

s, p - d coupling in ZnO doped with 3d transition metal impurities

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Motivations

- The s, p d coupling between free carriers and the localized delectrons of the TM dopants constitutes the basic feature of diluted magnetic semiconductors. The magnetic properties of ZnO:TM are explored because of possible application in spintronics or solotronics [1].
- We study electronic and magnetic properties of ZnO:TM and determine dominant mechanisms of coupling with localized spins of

sp-d magnetic coupling

Exchange constants are calculated from

 $N_0 \alpha = \Delta \varepsilon_c / (x \langle S \rangle), \qquad N_0 \beta = \Delta \varepsilon_v / (x \langle S \rangle),$ where $\Delta \varepsilon_c$ and $\Delta \varepsilon_v$ - spin splittings of CBM and VBM, x - concentration of the TM ions, $\langle S \rangle = M/2$ - total magnetic moment.

 $N_0 \alpha$ and $N_0 \beta$

$N_0\beta$ constant and p-d coupling

VBM: p(O)+p(Zn)+d(TM)- kinetic exchange caused by strong p(O)-d(TM) hybridization - sp - d can be FM $(N_0\beta > 0)$ or AFM $(N_0\beta < 0)$



magnetic atoms for both electrons and holes [2].

DFT methods

The calculations are performed within the DFT within the GGA+Uapproximation [3]. We use the pseudopotential method implemented in the Quantum ESPRESSO code [4]. Analysis of a single TM impurity in ZnO is performed using $3 \times 3 \times 2$ supercells with 72 atoms.

U corrections for 3d(TM), 3d(Zn) and 2p(O) electrons:

The values of U corrections give a correct band structure of ZnO and lead to an excelent agreement with experimental optical properties of TM ions in ZnO [5].

TM levels in ZnO

Effect of intrashell Coulomb repulsion [6]: TM levels increase with increasing occupation



Figure 3: The exchange integrals $N_0\alpha$ and $N_0\beta$ for ZnO with TM in 2+ and 3+ charge state.

Spin density and xc perturbation In spin polarized system $\Delta n = n_{\uparrow} - n_{\downarrow} \neq 0$ and thus

 $\Delta V_{xc} = V_{xc\downarrow} - V_{xc\uparrow} \approx A(n_{\downarrow}^{1/3} - n_{\uparrow}^{1/3})$ $\Delta \varepsilon_c = \langle CBM \downarrow | V_{xc\downarrow} | CBM \downarrow \rangle - \langle CBM \uparrow | V_{xc\uparrow} | CBM \uparrow \rangle$ if $|CBM\uparrow\rangle = |CBM\downarrow\rangle$ then simply:

 $\Delta \varepsilon_c = \langle CBM | \Delta V_{xc} | CBM \rangle$

Both Δn and ΔV_{xc} dominated by d(Co)with contribution of p(O-nn)



Figure 1: (left) Spin splitting and CF splitting of TM level in ZnO. (right) The energy levels of TM in various charge state.



Figure 2: Energy levels of TM^{2+} ions in ZnO.



Figure 4: Two-dimensional plot of $n_{\uparrow}^{1/3} - n_{\downarrow}^{1/3}$ of ZnO with Co 2+ in the plane containing Co-O bonds.

$N_0\alpha$ constant and s-d coupling

CBM: s(O)+s(Zn)+s(TM) states - on-site direct exchange on TM ion (70%, Liu picture) and - on O-nn caused by p(O)-d(TM) hybridization (30%, new effect) - s - d coupling is always FM, $N_0 \alpha$ silimar for all ions



Figure 5: The wave functions squared of the CBM of $ZnO:Co^{2+}$.

Conclusions

• The s - d coupling between magnetic moment of TM ion and conduction band electrons is always FM and the exchange constant $N_0 \alpha \approx 0.5$ eV. It is caused by on-site direct exchange on TM ion as well as on nearest-neighbour O ions, which magnetic properties result from p(O)-d(TM) hybridization.

• In turn, the p - d coupling with the valence band holes strongly depends on the position of TM levels in the gap. For some ions like Fe, Co, Ni, the coupling changes the sign from FM $(N_0\beta > 0)$ to AFM $(N_0\beta < 0)$ going from 2+ to 3+ charge state.

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References

[1] F.Pan et al., Mat. Sci Eng. R **62**, 1 (2008). [2] A. Ciechan and P. Bogusławski, Sci. Reports **10**, 113848 (2021). [3] P. Hohenberg and W. Kohn, Phys. Rev. **136**, 864 (1964); J. P. Perdew, K. Burke and M. Ernzerhof, Phys. Rev. Lett. 77, 3865 (1996); M. Cococcioni and S. de Gironcoli, Phys. Rev. B **71**, 035105 (2005). [4] QUANTUM ESPRESSO, www.quantum-espresso.org. [5] A. Ciechan et al., Phys. Rev. B **94**, 165143 (2016); J. Papierska et al., Phys. Rev. B **94**, 224414 (2016); A. Ciechan and P. Bogusławski, J. Phys.: Condens. Matter **31** 255501 (2019); O. Volnianska and P. Bogusławski, J. Alloys Compd. 782, 1024 (2029). [6] A. Ciechan and P. Bogusławski, Opt. Mat. **79**, 264 (2018). [7] K. Ando et al., J. Appl. Phys. 89, 7284 (2001). [8] W. Pacuski et al., Phys. Rev. B 84, 035214 (2011). [9] W. Pacuski et al., Phys. Rev. B **73**, 035214 (2006).

Comparison with experiment

Ti, V, Cr experiment: $N_0 \alpha = N_0 \beta \approx 0$ [7] theory: Ti - non-magnetic; V, Cr - very weak coupling

Mn - good agreement experiment: $N_0 |\beta - \alpha| = (0.2 \pm 0.1) \text{ eV} [8]$ theory: $N_0 |\beta - \alpha| = 0.1 \text{ eV for } \text{Mn}^{2+}$

Co - close to experiment experiment: $N_0 |\beta - \alpha| = 0.8 \text{ eV} [9]$ theory: $N_0 |\beta - \alpha| = 2.6 \text{ eV for } \mathrm{Co}^{2+}$ $N_0 |\beta - \alpha| = 1.2 \text{ eV for } \text{Co}^{3+}$