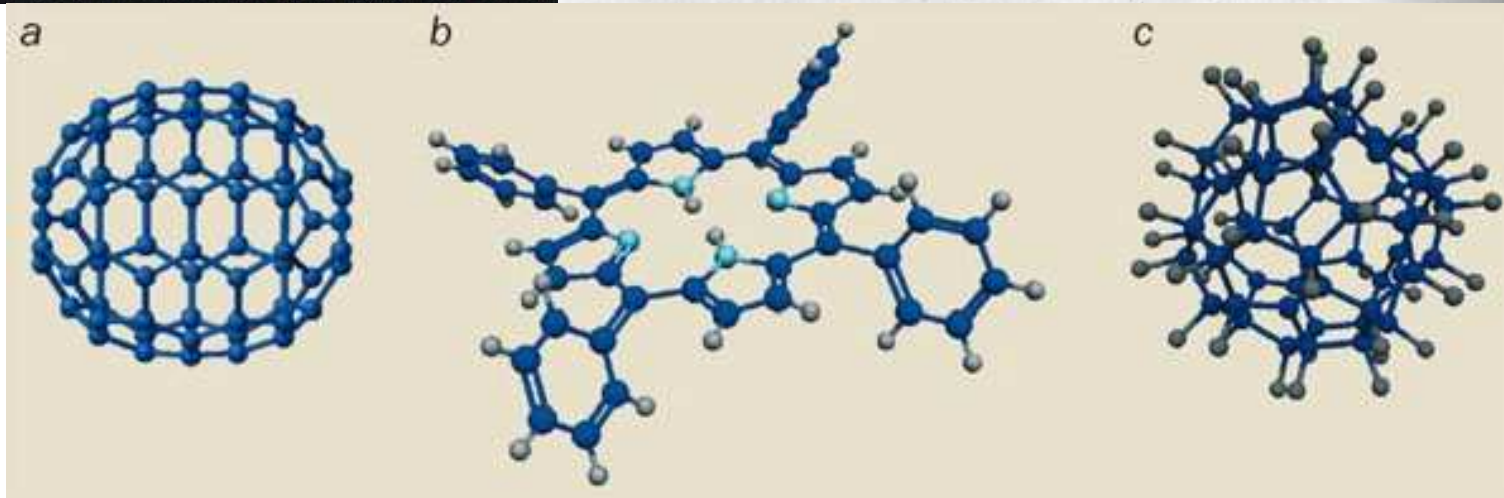
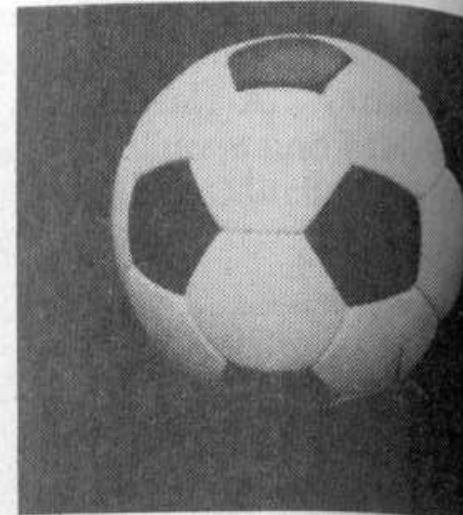
- Which-way information induces decoherence

Also true for significantly larger objects



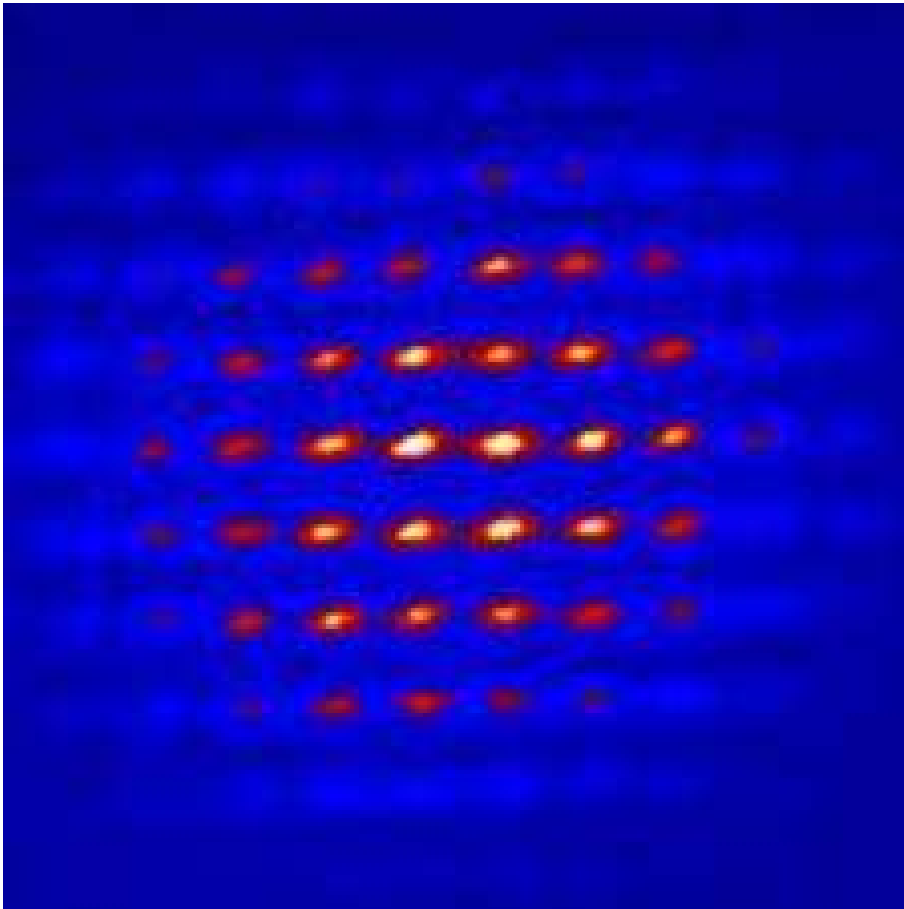
NATURE

Fig. 1 A football (in the United States, a soccerball) on Texas grass. The C_{60} molecule featured in this letter is suggested to have the truncated icosahedral structure formed by replacing each vertex on the seams of such a ball by a carbon atom.



What about many-particle systems, e.g. (but not only) quantum computers?

N cold atoms in a matrix [Birkel et al.]



many particle quantum register

$|011001011101010001\rangle$

rely on coherent
superposition/dynamics of 2^N many
particle states

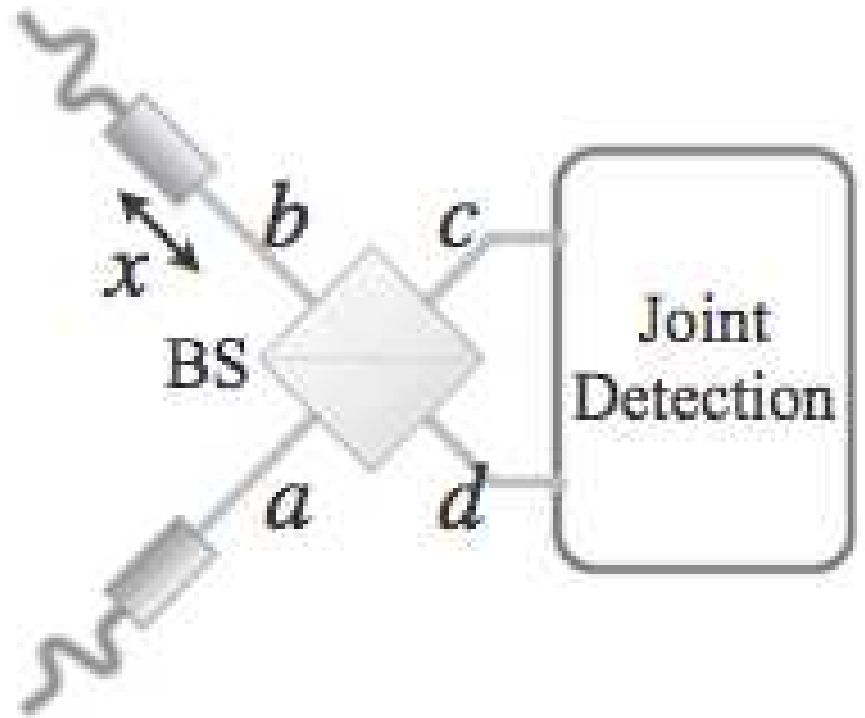
$a|011001011101010001\rangle$
 $+b|001001011101010101\rangle$
 $+c|011001001101110011\rangle$
 $+d|111001010001010001\rangle$
 $+ \dots$

what about (in-)distinguishability and/or (de-)coherence?

The Hong-Ou-Mandel effect

one photon in each mode a and b – distinguishability controlled by path delay x

coincident detection in output modes c and d



- coincidence probability if **distinguishable**: $P(2; 1, 1) = 1/2$
- coincidence probability if **indistinguishable**: $P(2; 1, 1) = 0$

if indistinguishable: destructive interference of two two-particle trajectories

[Shi & Alley (1986, 1988); Hong, Ou & Mandel (1987)]

What happens “in between”?

Continuous distinguishability-transition

inject (Gaussian) photon wave packet state $|1_{t_1}\rangle$ ($|1_{t_2}\rangle$) with beam splitter arrival time t_1 (t_2) in mode a (b) – then (Gram-Schmidt):

$$|1_{t_2}\rangle = \alpha|1_{t_1}\rangle + \sqrt{1 - \alpha^2}|\tilde{1}_{t_1}\rangle, \quad \alpha^2 = \exp\left(-(\Delta\omega)^2(t_2 - t_1)^2/2\right)$$

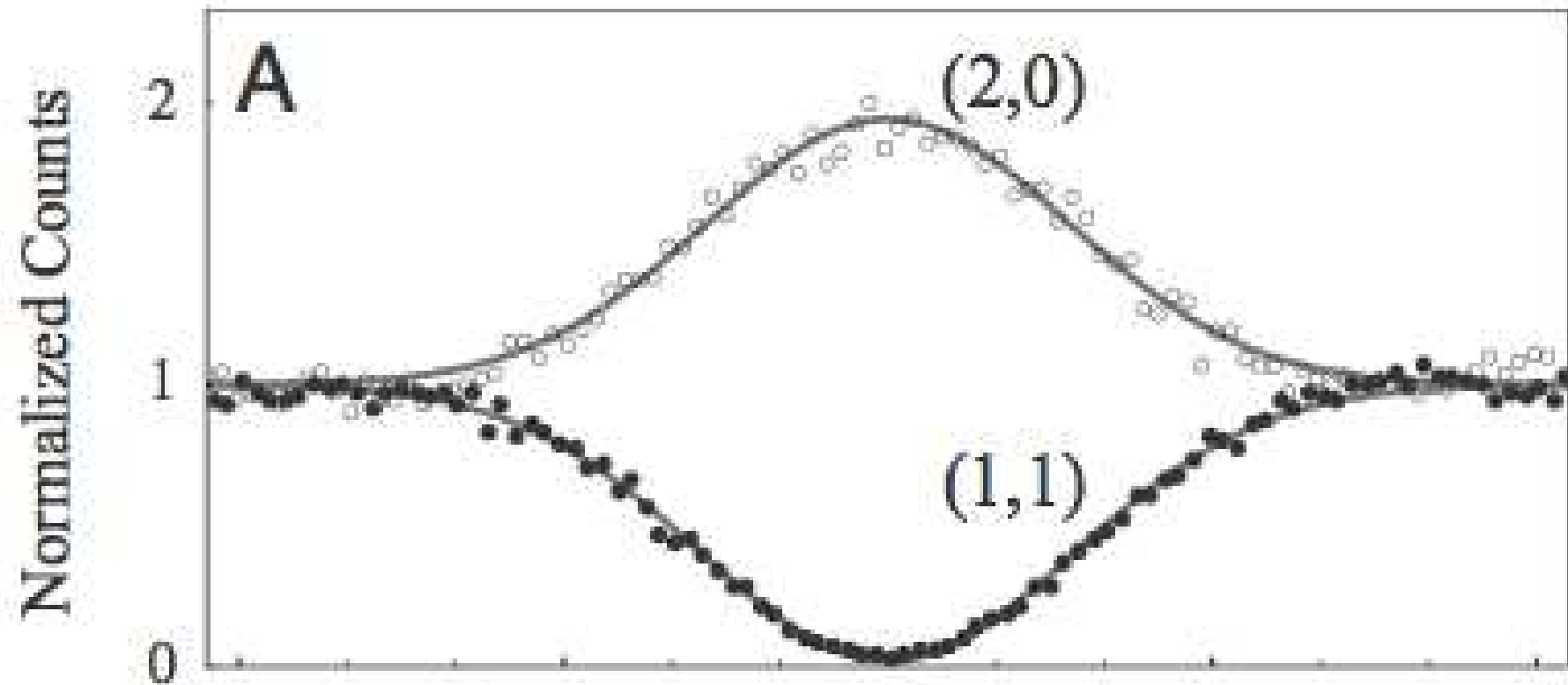
HOM two-photon state impinging on beam splitter:

$$a_{t_1}^\dagger b_{t_2}^\dagger |0\rangle = \alpha|1, 1\rangle + \sqrt{1 - \alpha^2}|1, \tilde{1}\rangle,$$

with the photons in $|1, 1\rangle$ ($|1, \tilde{1}\rangle$) in(-distinguishable) – $\alpha \in \{0, 1\}$ quantifies distinguishability (fully indistinguishable for $\alpha = 1$).

Ergo: **Monotonic** disappearance of interference signal
in coincident detection as α decreases!

Experimental confirmation



The more two-particle which-way information, the less interference
– as for single-particle scenario!

Something new happens for more particles!

Inject **two photons per mode**, and write in Gram-Schmidt decomposition with respect to temporal overlap:

$$\frac{1}{2}(a_{t_1}^\dagger)^2(b_{t_2}^\dagger)^2|0\rangle = \alpha^2|2, 2\rangle + \sqrt{2}\alpha\sqrt{1 - \alpha^2}|2, 1, \tilde{1}\rangle + (1 - \alpha^2)|2, \tilde{2}\rangle ,$$

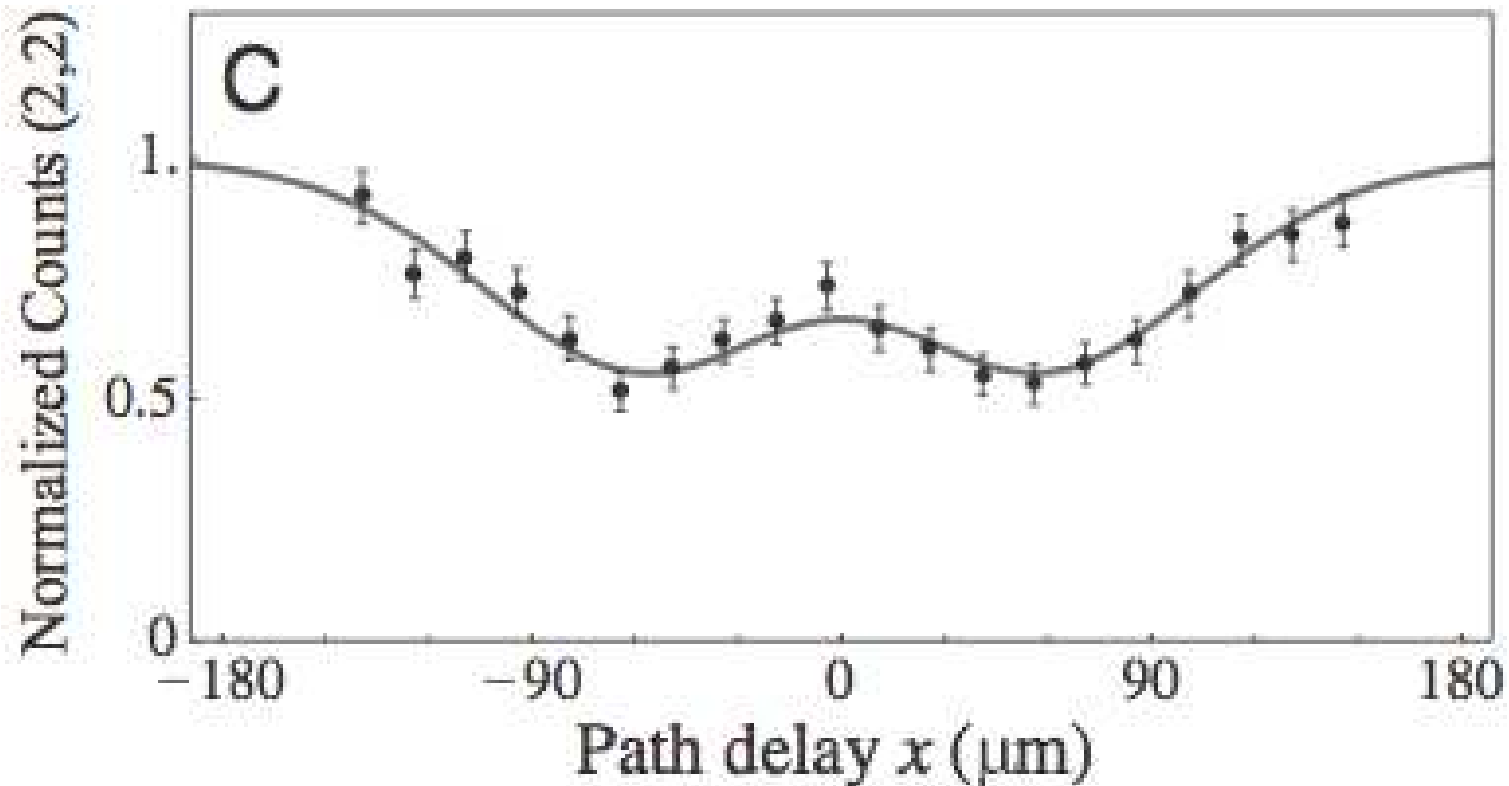
For more than one photons per mode, there **appear partially distinguishable contributions** – here, $|2, 1, \tilde{1}\rangle$ implies the **interference of three-particle rather than four-particle paths**, with weight $W_{\text{type}}^{(N;m,n)}$ non-monotonic in α !

General expression for (m,n)-event probability:

$$P(N; m, n) = \sum_{\text{type}} p_{\text{type}}^{(N;m,n)} W_{\text{type}}^{(N;m,n)}$$

“type” – fully indistinguishable, fully distinguishable, partially distinguishable
 p_{type} determined by mode mapping $a^\dagger \rightarrow (c^\dagger + d^\dagger)/\sqrt{2}$, $b^\dagger \rightarrow (c^\dagger - d^\dagger)/\sqrt{2}$.

Non-monotonic quantum-to-classical transition



With increasing distinguishability of the particles, the order of many-particle interferences is reduced from four- over three- to two-particle paths, with non-monotonic weights! Gaining which-path information at different multi-particle orders generically leads to a *non-monotonic* quantum-to-classical transition!

Conclusion/Perspective

- The distinguishability transition affects different orders of many-particle interferences differently!
- Quantum interference in correlation functions thus fades away non-monotonically, in general!
- A probe for different types of system-environment interactions?

– Phys. Rev. A 83, 062111 (2011); PNAS 110, 1227 (2013) –
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