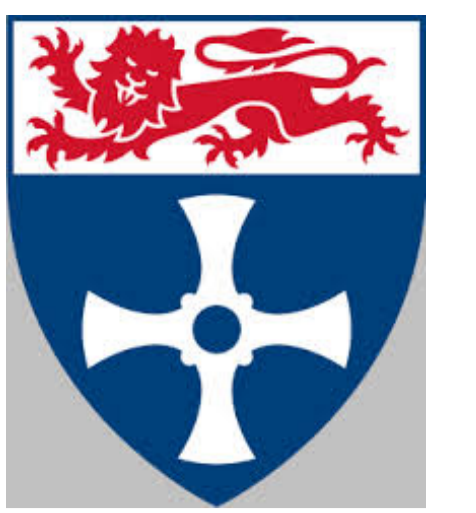




The Wigner Stochastic Gross-Pitaevskii Equation: a c-field theory that includes both thermal and quantum fluctuations



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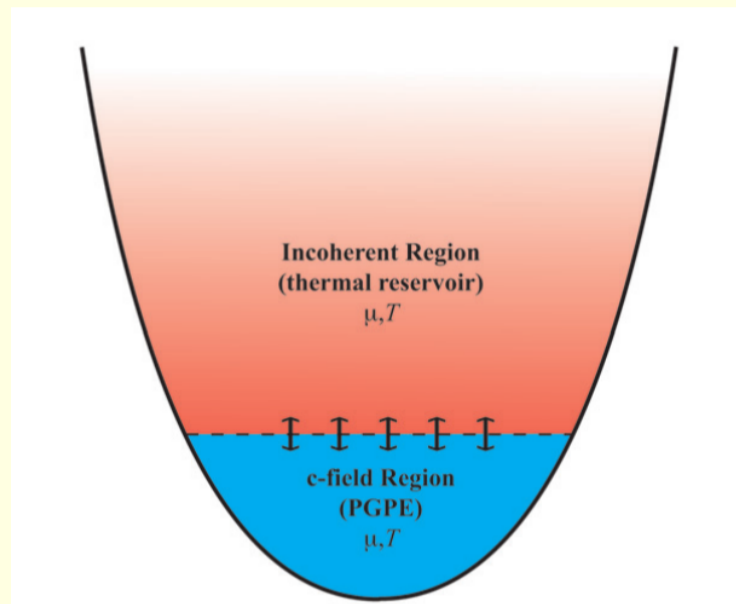


Aim: A c-field description that includes thermal and quantum fluctuations on an even footing
 To be used for: (*) Generating a thermal ensemble of single realizations
 (**) Calculating the quantum dynamics

Concept

- Re-derive SGPE equations from a Wigner representation of the Bose field + thermal bath, $\psi_W(\mathbf{x})$ but this time: without explicitly assuming high occupations
- Begin like Gardiner+Davis, *J. Phys. B* **36**, 4732 (2003).
- Still assume relatively high T : $\exp\left[\frac{\mathcal{L}_{GP}}{k_B T}\right] \approx 1 + \frac{\mathcal{L}_{GP}}{k_B T}$
- Obtain trajectories related to a "pre-SGPE" Fokker-Planck equation of Duine+Stoof, *PRA* **65**, 013603 (2001) that preserved fluctuation-dissipation relations, without assuming $N(\mathbf{k}) \gg 1$.
- To our knowledge, it has not been previously simulated.

The baseline: GPE $\hbar \frac{\partial \psi(\mathbf{x})}{\partial t} = -i\mathcal{L}_{GP}\psi(\mathbf{x})$ $\mathcal{L}_{GP} = -\frac{\hbar^2}{2m}\nabla^2 + V(\mathbf{x}) - \mu + g|\psi(\mathbf{x})|^2$



Blakie, Bradley, Davis, Ballagh, Gardiner, *Adv. Phys.* **57**, 363 (2008)

SGPE (Stochastic GP Equation)

$$\hbar \frac{\partial \psi(\mathbf{x})}{\partial t} = -i(1-i\gamma)\mathcal{L}_{GP}\psi(\mathbf{x}) + \sqrt{2\gamma\hbar k_B T}\eta(\mathbf{x}, t)$$

Stable \rightarrow equilibrium; No quantum fluctuations, assumes macroscopic occupation

Truncated Wigner

$$\hbar \frac{\partial \psi_W(\mathbf{x})}{\partial t} = -i\mathcal{L}_{GP}\psi_W(\mathbf{x})$$

Stable numerically, Quantum fluctuations at short time; Unclear crossover at long time to GPE with no quantum fluctuations

Positive P $\mathcal{L}_{PP} = -\frac{\hbar^2}{2m}\nabla^2 + V(\mathbf{x}) - \mu + g\tilde{\psi}(\mathbf{x})^*\psi(\mathbf{x})$ Świsłocki, Deuar, arXiv:1409.0146

$$\hbar \frac{\partial \psi(\mathbf{x})}{\partial t} = -i(1-i\gamma)\mathcal{L}_{PP}\psi(\mathbf{x}) + \sqrt{i\hbar g(1-2i\gamma)}\xi(\mathbf{x}, t)\psi(\mathbf{x}) + \sqrt{2\gamma\hbar k_B T}\eta(\mathbf{x}, t)$$

$$\hbar \frac{\partial \tilde{\psi}(\mathbf{x})}{\partial t} = -i(1-i\gamma)\mathcal{L}_{PP}^*\tilde{\psi}(\mathbf{x}) + \sqrt{i\hbar g(1-2i\gamma)}\tilde{\xi}(\mathbf{x}, t)\tilde{\psi}(\mathbf{x}) + \sqrt{2\gamma\hbar k_B T}\eta(\mathbf{x}, t)$$

Unstable numerically \rightarrow equilibrium not achievable; Full quantum mechanics while it lasts

Prior candidates

Expectation values

$$\langle \hat{\Psi}^\dagger(\mathbf{x})\hat{\Psi}(\mathbf{x}') \rangle = \langle \psi_W(\mathbf{x})^*\psi_W(\mathbf{x}') \rangle_{\text{ensemble}} - \frac{1}{2}\delta(\mathbf{x} - \mathbf{x}')$$

Symmetrically ordered moments

$$\langle \hat{\Psi}^\dagger(\mathbf{x})\hat{\Psi}^\dagger(\mathbf{x}')\hat{\Psi}(\mathbf{x})\hat{\Psi}(\mathbf{x}') \rangle = \langle |\psi_W(\mathbf{x})|^2|\psi_W(\mathbf{x}')|^2 \rangle_{\text{ensemble}} - \frac{1}{2}\left[\frac{1}{\Delta v} + \delta(\mathbf{x} - \mathbf{x}')\right] \langle |\psi_W(\mathbf{x})|^2 + |\psi_W(\mathbf{x}')|^2 \rangle_{\text{ensemble}} + \frac{1}{4\Delta v}\left[\frac{1}{\Delta v} + \delta(\mathbf{x} - \mathbf{x}')\right]$$

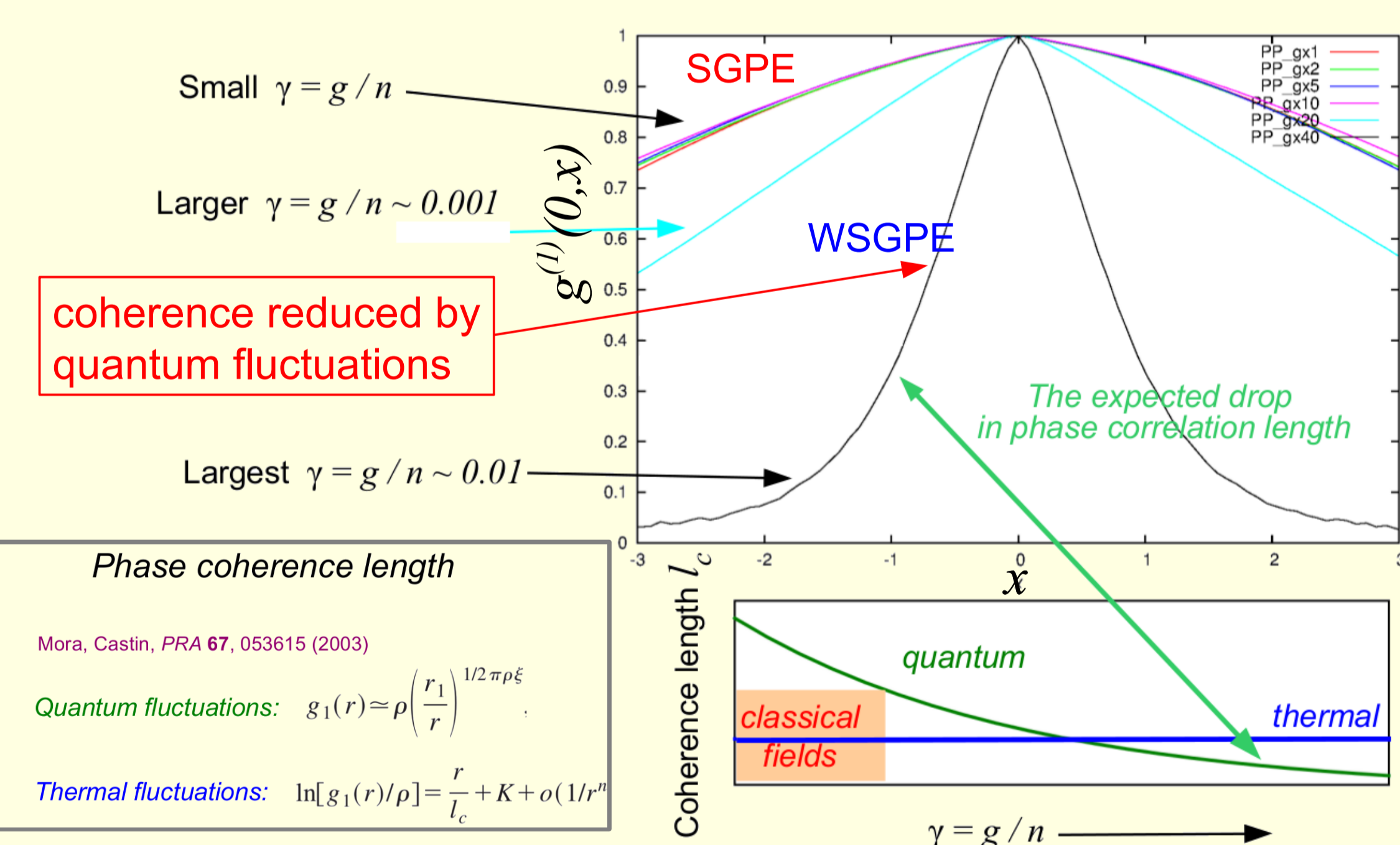
WSGPE evolution equation:

$$\hbar \frac{\partial \psi_W(\mathbf{x})}{\partial t} = -i(1-i\gamma)\left[\mathcal{L}_{GP} - \frac{g}{\Delta v}\right]\psi_W(\mathbf{x}) + \sqrt{\gamma\hbar\left[2k_B T + \mathcal{L}_{GP} - \frac{g}{\Delta v}\right]}\eta(\mathbf{x}, t) + i\sqrt{\gamma\hbar g}\xi(\mathbf{x}, t)\psi_W(\mathbf{x})$$

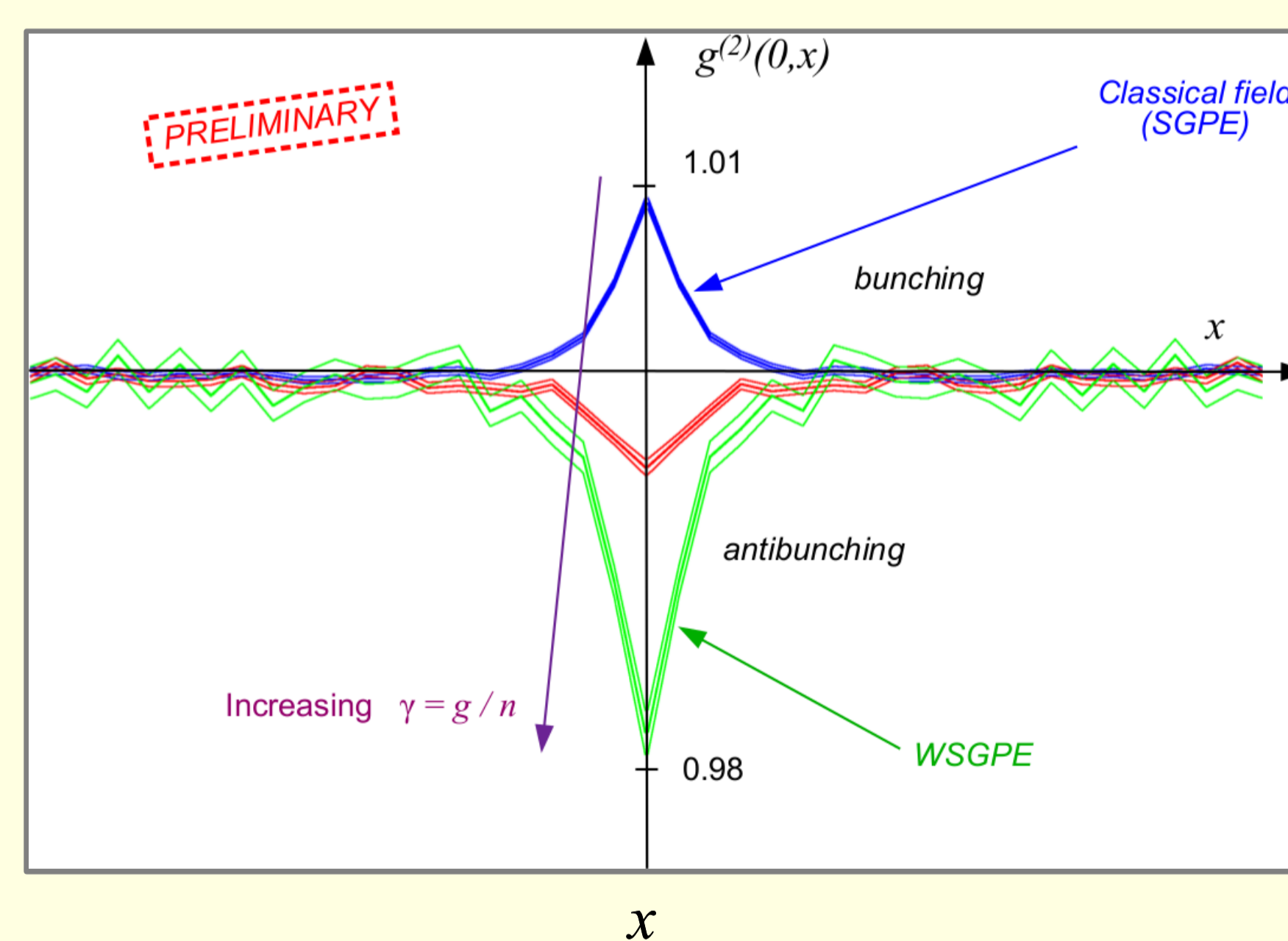
Extra thermal noise at high energy
 Complex thermal noise
 Strange positive-P-like term with real noise from bath coupling
 Terms with explicit dependence on particle number (not just gn)

TEST CASE: Trapped 1D Bose gas $\mu = 22.4$, $T = 139 = 0.16T_\phi$, $g=0.01 \rightarrow$ centrally $\gamma=4.5 \times 10^{-6}$ $\tau = 5.5 \times 10^{-5}$, $\tau/\sqrt{\gamma} = 0.026$ [cold quasicondensate]
 Then, we change g , and $T \sim 1/g$, to keep SGPE, μ and $\tau/\sqrt{\gamma}$ constant. Reach $\gamma \sim 0.007$ and $\tau \sim 0.0022$

Phase correlations

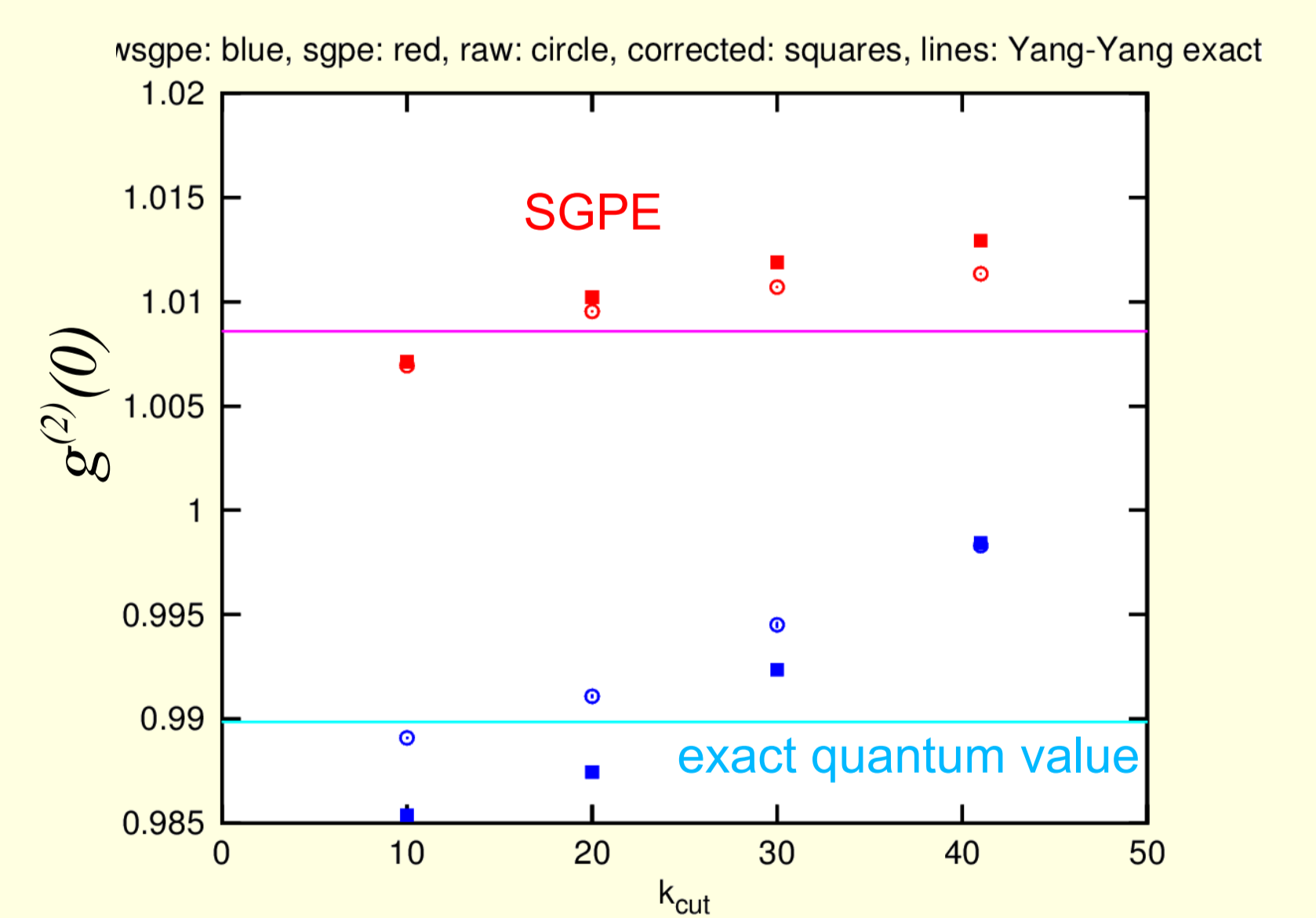


Antibunching

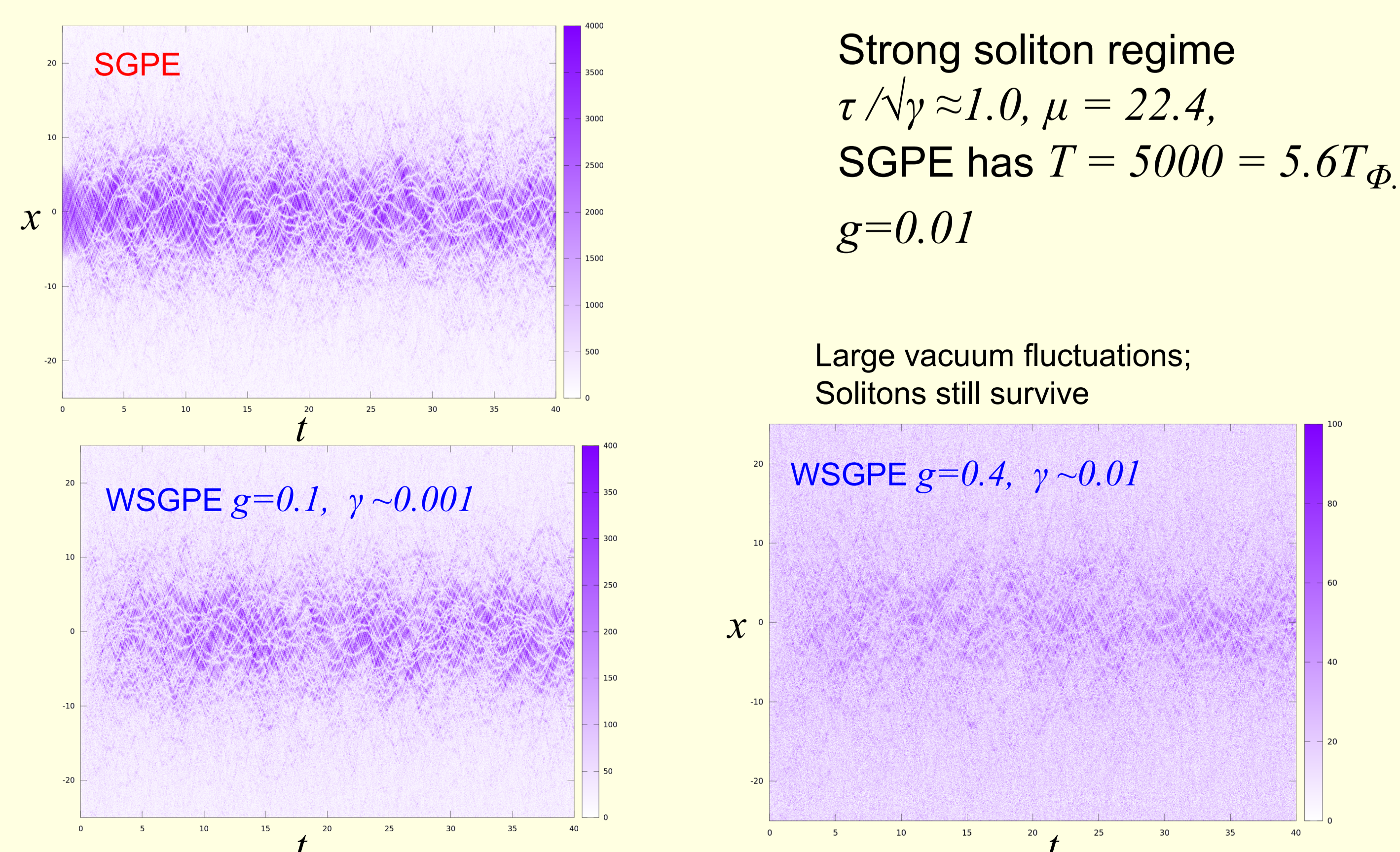


Interaction strength $\gamma = g/n$
 Relative temperature $\tau = k_B T / 4\pi T_d$

Cutoff dependence



Single shot dynamics

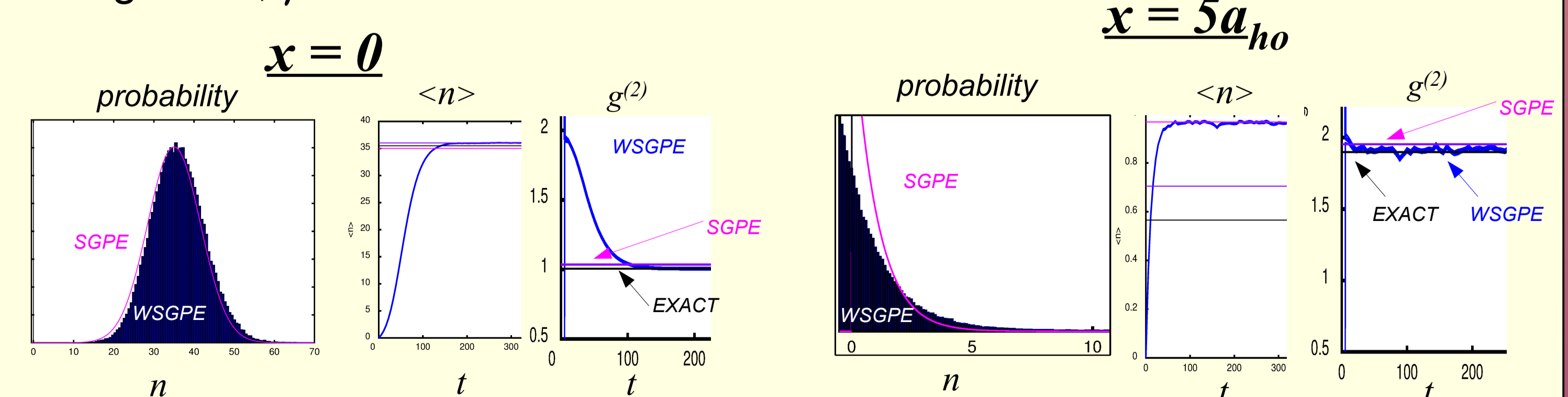


1 mode toy model

$$\mathcal{L}_{GP} = \frac{m\omega^2 x^2}{2} - \mu + g|\psi_W|^2$$

$$\hbar \frac{\partial \psi_W}{\partial t} = -i(1-i\gamma)\left[\mathcal{L}_{GP} - \frac{g}{\Delta x}\right]\psi_W + \sqrt{\gamma\hbar\left[2k_B T + \mathcal{L}_{GP} - \frac{g}{\Delta x}\right]}\eta(t) + i\sqrt{\gamma\hbar g}\frac{\xi(t)}{\Delta x}\psi_W$$

Long times, $\gamma = 1.1 \times 10^{-4}$ $\tau = 2.8 \times 10^{-4}$ in centre



Questions

- What distribution is reached in equilibrium?
- What do the multiplicative phase noise terms do?
- Is the cutoff dependence the same as in the SGPE?
- At what point does Wigner truncation affect results?
- A better form of the noise can be obtained that applies also at low $k_B T$